

New prospects in synthesis and study of superheavy nuclei

- **Superheavy dreams of Walter Greiner**
- **Epoch of “cold” and “hot” fusion reactions ($100 < Z \leq 118$)**
- **Multi-nucleon transfer reactions**
- **Non-accelerative SHE production**
- **Summary**

Valery Zagrebaev and Walter Greiner

for XIII WONP – VII NURT, February 08, 2011, Havana, Cuba



JINR (*Dubna*)



FIAS (*Frankfurt*)

Mendeleev's Table 140 years ago

Система элементов

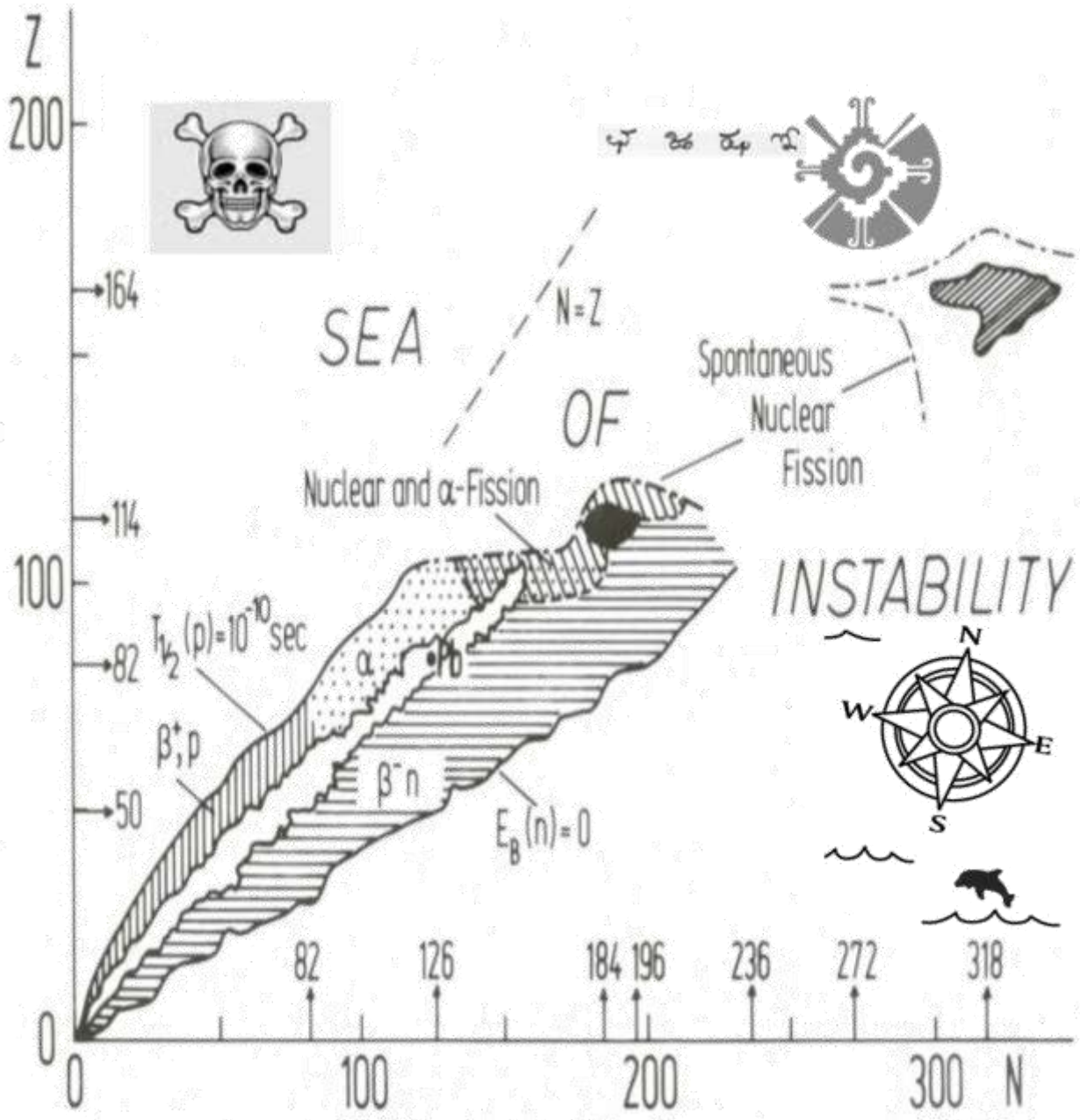
по атомному весу

и химическим свойствам

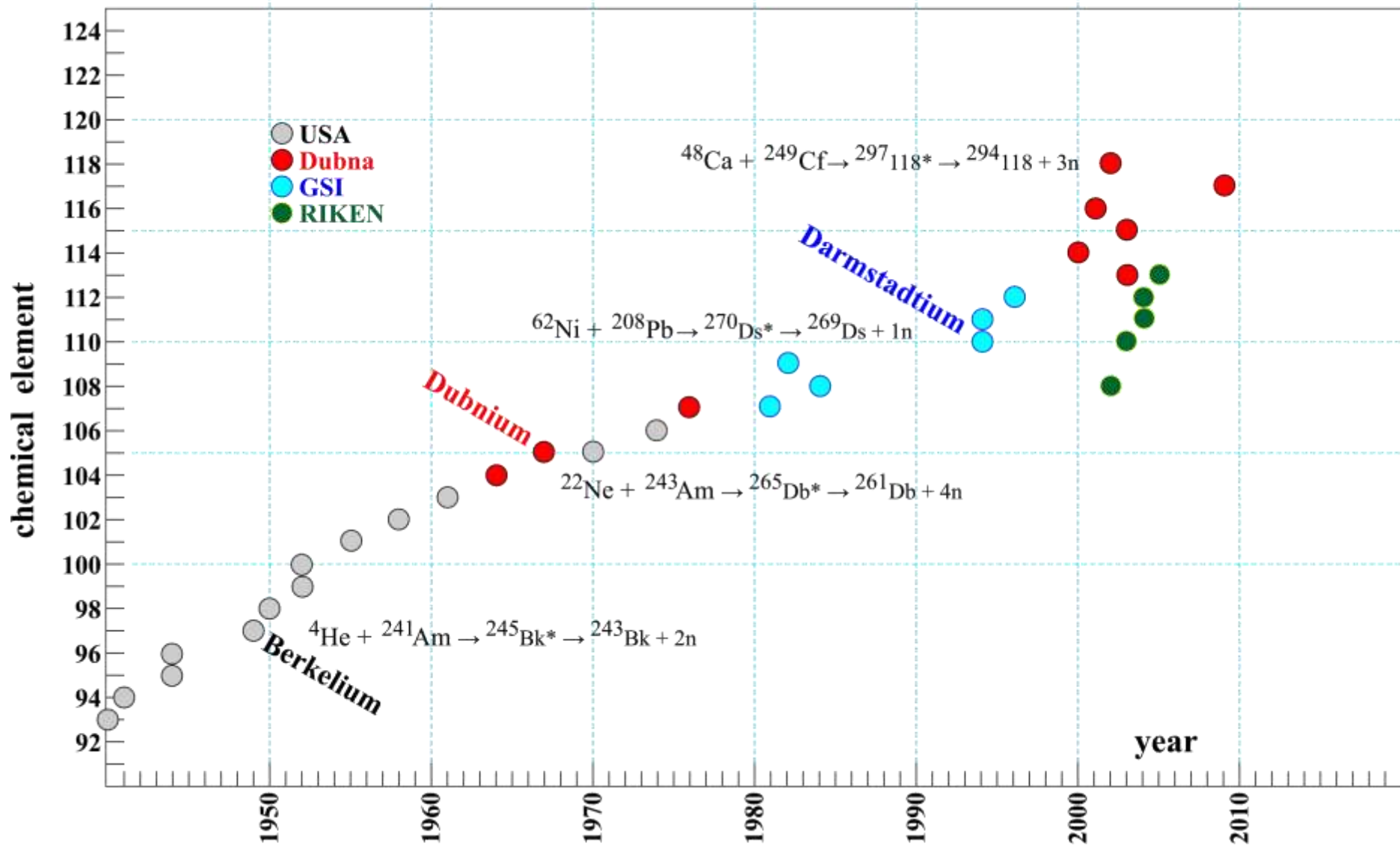
Д. Менделѣевъ.

			<i>Li=70</i>	<i>Be=90</i>	<i>B=100</i>
			<i>V=57</i>	<i>Nb=94</i>	<i>Ta=180</i>
			<i>C=50</i>	<i>Sn=76</i>	<i>W=186</i>
			<i>As=55</i>	<i>Hf=104</i>	<i>Rh=174</i>
			<i>Se=56</i>	<i>Ru=104</i>	<i>Ir=191</i>
			<i>K=39</i>	<i>Pd=106</i>	<i>Ag=179</i>
			<i>Ca=40</i>	<i>Cu=63</i>	<i>Zn=72</i>
			<i>Sc=44</i>	<i>Ni=58</i>	<i>Cd=112</i>
			<i>Ti=78</i>	<i>Co=59</i>	<i>Hg=200</i>
			<i>Zr=90</i>	<i>Ni=58</i>	<i>Pb=207</i>
			<i>Hf=104</i>	<i>As=75</i>	<i>Bi=208</i>
			<i>Ru=104</i>	<i>Sn=76</i>	<i>Po=210</i>
			<i>Rh=106</i>	<i>Pb=77</i>	<i>At=210</i>
			<i>Ir=191</i>	<i>Te=78</i>	<i>Fr=211</i>
			<i>Os=192</i>	<i>I=79</i>	<i>Ra=226</i>
			<i>Pt=195</i>	<i>X=80</i>	
			<i>Au=197</i>	<i>Y=81</i>	
			<i>Hg=200</i>	<i>Zr=82</i>	
			<i>Tl=203</i>	<i>Nb=83</i>	
			<i>Pb=207</i>	<i>Mo=84</i>	
			<i>Bi=208</i>	<i>Ru=85</i>	
			<i>Po=210</i>	<i>Rh=86</i>	
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			<i>Fr=211</i>	<i>Ag=88</i>	
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				<i>I=94</i>	
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				<i>Y=96</i>	
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				<i>Mo=99</i>	
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				<i>Te=108</i>	
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				<i>Y=111</i>	
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				<i>X=260</i>	
				<i>Y=261</i>	
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				<i>Rh=311</i>	
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				<i>Ag=313</i>	
				<i>Cd=314</i>	
				<i>In=315</i>	
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Ancient Nuclear Map and dreams of Walter Greiner



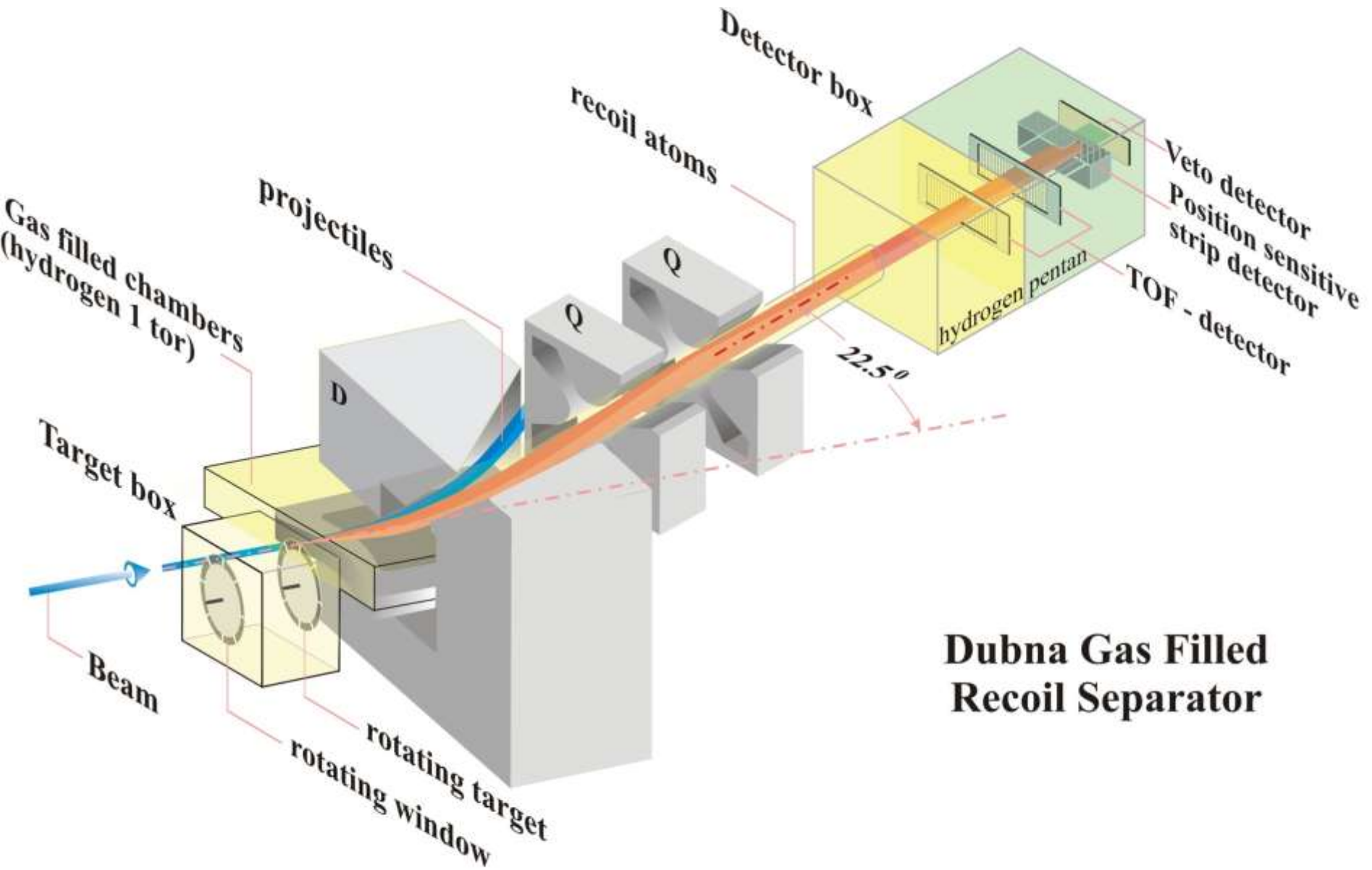
Synthesis of new elements (history)



Synthesis of superheavy elements (experiment)

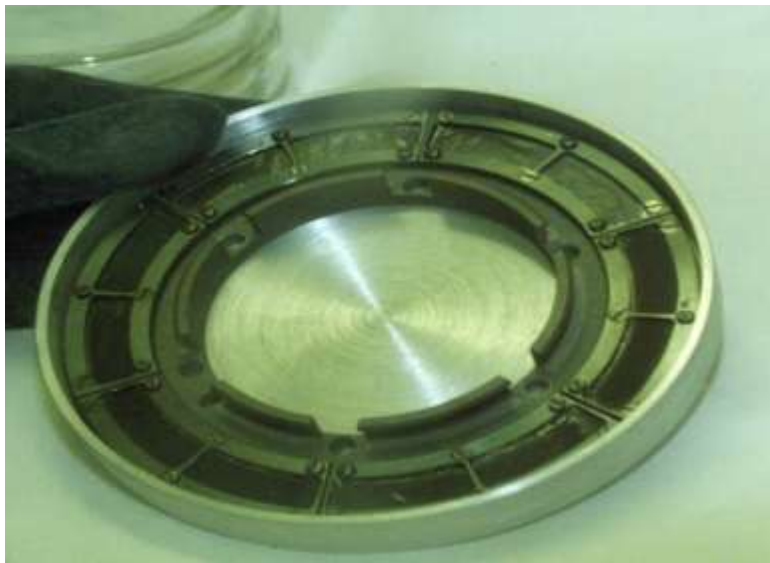


Synthesis of superheavy elements (experiment)

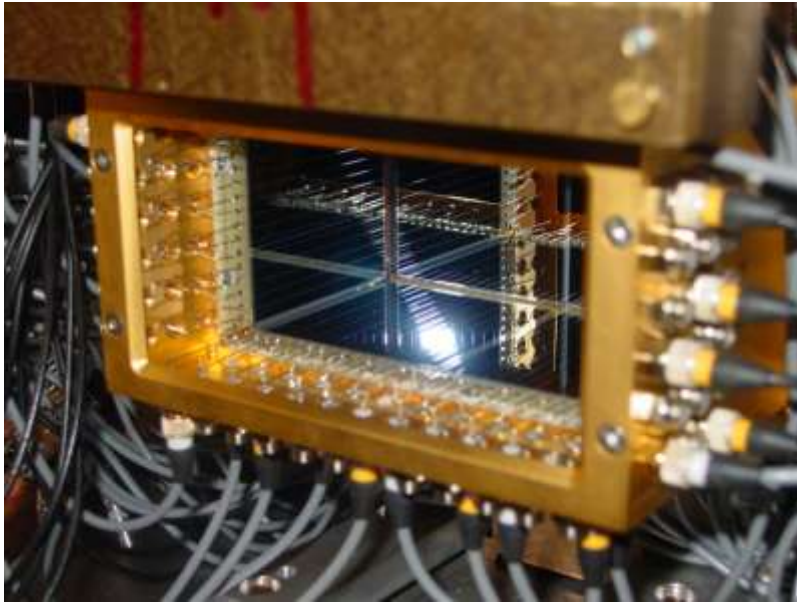


**Dubna Gas Filled
Recoil Separator**

Target made of ^{249}Cf (half-life is 350 years)



Detectors for SH nucleus recoil and products of its decay



Significance of Theory

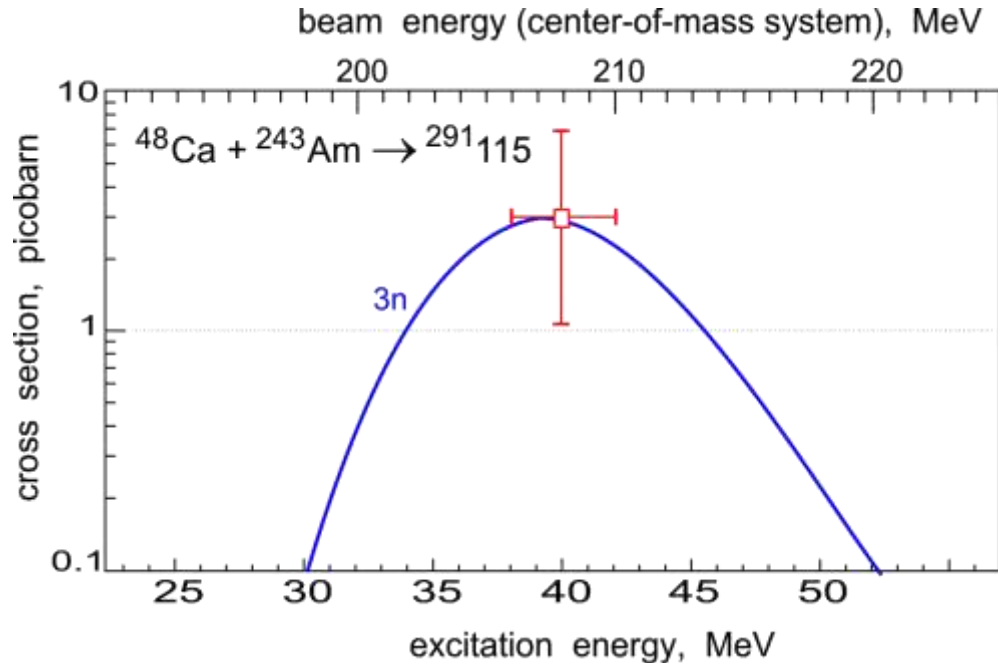
Which nuclei are to be fused ?

«cold» synthesis: $^{208}\text{Pb} + ^{64}\text{Ni}, ^{70}\text{Zn}, \dots \rightarrow ^{272}\text{110}, ^{278}\text{112}, \dots$ (GSI, Germany)

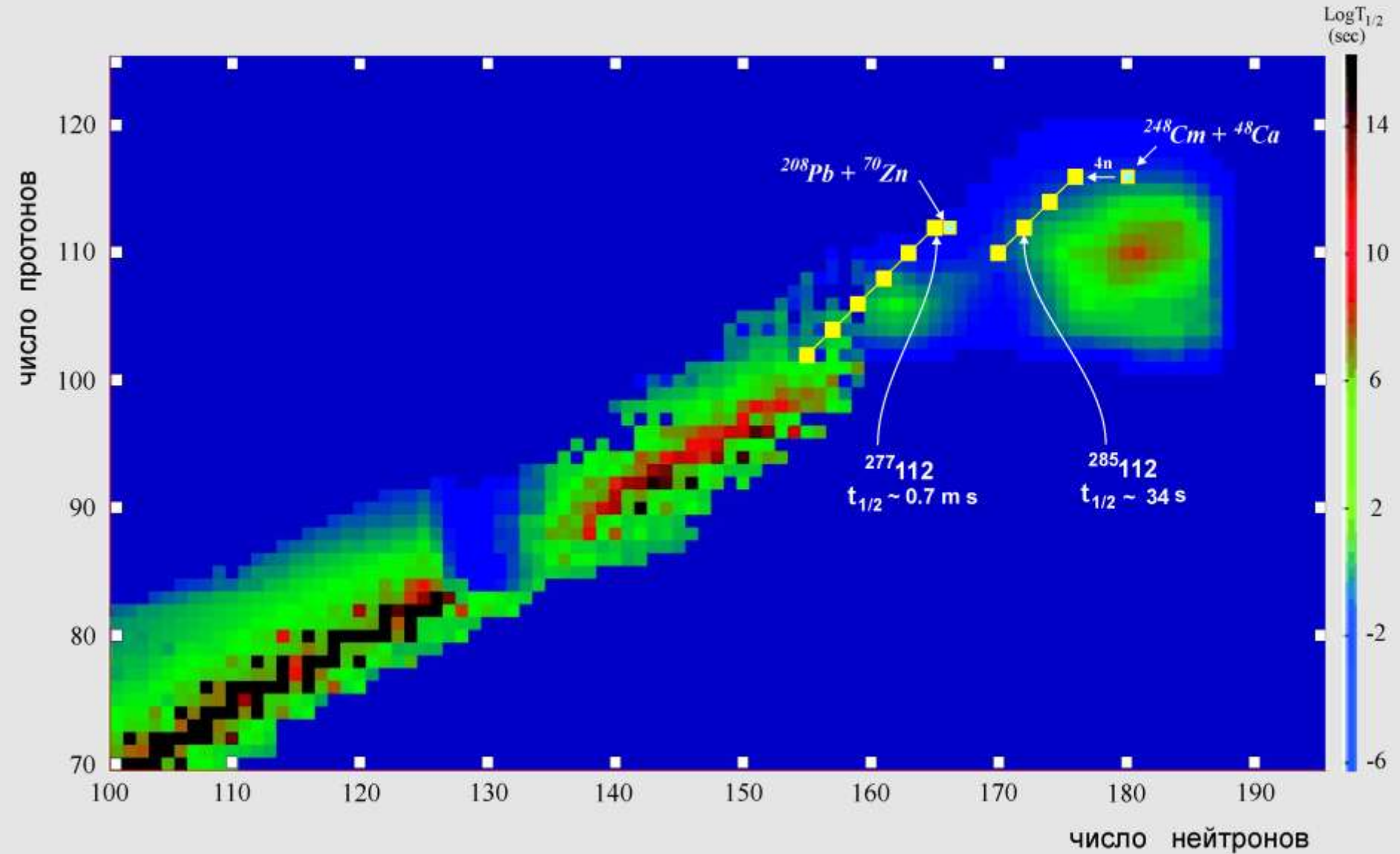
«hot» synthesis: $^{238}\text{U}, ^{244}\text{Pu}, ^{248}\text{Cm}, ^{249}\text{Cf} + ^{48}\text{Ca} \rightarrow ^{286}\text{112}, ^{292}\text{114}, ^{296}\text{116}, ^{297}\text{118}$ (Dubna)

Symmetric combinations: $^{148}\text{Nd} + ^{154}\text{Sm} \rightarrow ^{302}\text{122} ?$ Radioactive beams of ^{132}Sn , etc. ?

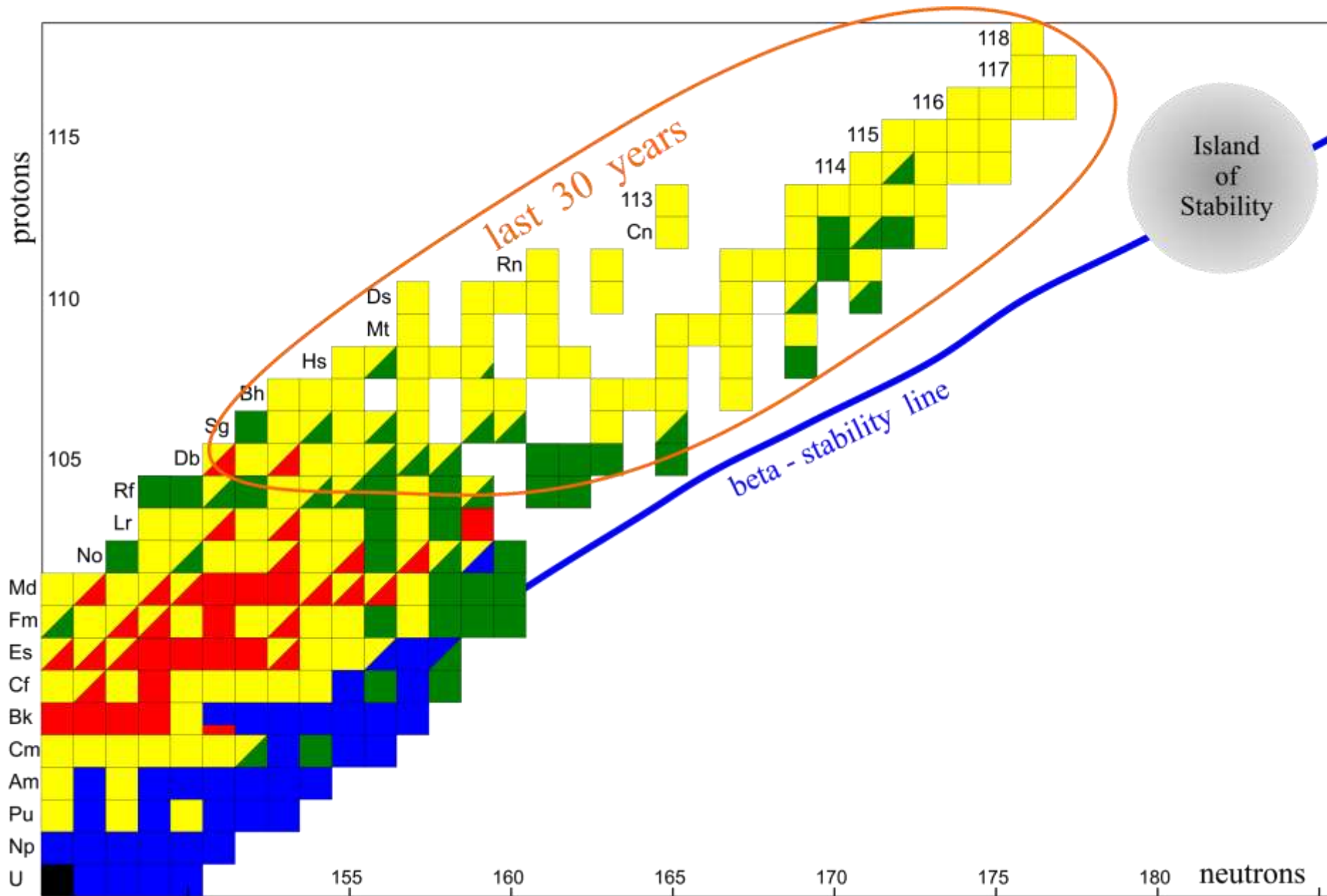
How to choose colliding energy ?



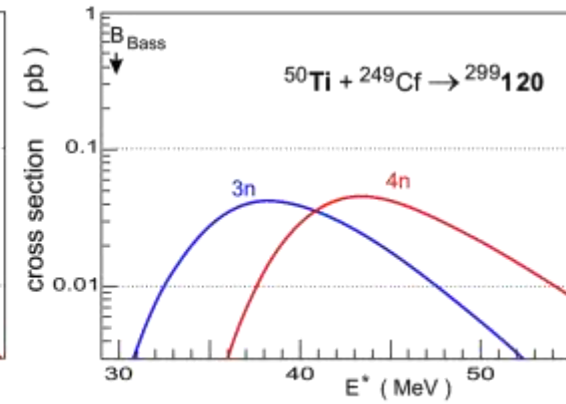
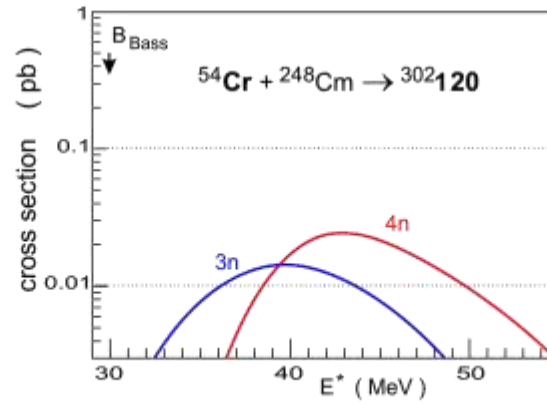
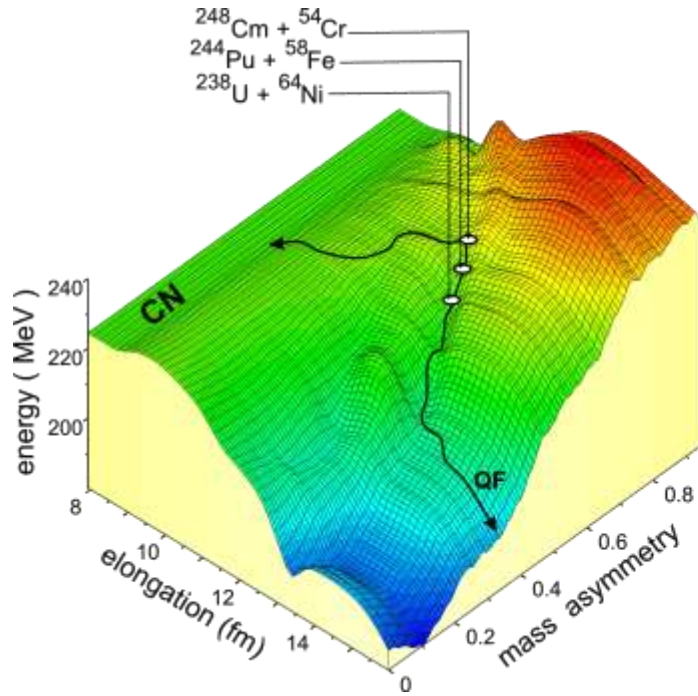
Superheavy Elements (Island of Stability)



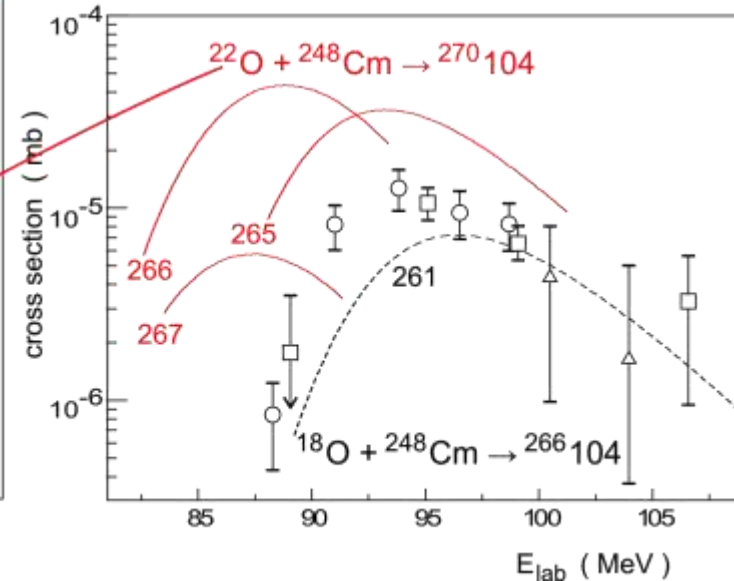
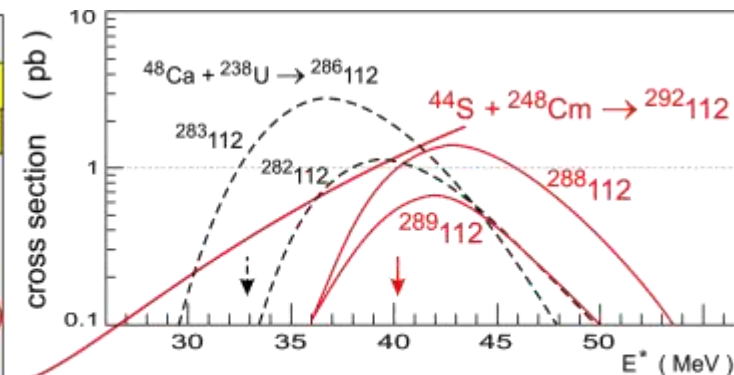
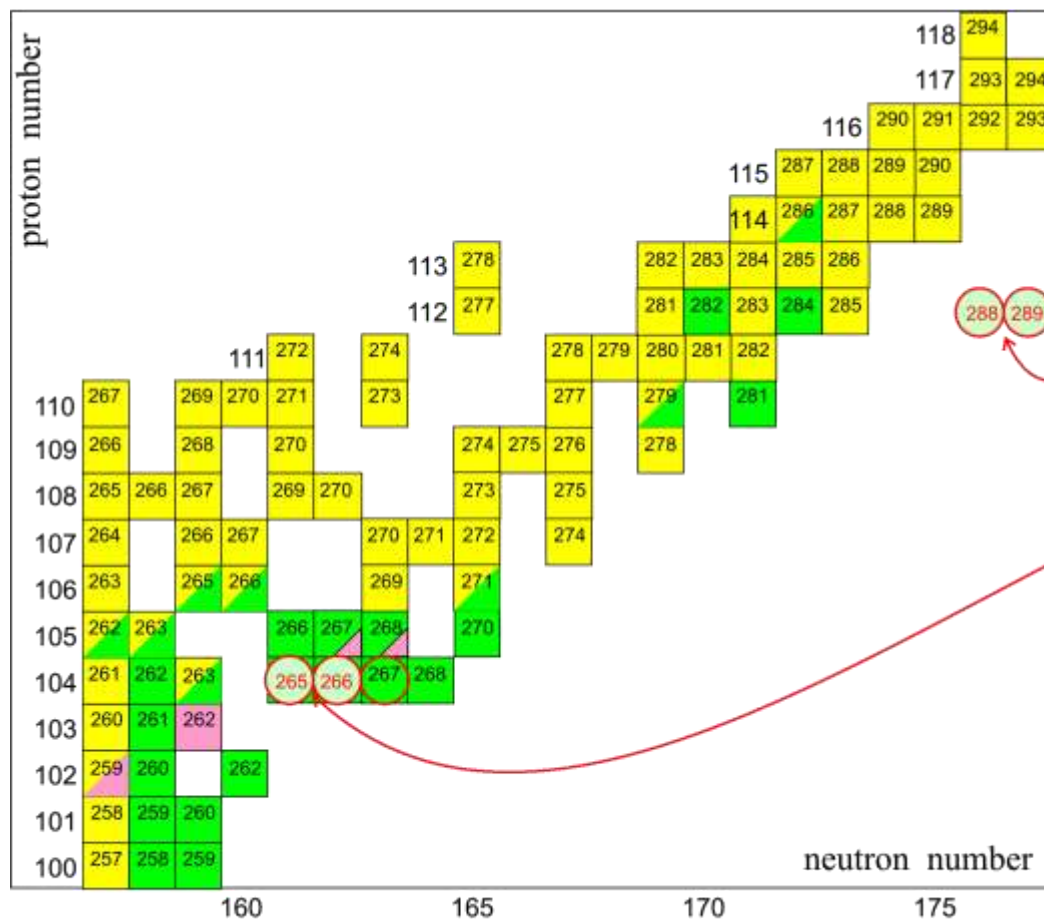
Great progress in synthesis of superheavy nuclei



Beyond ^{48}Ca : Pursuit of 120



Radioactive Ion Beams for production of neutron rich superheavy nuclei ?

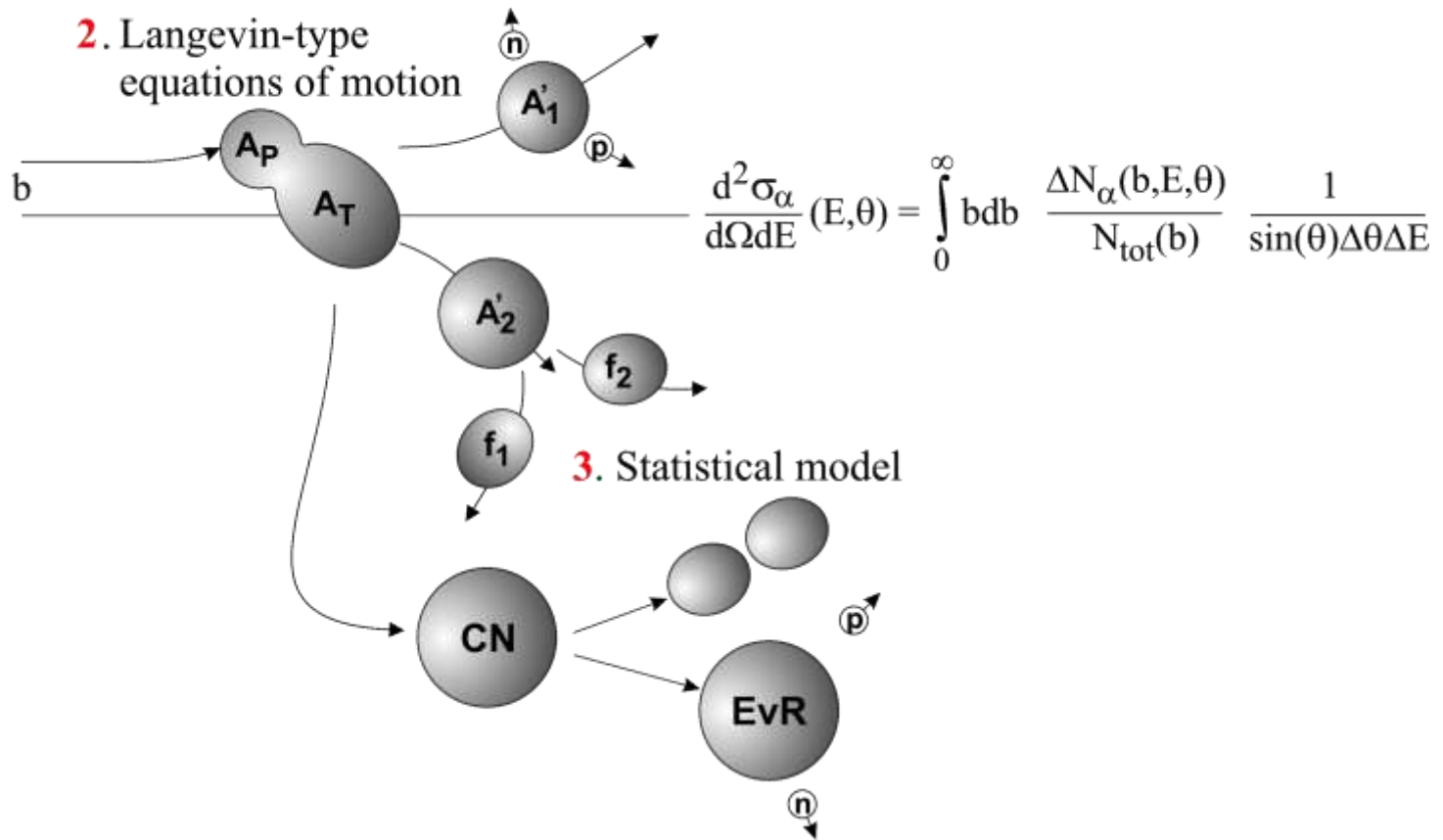


**Multi-nucleon transfer reactions
in low-energy heavy ion collisions**

Simulation of experiment and cross sections

1. Time-dependent driving potential $V(r, \xi; t)$:
Folding \rightarrow Adiabatic Two-Center Shell Model

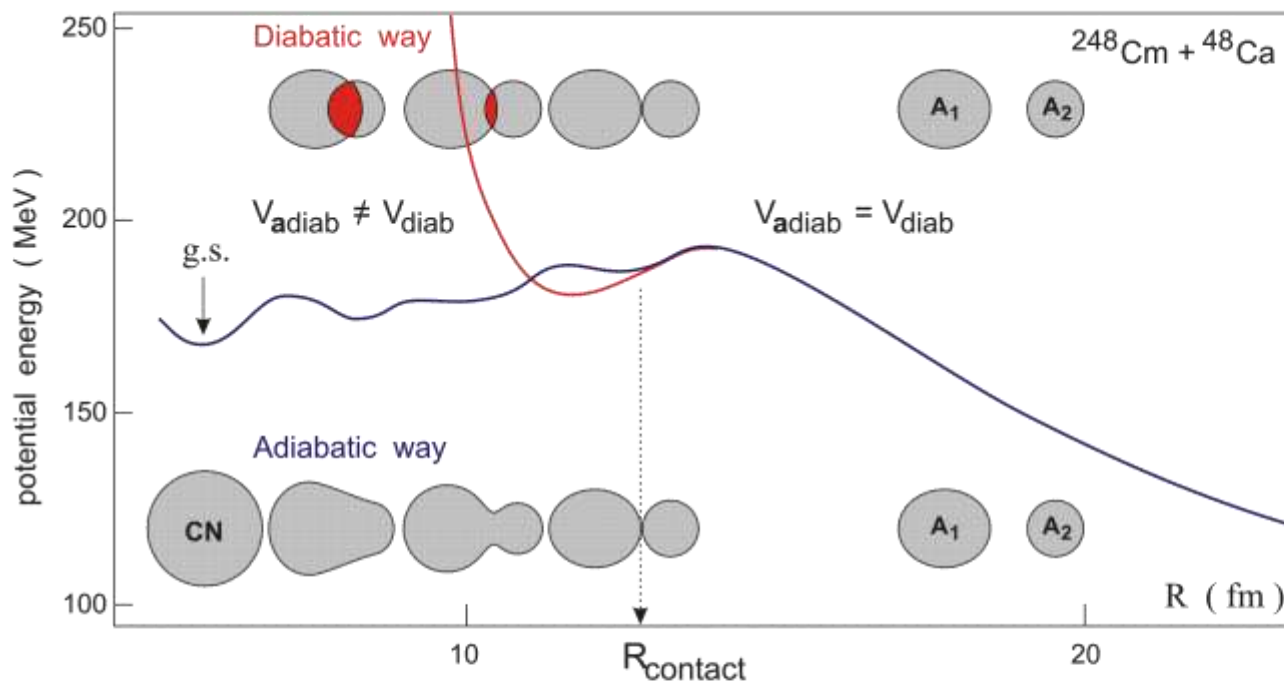
2. Langevin-type equations of motion



Dynamics: 10^6 tested events (trajectories),
 Statistical model: 10^{-6} ($3n$), 10^{-7} ($4n$) survival probability
 cross sections up to **0.1 pb** can be calculated

Time-dependent Driving Potential

$$V_{\text{diabat}}(R, \beta_1, \beta_2, \alpha, \dots) = V_{12}^{\text{folding}}(Z_1, N_1, Z_2, N_2; R, \beta_1, \beta_2, \dots) + M(A_1) + M(A_2) - M(\text{Proj}) - M(\text{Targ})$$



$$V_{\text{adiabat}}(R, \beta_1, \beta_2, \eta, \dots) = M_{\text{TCSM}}(R, \beta_1, \beta_2, \eta, \dots) - M(\text{Proj}) - M(\text{Targ})$$

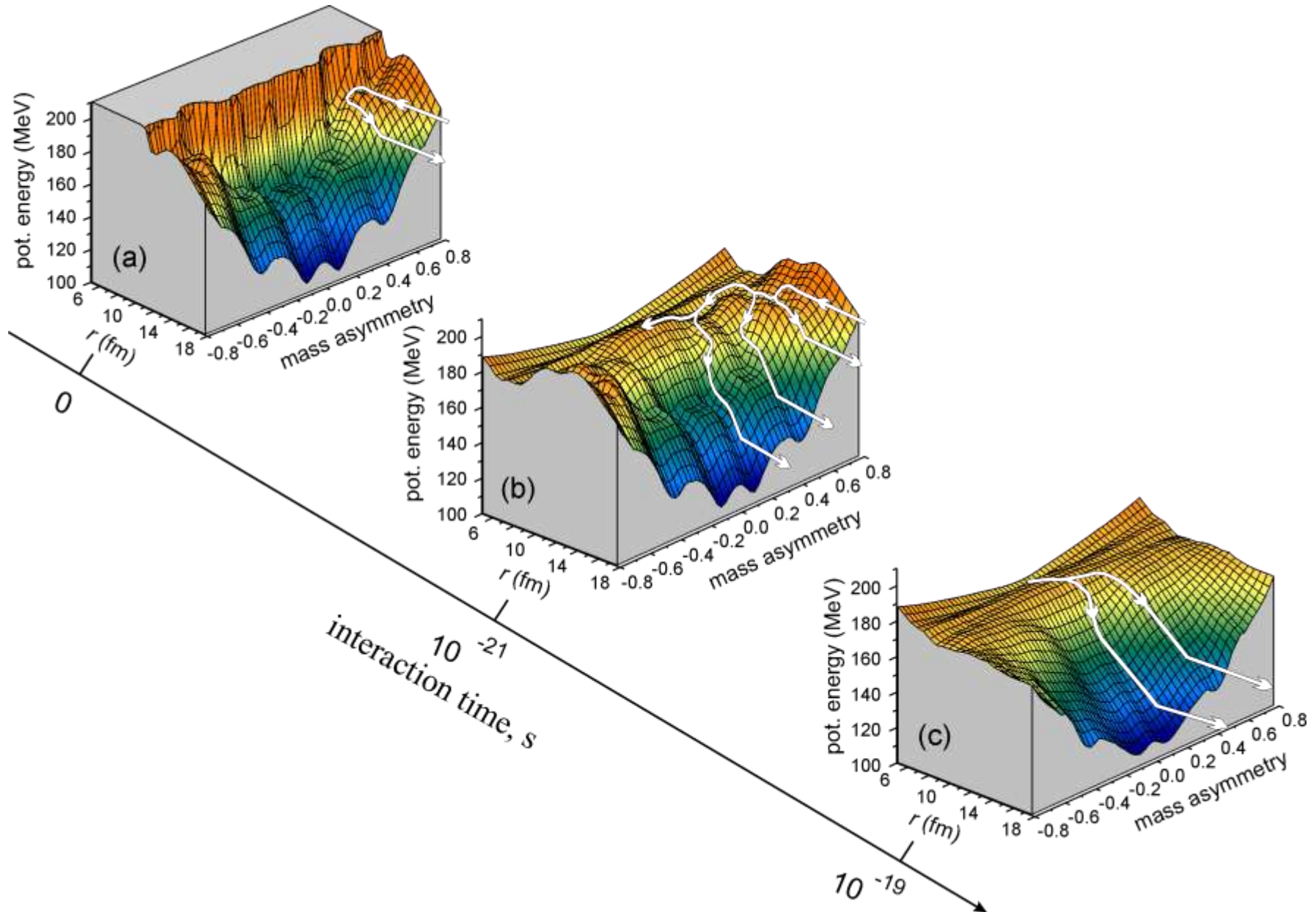
Time -dependent driving potential has to be used

$$V(t) = V_{\text{diab}}(\xi) \cdot \exp\left(-\frac{t_{\text{int}}}{\tau_{\text{relax}}}\right) + V_{\text{adiab}}(\xi) \cdot \left[1 - \exp\left(-\frac{t_{\text{int}}}{\tau_{\text{relax}}}\right)\right]$$

$$\tau_{\text{relax}} \sim 10^{-21} \text{ s}$$

*the same degrees of freedom ($\xi = R, \theta, \varphi_1, \varphi_2, \beta_1, \beta_2, \eta_Z, \eta_N$) !
All forces, $F_i(t) = -\partial V / \partial \xi_i$, are quite smooth*

Time-dependent Driving Potential





Nucleon Exchange

(L. Moretto, 1974)

Distribution function $\varphi(A_1, t) \rightarrow$ Master equation $\frac{\partial \varphi}{\partial t} = \sum_{A_1' = A_1 \pm 1} \lambda(A_1' \rightarrow A_1) \cdot \varphi(A_1') - \lambda(A_1 \rightarrow A_1') \cdot \varphi(A_1)$

$$\frac{\partial \varphi}{\partial t} = -\frac{\partial}{\partial A_1} (D^{(1)} \varphi) + \frac{\partial^2}{\partial A_1^2} (D^{(2)} \varphi) \quad \text{Fokker - Planck (W. Nörenberg, 1974)}$$

$$\eta = \frac{A_1 - A_2}{A_{CN}} = \frac{A_1 - (A_{CN} - A_1)}{A_{CN}} = \frac{2A_1 - A_{CN}}{A_{CN}}$$

$$\frac{d\eta}{dt} = \frac{2}{A_{CN}} D_A^{(1)} + \frac{2}{A_{CN}} \sqrt{D_A^{(2)}} \Gamma(t)$$

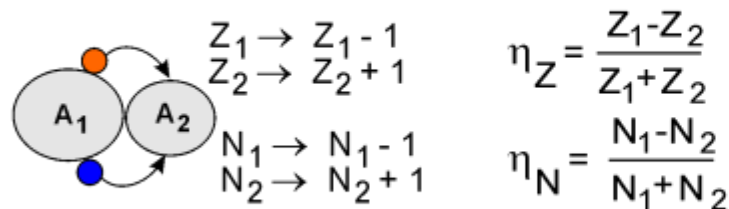
$$\frac{dA_1}{dt} = D^{(1)} + \sqrt{D^{(2)}} \Gamma(t) \quad \text{Langevin type eq.}$$

at $A' = A \pm 1$

$$D^{(1)} = \lambda(A_1 \rightarrow A_1 + 1) - \lambda(A_1 \rightarrow A_1 - 1)$$

$$D^{(2)} = \frac{1}{2} [\lambda(A_1 \rightarrow A_1 + 1) + \lambda(A_1 \rightarrow A_1 - 1)]$$

transition probability $\lambda^{(\pm)} = \lambda_0 \sqrt{\frac{\rho(A \pm 1)}{\rho(A)}} P_{tr}(R; A \rightarrow A \pm 1), \quad \rho \sim \exp(2\sqrt{aE^*}), \quad E^* = E_{c.m.} - V(R, \beta_1, \beta_2, \eta)$



$$\eta_Z = \frac{Z_1 - Z_2}{Z_1 + Z_2}$$

$$\eta_N = \frac{N_1 - N_2}{N_1 + N_2}$$

$$D_{N,Z}^{(1)} = \lambda_{N,Z}(A \rightarrow A + 1) - \lambda_{N,Z}(A \rightarrow A - 1)$$

$$D_{N,Z}^{(2)} = \frac{1}{2} [\lambda_{N,Z}(A \rightarrow A + 1) + \lambda_{N,Z}(A \rightarrow A - 1)]$$

$$\lambda_{N,Z}^{(\pm)} = \lambda_{N,Z}^0 \sqrt{\frac{\rho(A \pm 1)}{\rho(A)}} P_{tr}(R; A \rightarrow A \pm 1)$$

System of coupled Langevin type Equations of Motion

$$\frac{dR}{dt} = \frac{p_R}{\mu_R}$$

$$\frac{d\vartheta}{dt} = \frac{\ell}{\mu_R R^2}$$

$$\frac{d\varphi_1}{dt} = \frac{L_1}{\mathfrak{I}_1}, \quad \frac{d\varphi_2}{dt} = \frac{L_2}{\mathfrak{I}_2}$$

$$\frac{d\beta_1}{dt} = \frac{p_{\beta_1}}{\mu_{\beta_1}}$$

$$\frac{d\beta_2}{dt} = \frac{p_{\beta_2}}{\mu_{\beta_2}}$$

$$\frac{d\eta_Z}{dt} = \frac{2}{Z_{CN}} D_Z^{(1)} + \frac{2}{Z_{CN}} \sqrt{D_Z^{(2)}} \Gamma_Z(t)$$

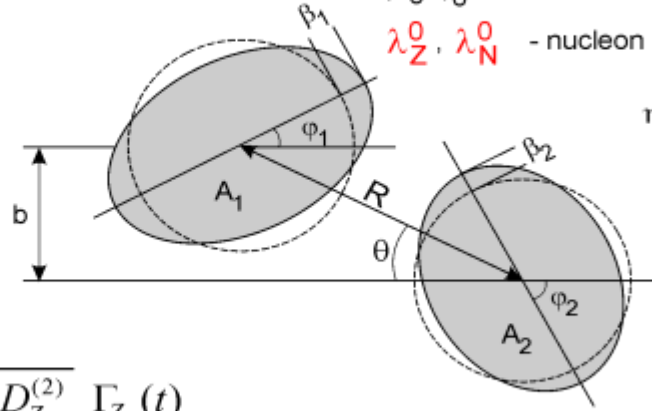
$$\frac{d\eta_N}{dt} = \frac{2}{N_{CN}} D_N^{(1)} + \frac{2}{N_{CN}} \sqrt{D_N^{(2)}} \Gamma_N(t)$$

Variables: $\{R, \theta, \varphi_1, \varphi_2, \beta_1, \beta_2, \eta_Z, \eta_N\}$

Most uncertain parameters:

μ_0, γ_0 - nuclear viscosity and friction,

λ_Z^0, λ_N^0 - nucleon transfer rate



$$\eta = \frac{A_1 - A_2}{A_1 + A_2}$$

$$\eta_Z = \frac{Z_1 - Z_2}{Z_1 + Z_2}$$

$$\eta_N = \frac{N_1 - N_2}{N_1 + N_2}$$

$$\lambda_Z^0 = \lambda_N^0 = \frac{\lambda^0}{2}$$

$$\frac{dp_R}{dt} = -\frac{\partial V}{\partial R} + \frac{\ell^2}{\mu_R R^3} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial R} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial R} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial R} - \gamma_R \frac{p_R}{\mu_R} + \sqrt{\gamma_R T} \Gamma_R(t)$$

$$\frac{d\ell}{dt} = -\frac{\partial V}{\partial \vartheta} - \gamma_{\text{tang}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) R + \sqrt{\gamma_{\text{tang}} T} \Gamma_{\text{tang}}(t)$$

$$\frac{dL_1}{dt} = -\frac{\partial V}{\partial \varphi_1} + \gamma_{\text{tang}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) a_1 - \frac{a_1}{R} \sqrt{\gamma_{\text{tang}} T} \Gamma_{\text{tang}}(t)$$

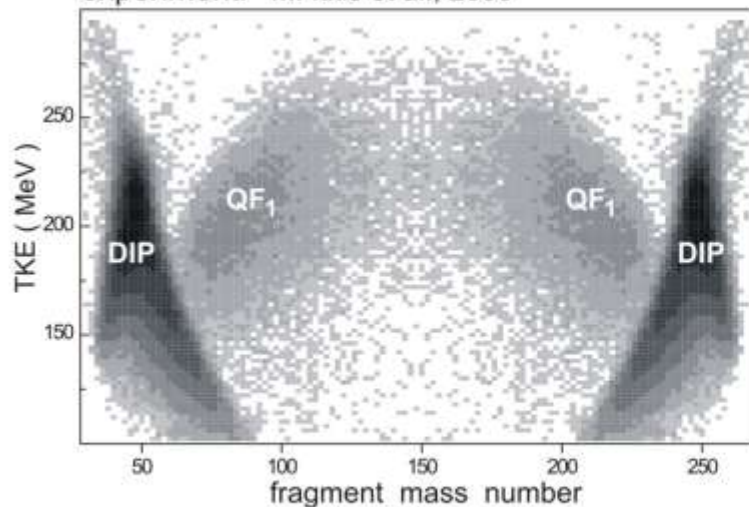
$$\frac{dL_2}{dt} = -\frac{\partial V}{\partial \varphi_2} + \gamma_{\text{tan}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) a_2 - \frac{a_2}{R} \sqrt{\gamma_{\text{tang}} T} \Gamma_{\text{tang}}(t)$$

$$\frac{dp_{\beta_1}}{dt} = -\frac{\partial V}{\partial \beta_1} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_1} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_1} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial \beta_1} - \gamma_{\beta} \frac{p_{\beta_1}}{\mu_{\beta_1}} + \sqrt{\gamma_{\beta_1} T} \Gamma_{\beta_1}(t)$$

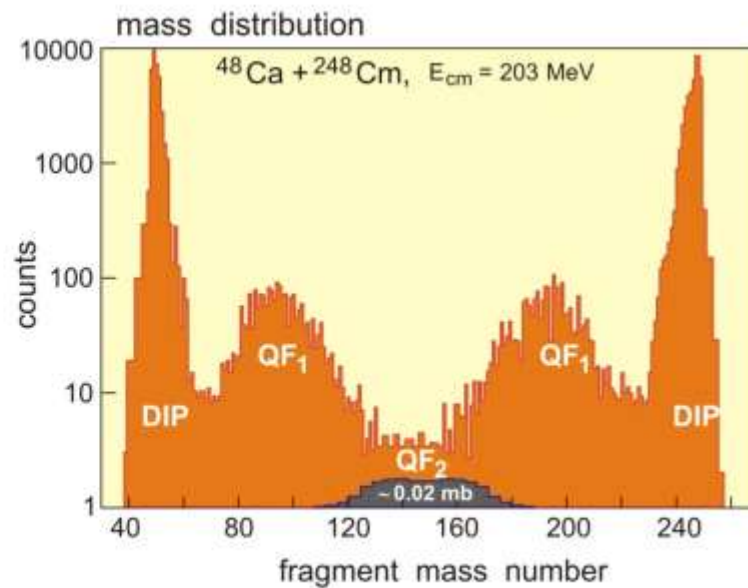
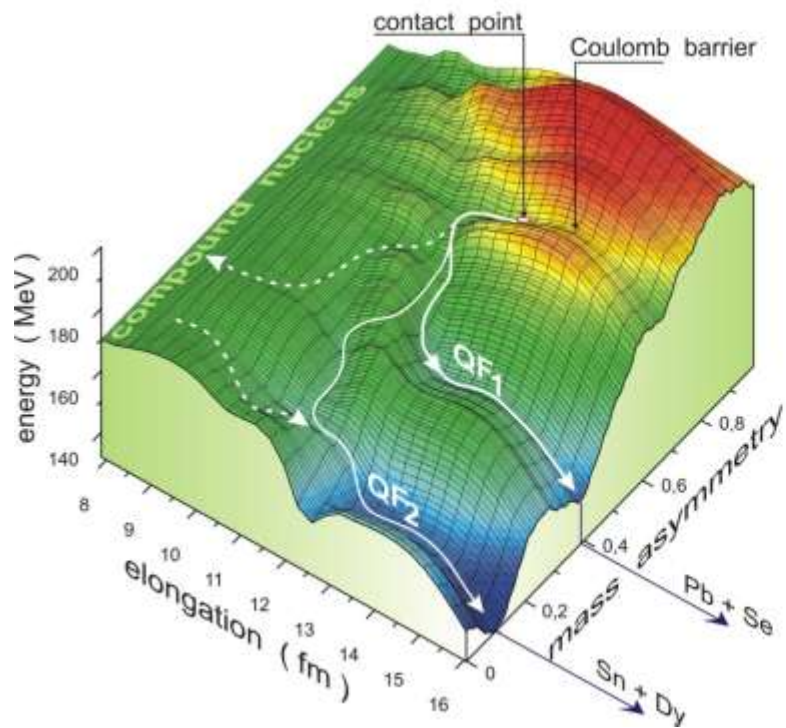
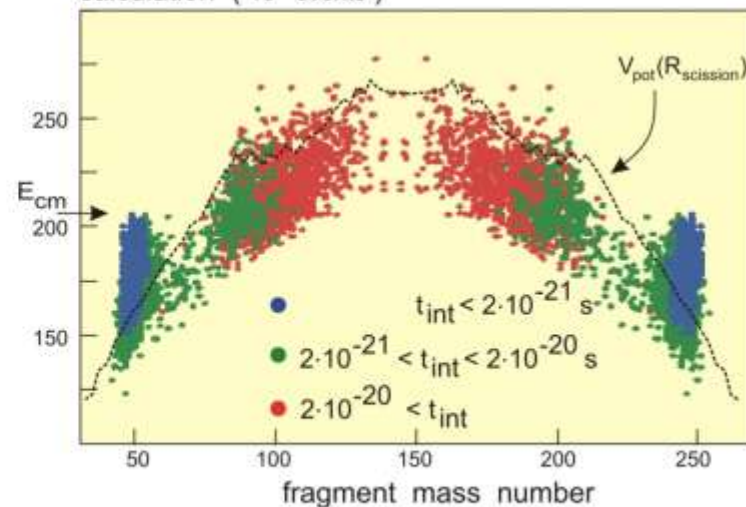
$$\frac{dp_{\beta_2}}{dt} = -\frac{\partial V}{\partial \beta_2} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_2} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_2} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial \beta_2} - \gamma_{\beta} \frac{p_{\beta_2}}{\mu_{\beta_2}} + \sqrt{\gamma_{\beta_2} T} \Gamma_{\beta_2}(t)$$

Good agreement with experiment: e.g. $^{48}\text{Ca} + ^{248}\text{Cm}$

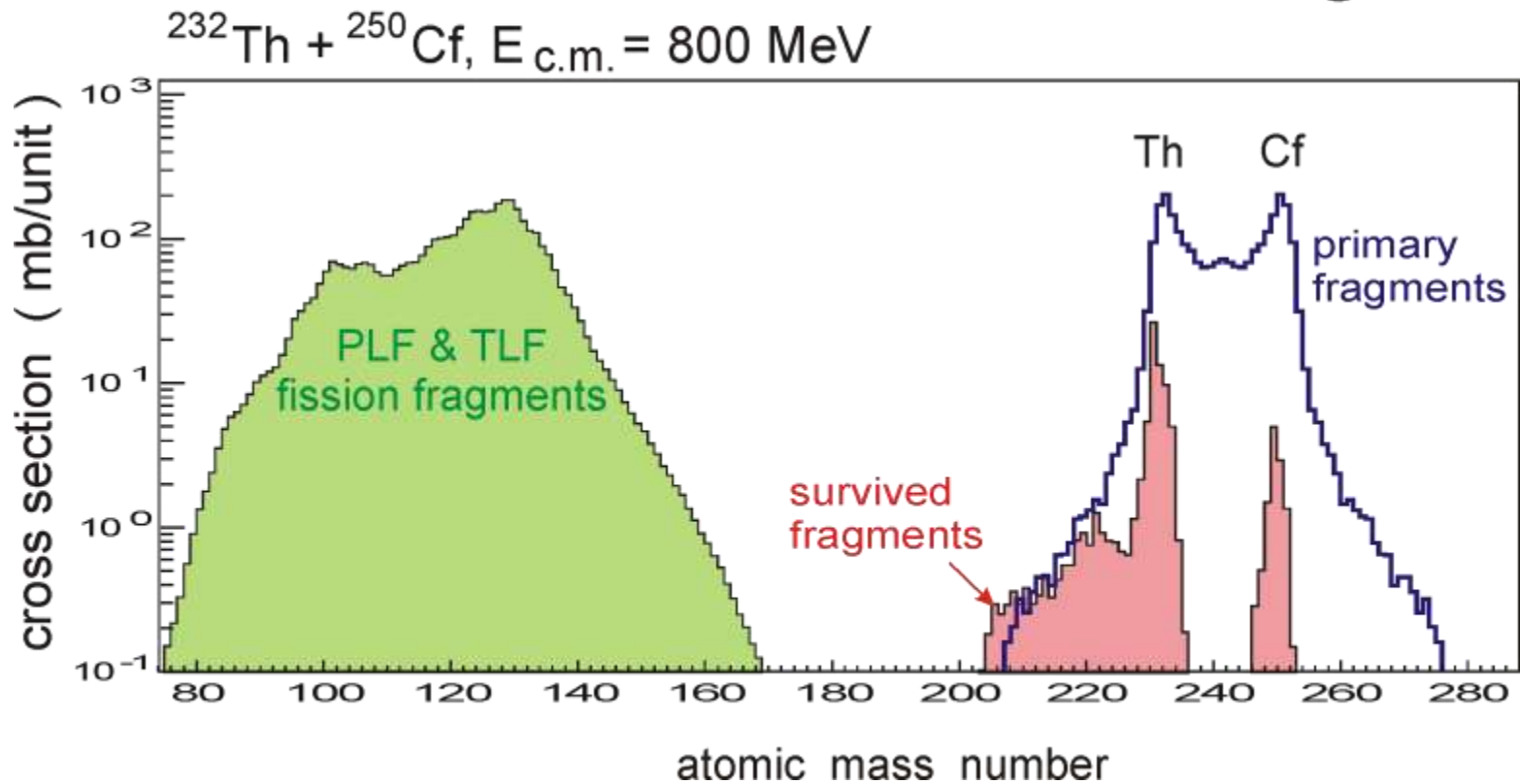
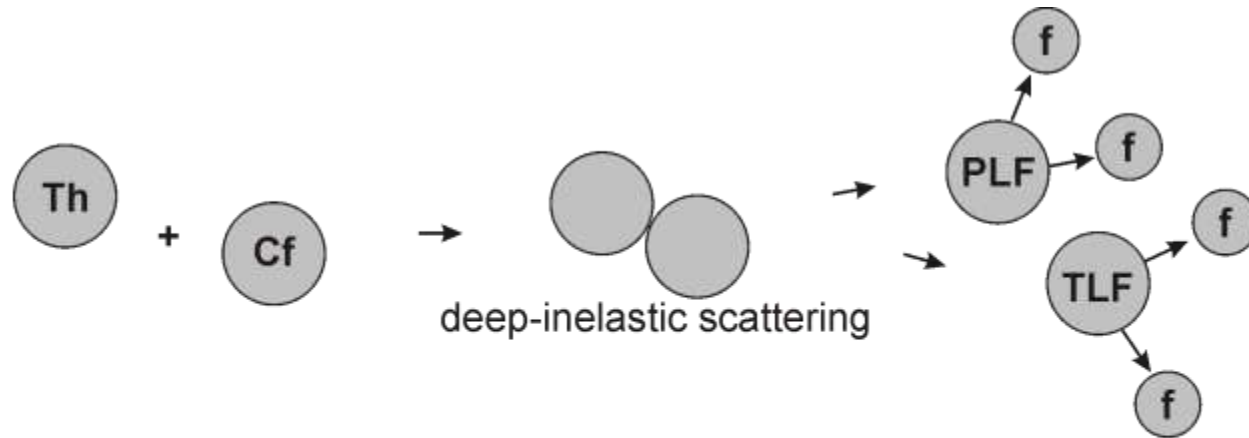
experiment: M. Itkis et al., 2000



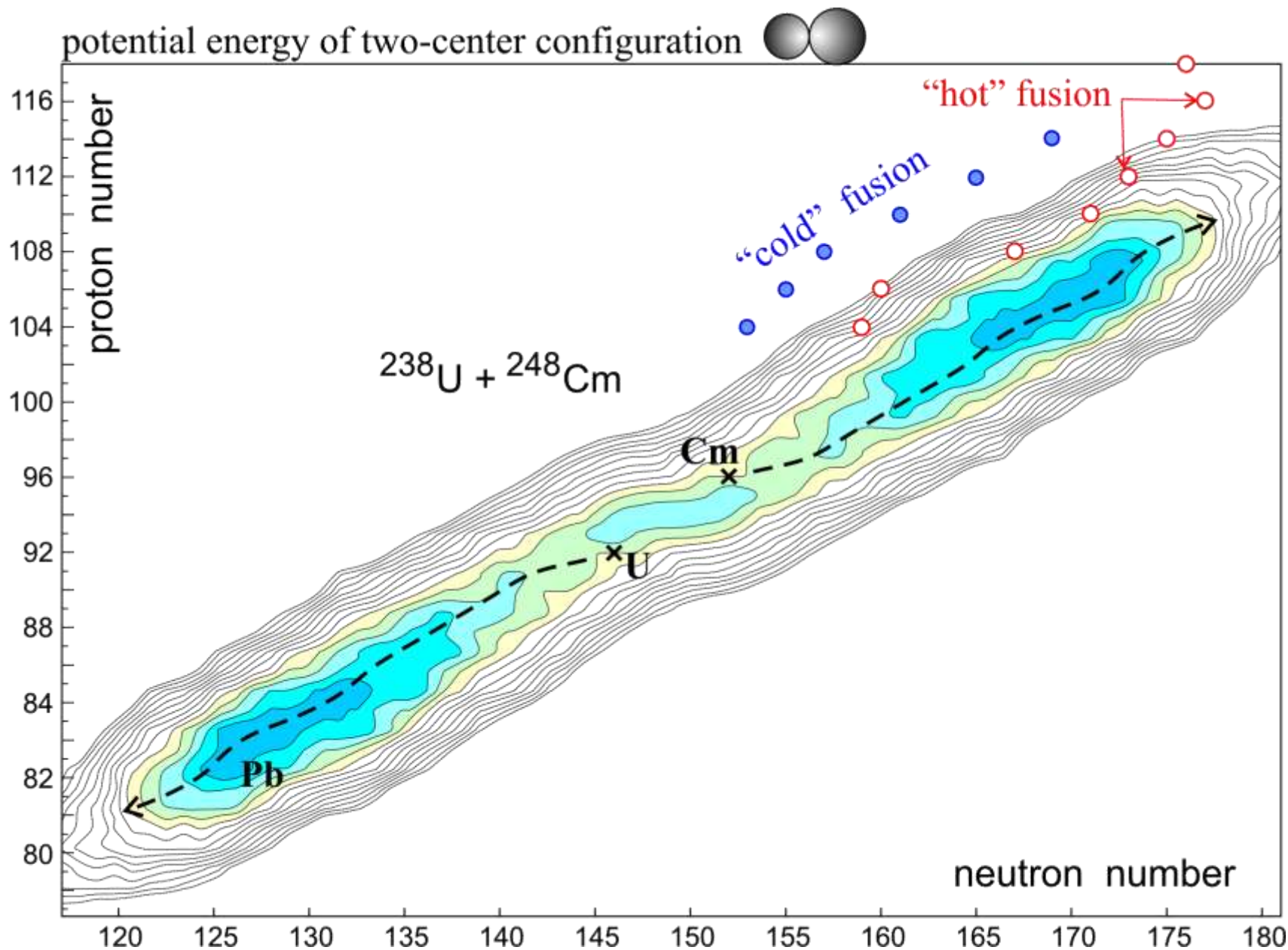
calculation (10^5 events)



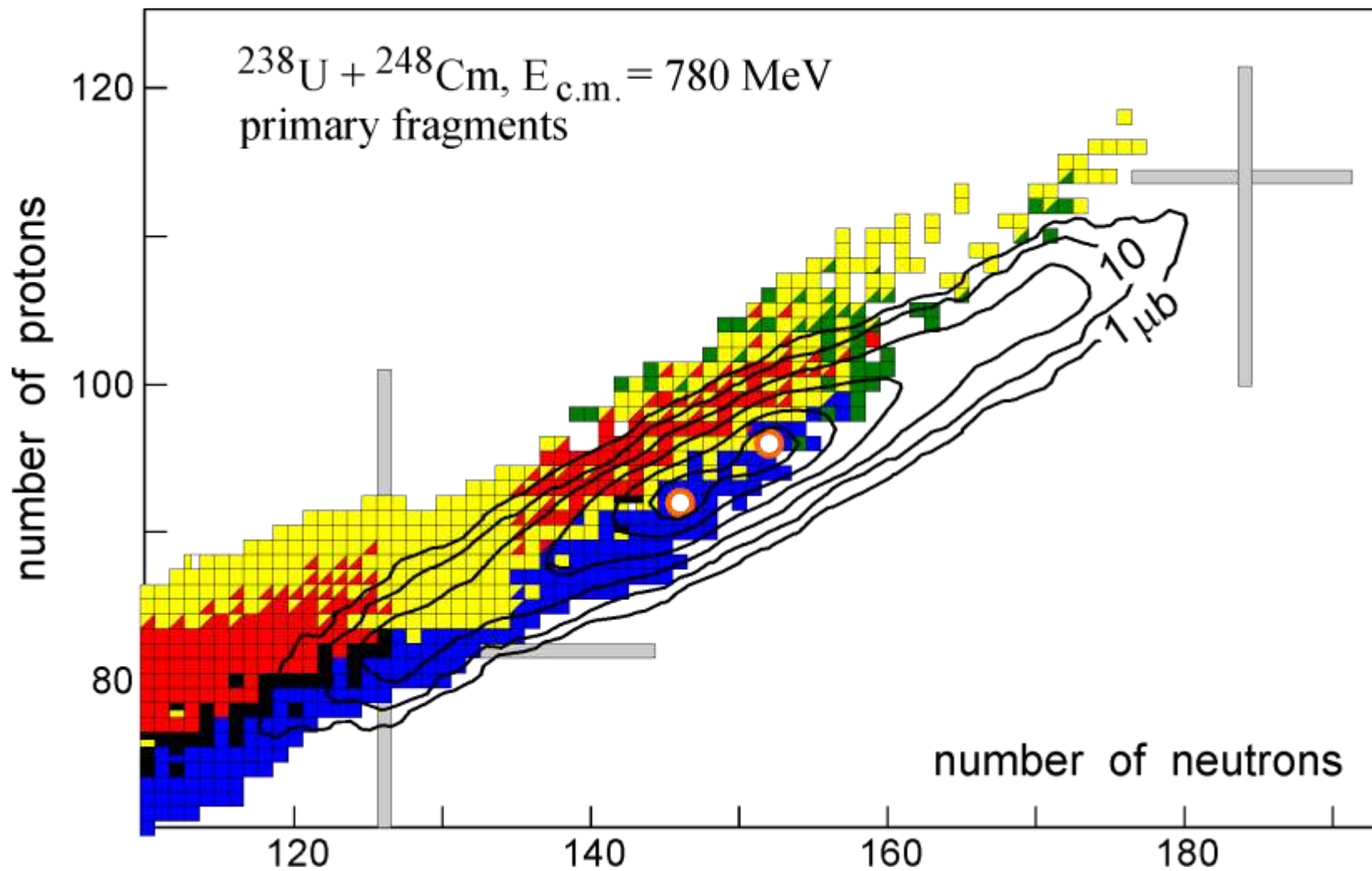
Transfer reactions in damped collision of very heavy nuclei ?



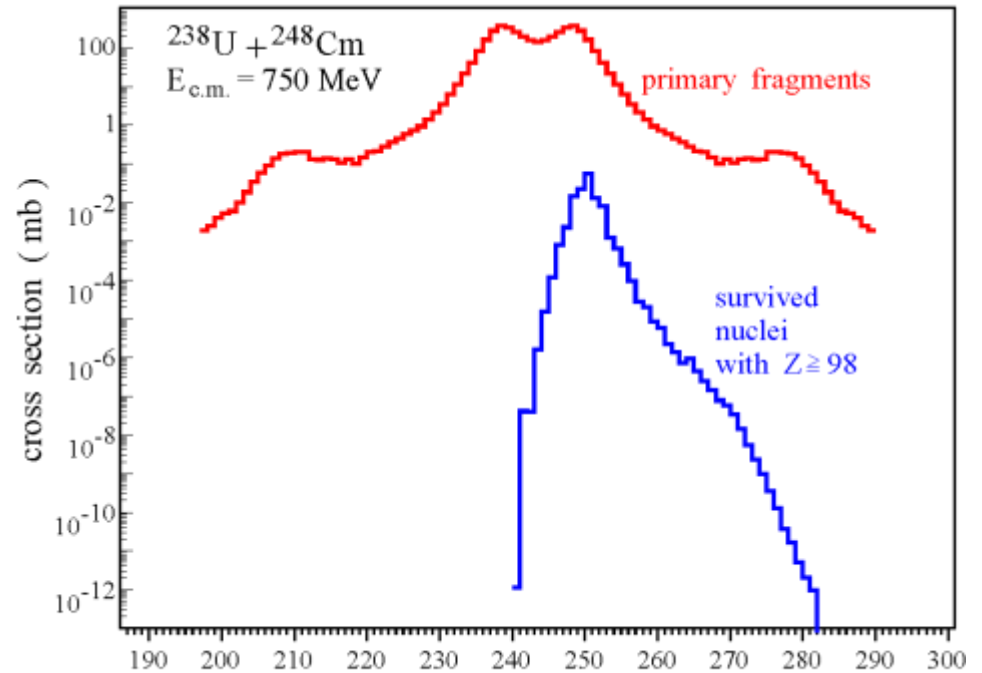
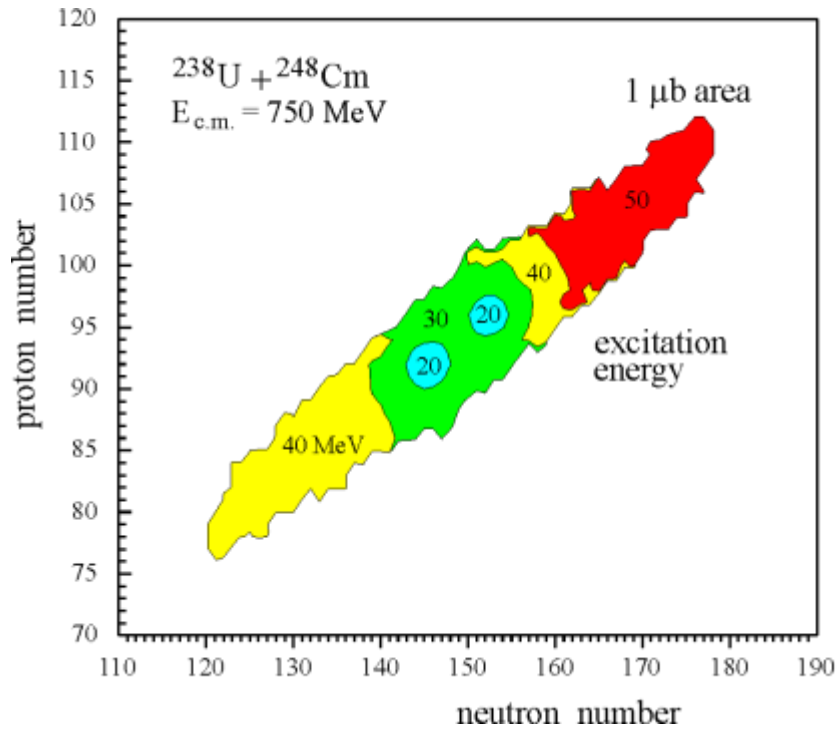
Most probable way of evolution of the giant nuclear system



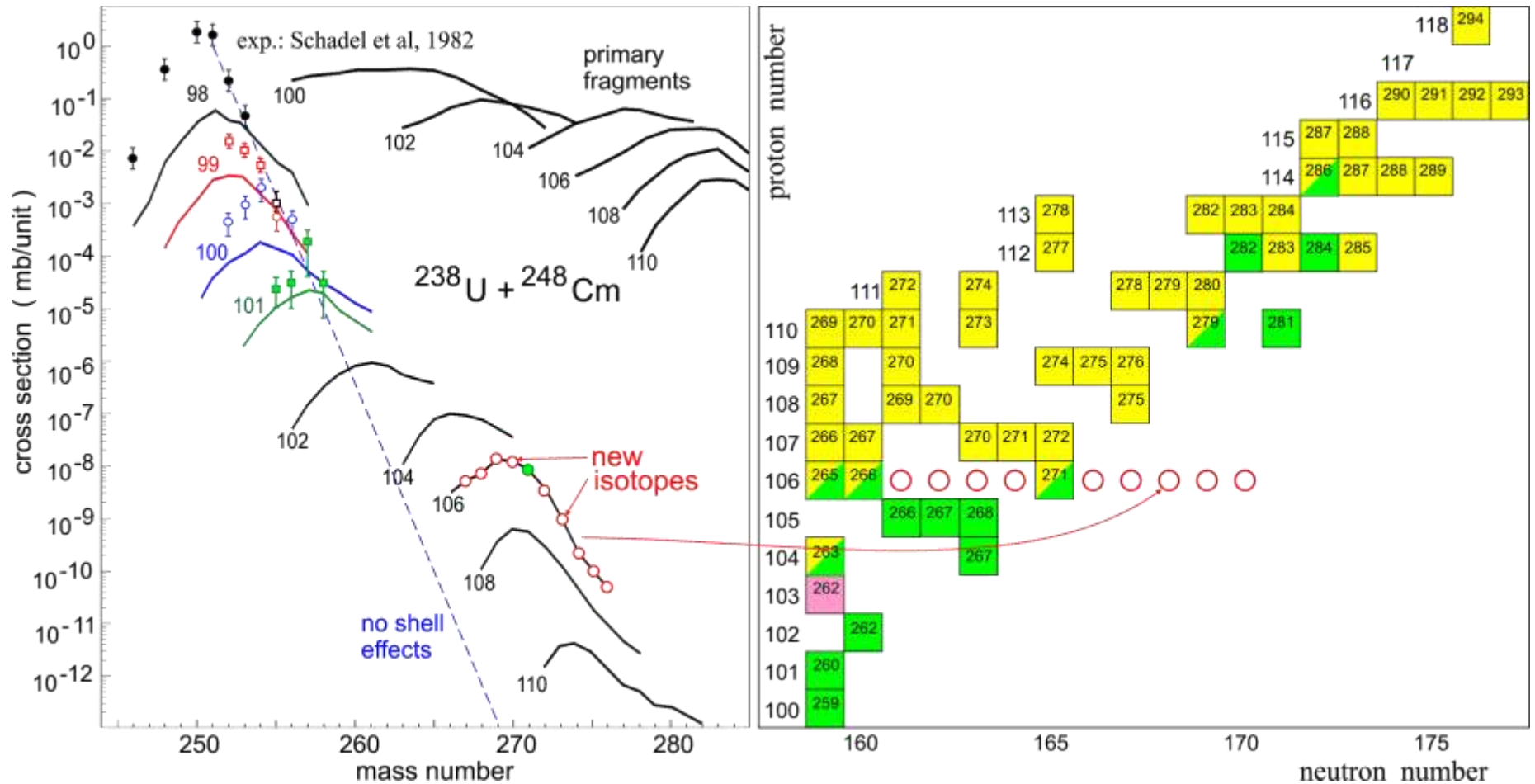
$^{238}\text{U} + ^{248}\text{Cm}$. Primary fragments



$^{238}\text{U} + ^{248}\text{Cm}$. Excitation energies and survival probability



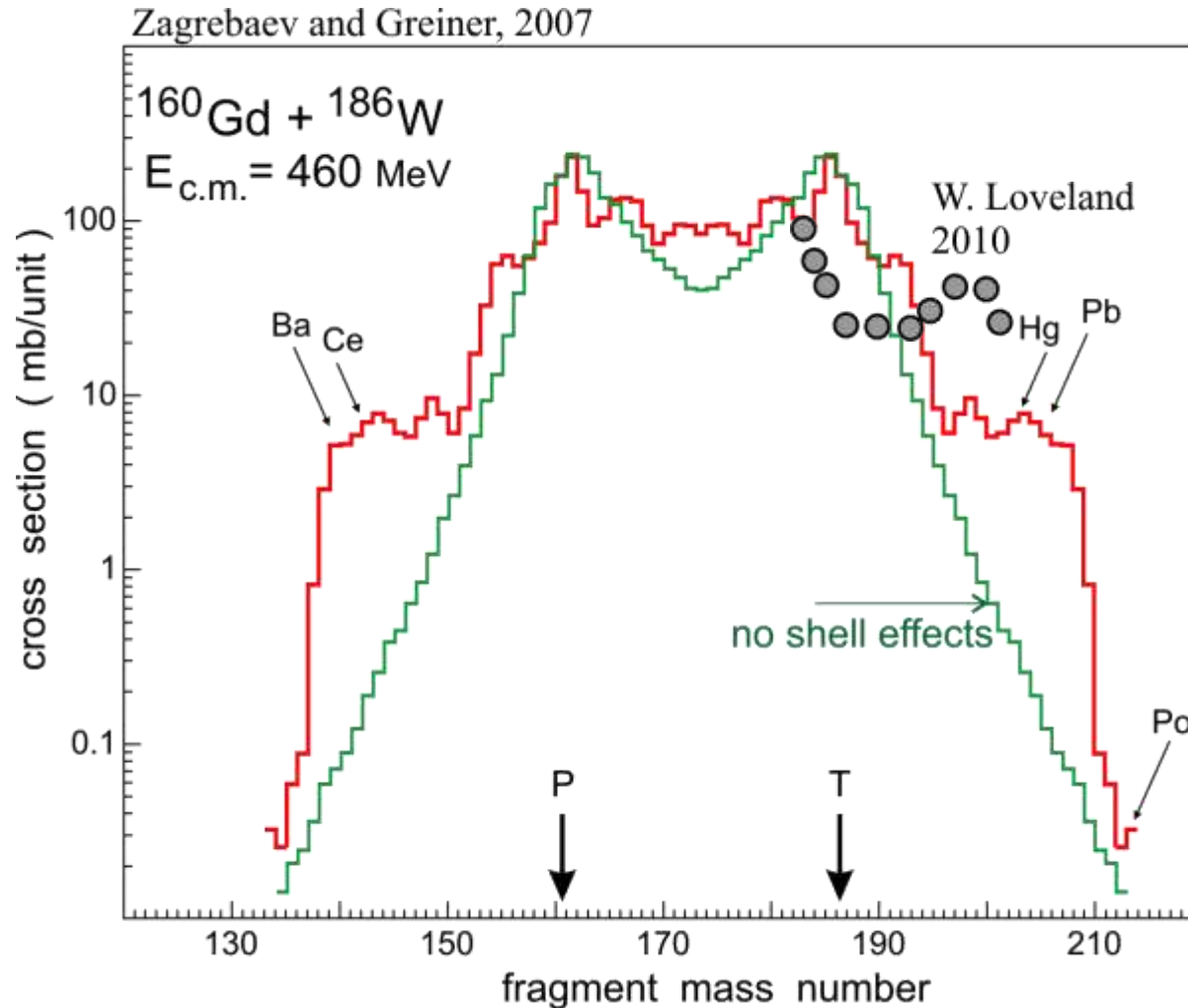
Isotopic yield of SHE in collisions of heavy actinide nuclei



How much is a role of the shell effects in damped collisions ?

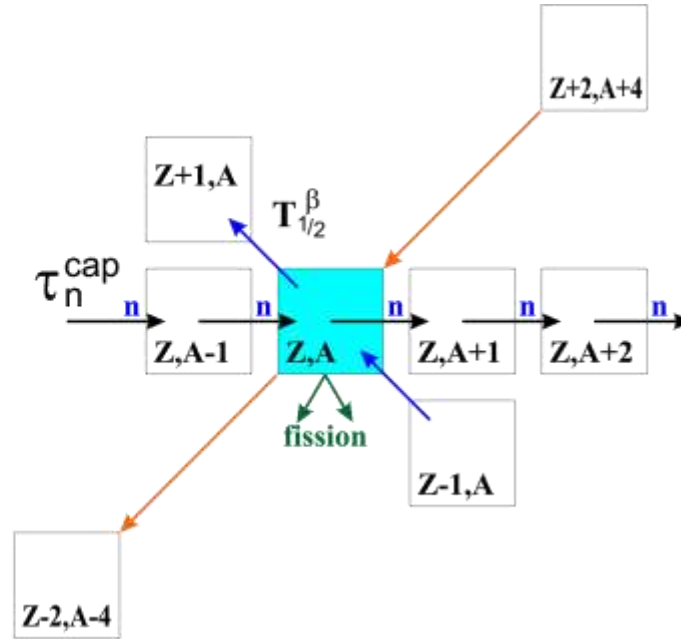


(proposal for a new experiment)



Non-accelerative production of superheavy nuclei

Nucleogenesis under the influence of neutron flux



$$\frac{dN_{ZA}}{dt} = N_{ZA-1} n_0 \sigma_{ZA-1}^{n\gamma} - N_{ZA} n_0 \sigma_{ZA}^{n\gamma} - N_{ZA} \frac{\ln 2}{T_{ZA}^\beta} - N_{ZA} \frac{\ln 2}{T_{ZA}^\alpha} - N_{ZA} \frac{\ln 2}{T_{ZA}^{\text{fis}}} + N_{Z-1A} \frac{\ln 2}{T_{Z-1A}^\beta} + N_{Z+2A+4} \frac{\ln 2}{T_{Z+2A+4}^\alpha}$$

time of neutron capture

$$\tau_n^{\text{cap}} = \frac{1}{n_0 \sigma(n, \gamma)}$$

n_0 is the neutron flux ($\frac{1}{\text{cm}^2 \cdot \text{sec}}$),
 $\sigma(n, \gamma)$ is the n-capture cross section ($\sim 1 \text{ barn} = 10^{-24} \text{ cm}^2$, $E_n = 0.5 \text{ MeV}$)

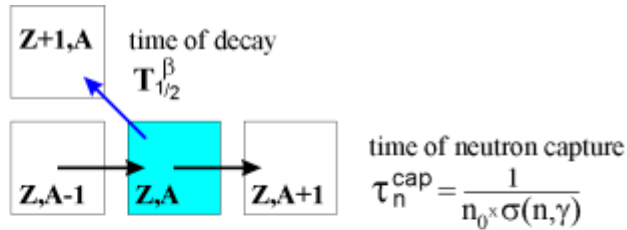
the shift to the right stops

when $T_{1/2}(Z, A) < \tau_n^{\text{cap}}$

$$n_0 (\text{reactor}) < 10^{19} \frac{1}{\text{cm}^2 \cdot \text{sec}}, \quad \tau_n^{\text{cap}} > 10^5 \text{ sec (1 day)}$$

$$n_0 (\text{explosion}) \sim 10^{30} \frac{1}{\text{cm}^2 \cdot \text{sec}}, \quad \tau_n^{\text{cap}} \sim 1 \mu\text{s}$$

Nucleogenesis in reactors and in nuclear explosion

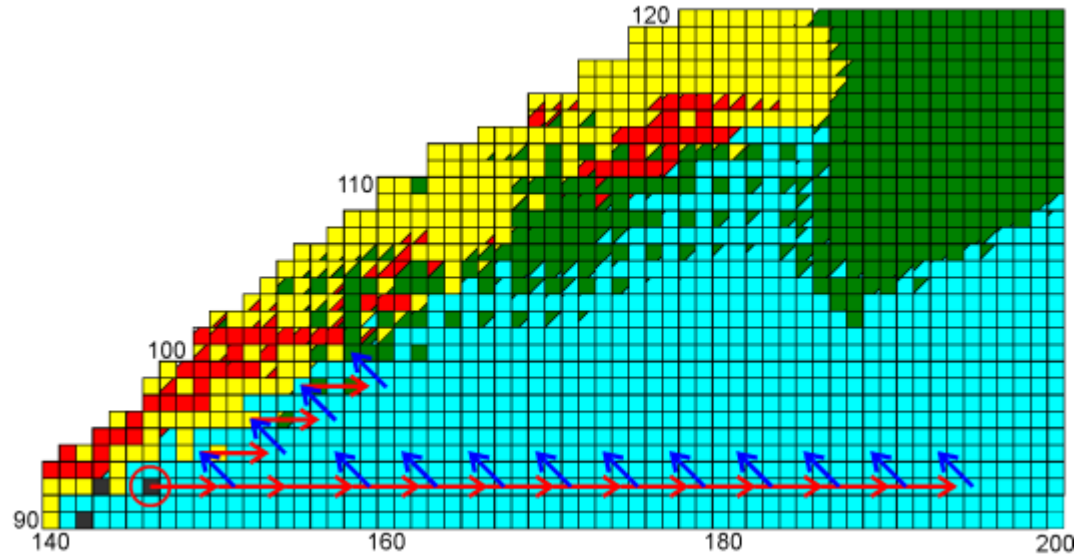
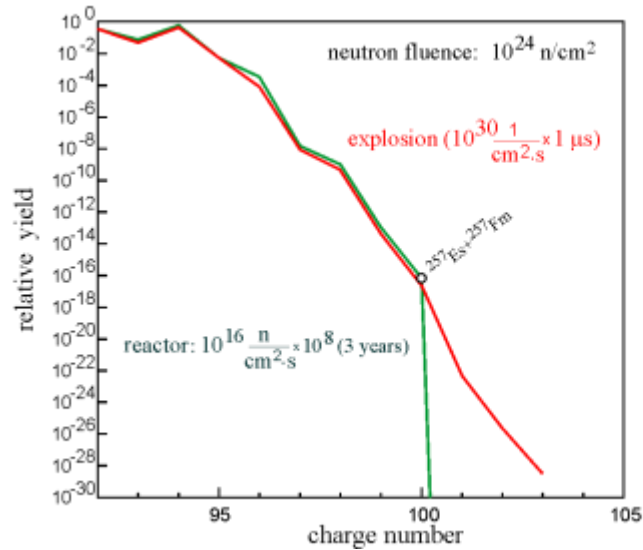
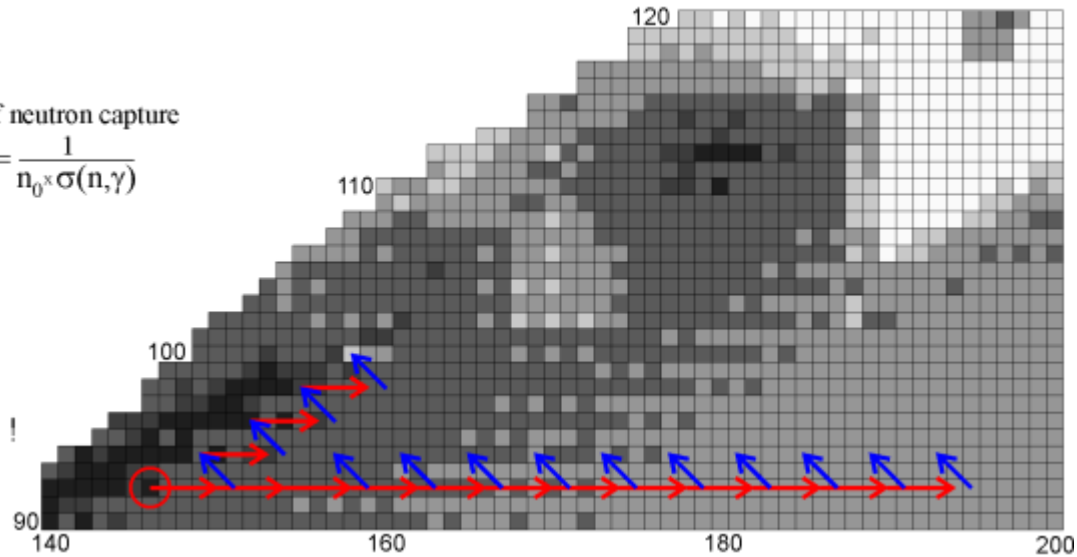


neutron capture stops

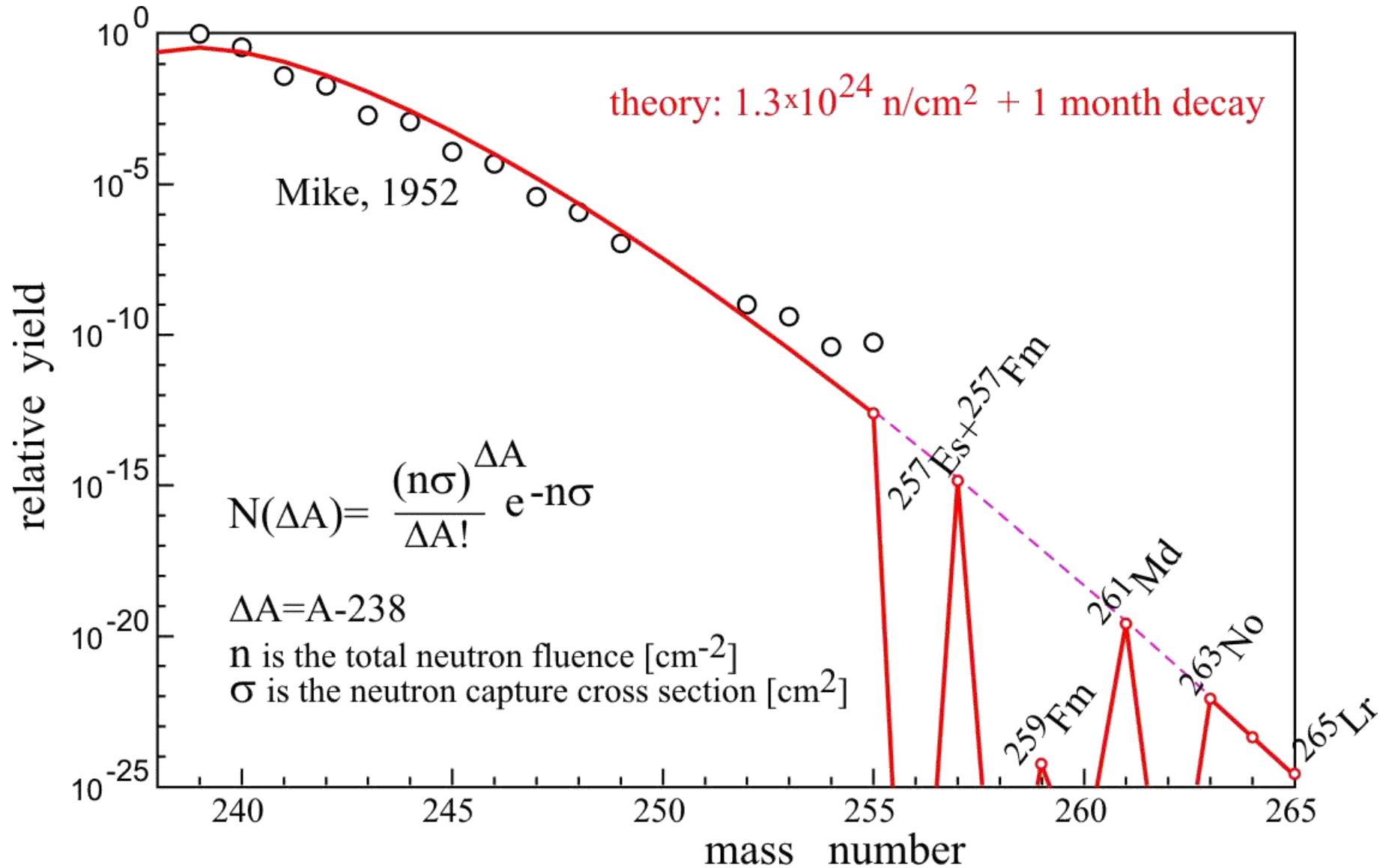
when $T_{1/2}^\beta(Z, A) < \tau_n^{cap}$

n_0 (reactor) $< 10^{19} \frac{1}{\text{cm}^2 \cdot \text{sec}}$, $\tau_n^{cap} > 10^5 \text{ sec}$ (1 day),
 but $T_{1/2}^{fis}$ (258-260 Fm) $\leq 1 \text{ sec}$!

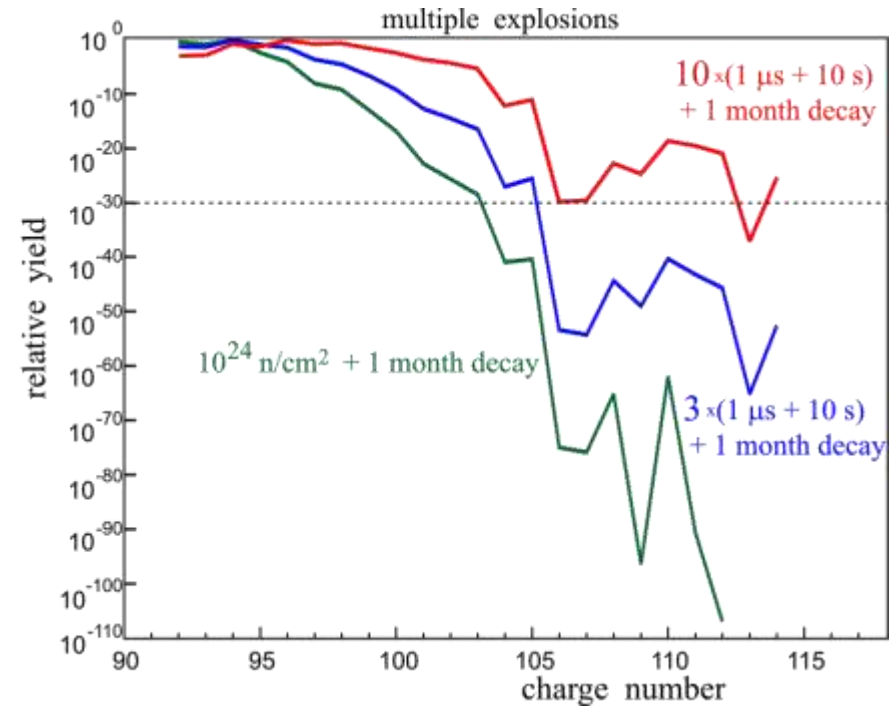
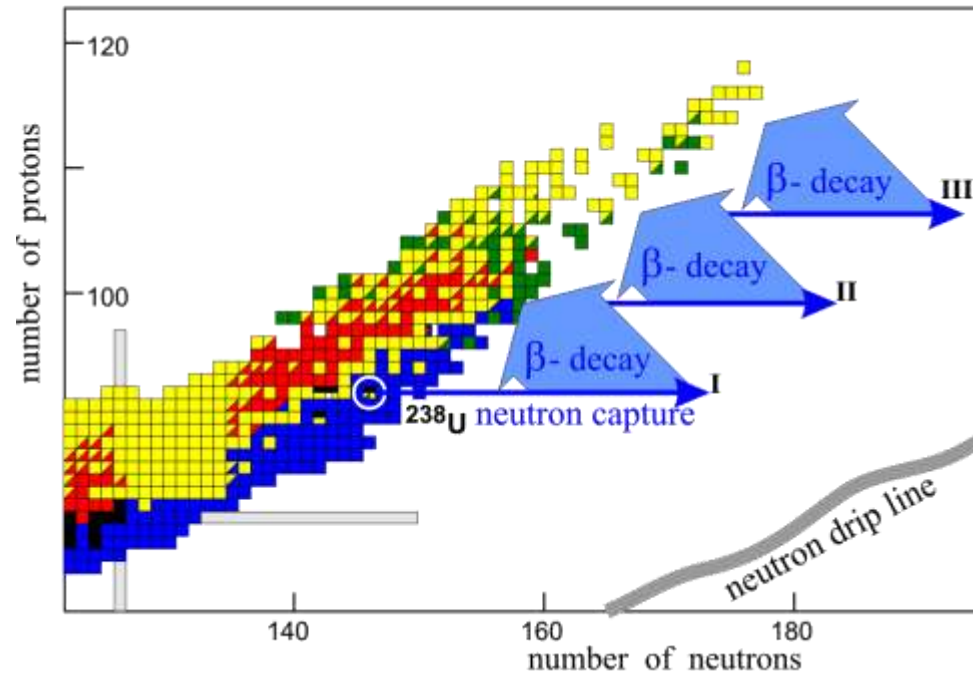
n_0 (explosion) $\sim 10^{30} \frac{1}{\text{cm}^2 \cdot \text{sec}}$, $\tau_n^{cap} \sim 1 \mu\text{s}$



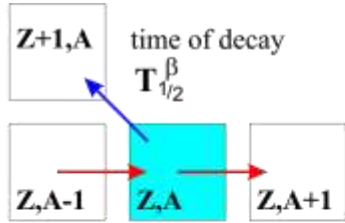
Rapid neutron capture in nuclear explosion



Multiple nuclear explosions



Pulsed reactors: Bypassing the gaps of instability



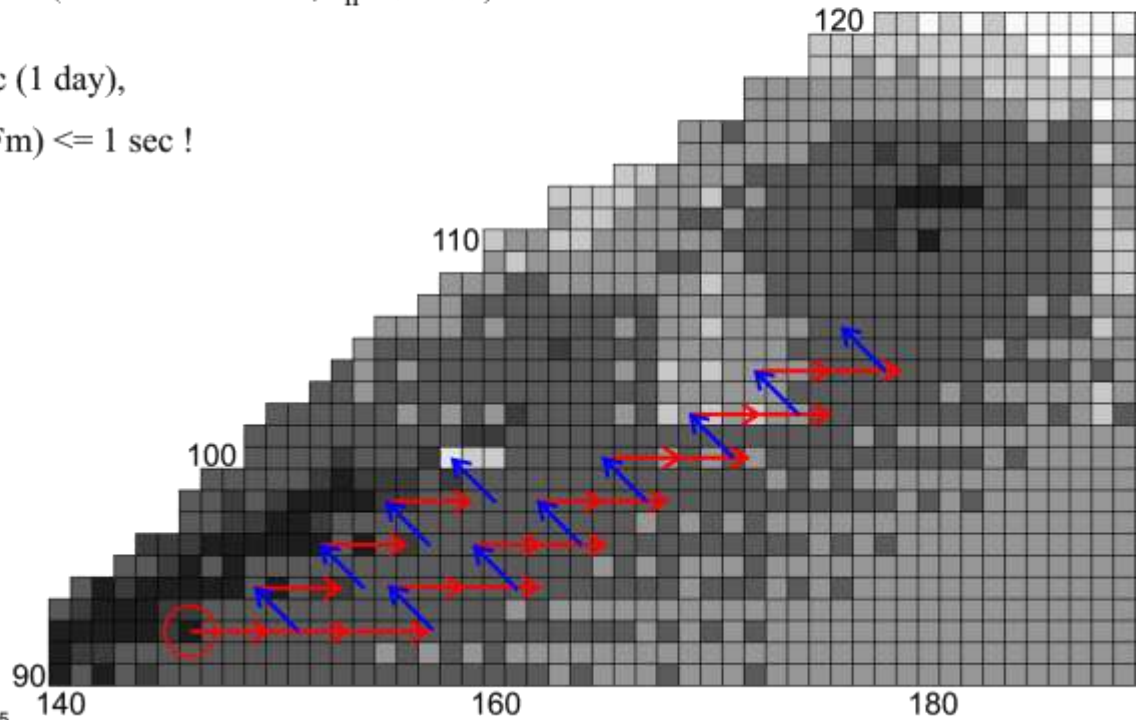
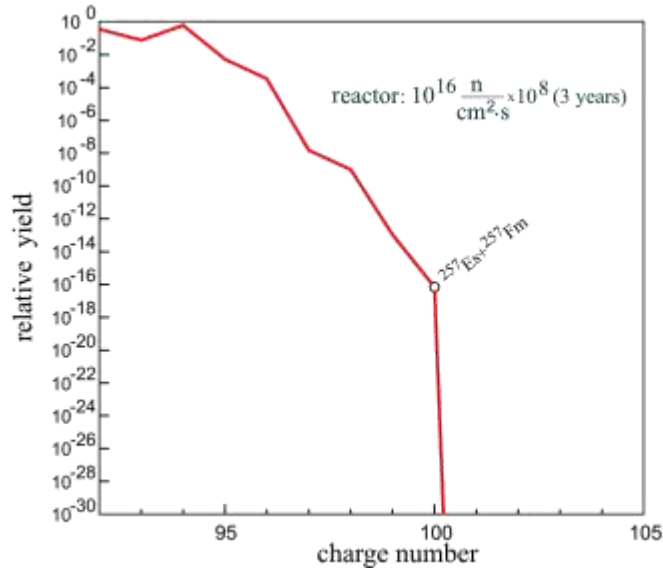
neutron capture stops
when $T_{1/2}^\beta(Z, A) < \tau_n^{\text{cap}}$

$$\tau_n^{\text{cap}} = \frac{1}{n_0 \times \sigma(n, \gamma)}$$

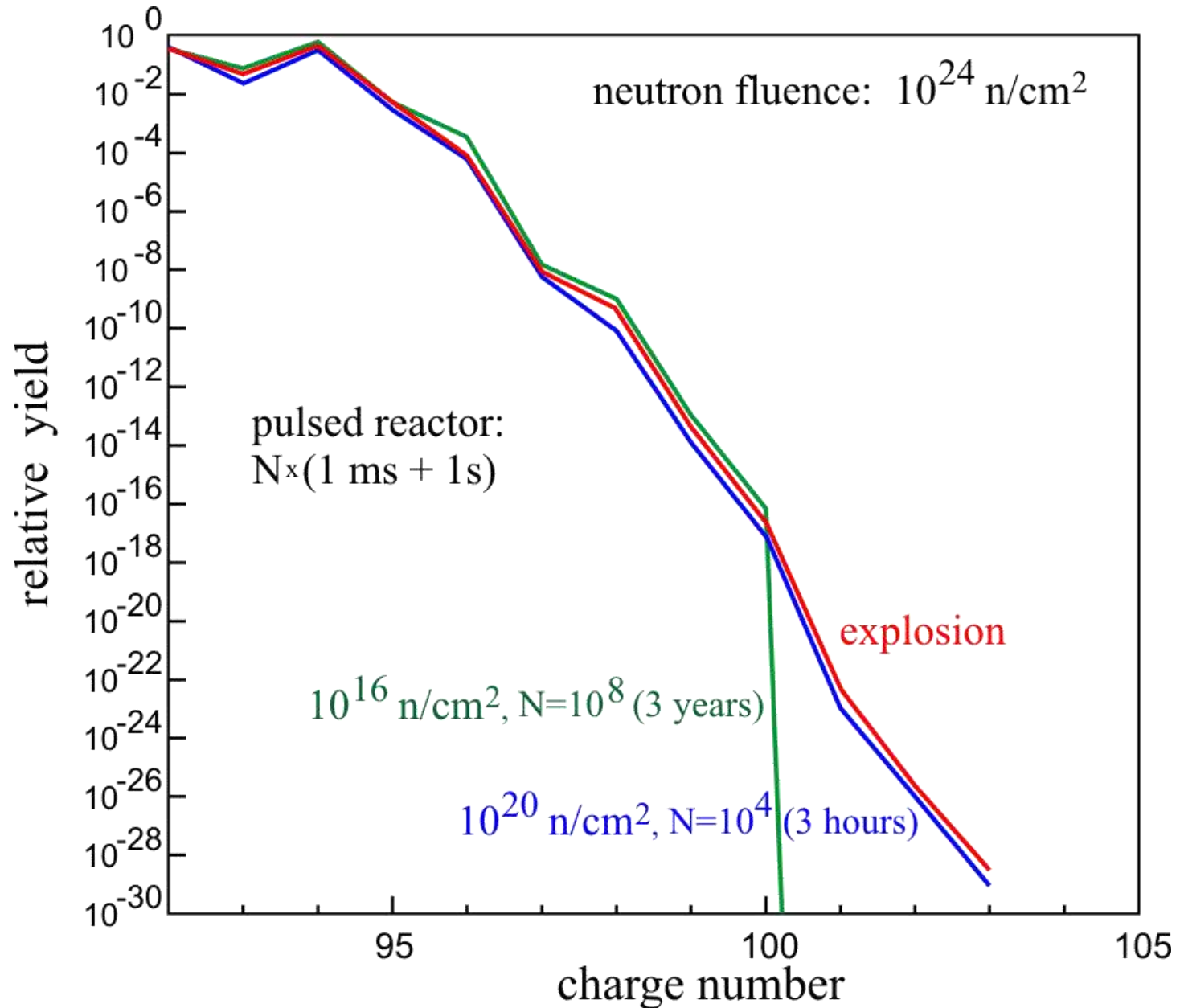
n_0 is the neutron flux ($\frac{1}{\text{cm}^2 \cdot \text{sec}}$),
 $\sigma(n, \gamma)$ is the n-capture cross section ($\sim 1 \text{ barn} = 10^{-24} \text{ cm}^2$, $E_n = 0.5 \text{ MeV}$)

$$n_0 (\text{reactor}) < 10^{19} \frac{1}{\text{cm}^2 \cdot \text{sec}}, \quad \tau_n^{\text{cap}} > 10^5 \text{ sec (1 day)},$$

but $T_{1/2}^{\text{fis}} (258-260 \text{ Fm}) \leq 1 \text{ sec} !$



New generation of Pulsed Reactors ?



Summary

- Element 120 is reachable in the Ti and/or Cr fusion reactions at the level of 0.02 pb
- Multi-nucleon transfer reactions are to be used for synthesis of neutron rich long living SH nuclei
- A macroscopic amount of the long-living SH nuclei located at the island of stability may be really produced in the multiple (rather “soft”) nuclear explosions
- This goal could be also reached by using the pulsed nuclear reactors of next generation

Problems

- Where is the island of stability ?
What is the most stable SH element to find it in Nature ?
- How much is the shell-effect enhancement in transfer reactions?
- How extensive is the Fm gap ? What is the nearest “blue” (beta-decayed) Fm isotope ?
How deep (short-living) is the gap in the region of $Z \sim 108$, $A \sim 270$?
- Is it possible to construct desired pulsed reactor to overcome the both gaps ?
Or “soft” multiple explosions are still cheaper?