New prospects in synthesis and study of superheavy nuclei

- Superheavy dreams of Walter Greiner
- Epoch of "cold" and "hot" fusion reactions (100 < Z <= 118)
- Multi-nucleon transfer reactions
- Non-accelerative SHE production
- Summary

Valery Zagrebaev and Walter Greiner



for XIII WONP – VII NURT, February 08, 2011, Havana, Cuba



FIAS (Frankfurt)

Mendeleev's Table 140 years ago

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Ancient Nuclear Map and dreams of Walter Greiner



Mendeleev's Table

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Synthesis of new elements (history)



Synthesis of superheavy elements (experiment)



Synthesis of superheavy elements (experiment)



Target made of ²⁴⁹Cf (half-life is 350 years)





Detectors for SH nucleus recoil and products of its decay



Significance of Theory

Which nuclei are to be fused ?

«cold» synthesis: ²⁰⁸Pb + ⁶⁴Ni, ⁷⁰Zn, ... → ²⁷²110, ²⁷⁸112, ... (GSI, Germany)
«hot» synthesis: ²³⁸U, ²⁴⁴Pu, ²⁴⁸Cm, ²⁴⁹Cf + ⁴⁸Ca → ²⁸⁶112, ²⁹²114, ²⁹⁶116, ²⁹⁷118 (Dubna)
Symmetric combinations: ¹⁴⁸Nd + ¹⁵⁴Sm → ³⁰²122 ? Radioactive beams of ¹³²Sn, etc. ?



Superheavy Elements (Island of Stability)



Great progress in synthesis of superheavy nuclei



Beyond ⁴⁸Ca: Pursuit of 120



Radioactive Ion Beams for production of neutron rich superheavy nuclei ?



Multi-nucleon transfer reactions in low-energy heavy ion collisions

Simulation of experiment and cross sections



Dynamics: 10^{6} tested events (trajectories), Statistical model: 10^{-6} (3n), 10^{-7} (4n) survival probability cross sections up to 0.1 pb can be calculated

Time-dependent Driving Potential

 $V_{\text{diabat}}(R,\beta_1,\beta_2,\alpha,...) = V_{12}^{\text{folding}}(Z_1,N_1,Z_2,N_2;R,\beta_1,\beta_2,...) + M(A_1) + M(A_2) - M(\text{Proj}) - M(\text{Targ})$



 $V_{\text{adiabat}}(\mathsf{R},\beta_1,\beta_2,\eta,...) = \mathsf{M}_{\mathsf{TCSM}}(\mathsf{R},\beta_1,\beta_2,\eta,...) - \mathsf{M}(\mathsf{Proj}) - \mathsf{M}(\mathsf{Targ})$

Time -dependent driving potential has to be used $V(t) = V_{\text{diab}}(\xi) \cdot \exp(-\frac{t_{\text{int}}}{\tau_{\text{relax}}}) + V_{\text{adiab}}(\xi) \cdot [1 - \exp(-\frac{t_{\text{int}}}{\tau_{\text{relax}}})]$ $\tau_{\text{relax}} \sim 10^{-21} \text{ s}$ the same degrees of freedom ($\xi = R, \theta, \phi_1, \phi_2, \beta_1, \beta_2, \eta_Z, \eta_N$) ! All forces, $F_i(t) = -\frac{\partial V}{\partial \xi_i}$, are quite smooth

Time-dependent Driving Potential





Nucleon Exchange

(L. Moretto, 1974) (L. Moretto, 1974) Distribution function $\varphi(A_1, t) \rightarrow \text{Master equation} \quad \frac{\partial \varphi}{\partial t} = \sum_{A_1'=A_1\pm 1} \lambda(A_1' \rightarrow A_1) \cdot \varphi(A_1') - \lambda(A_1 \rightarrow A_1') \cdot \varphi(A_1)$ $\frac{\partial \varphi}{\partial t} = -\frac{\partial}{\partial A_1} \left(D^{(1)} \varphi \right) + \frac{\partial^2}{\partial A_1^2} \left(D^{(2)} \varphi \right) \quad \text{Fokker - Planck}_{(W. \text{ Nörenberg, 1974})}$ $\eta = \frac{A_1 - A_2}{A_{CN}} = \frac{A_1 - (A_{CN} - A_1)}{A_{CN}} = \frac{2A_1 - A_{CN}}{A_{CN}}$ $\frac{dA_1}{dt} = D^{(1)} + \sqrt{D^{(2)}} \Gamma(t) \quad \text{Langevin type eq.}$ $\frac{d\eta}{dt} = \frac{2}{A_{\text{ev}}} D_A^{(1)} + \frac{2}{A_{\text{ev}}} \sqrt{D_A^{(2)}} \Gamma(t)$ at A' = A ± 1 $D^{(1)} = \lambda(A_1 \rightarrow A_1 + 1) - \lambda(A_1 \rightarrow A_1 - 1)$ $D^{(2)} = \frac{1}{2} [\lambda(A_1 \rightarrow A_1 + 1) + \lambda(A_1 \rightarrow A_1 - 1)]$ $\lambda^{(\pm)} = \lambda_0 \sqrt{\frac{\rho(A\pm 1)}{\rho(A)}} P_{\text{tr}}(R; A \to A\pm 1), \quad \rho \sim exp(2\sqrt{aE^*}), \quad E^* = E_{\text{c.m.}} - V(R, \beta_1, \beta_2, \eta)$ transition probability

$$\begin{array}{c} \sum_{\substack{A_1 \\ A_2 \\ N_2 \rightarrow N_2 + 1 \\ N_2 \rightarrow N_2 + 1 \end{array}} \sum_{\substack{R_1 \rightarrow R_2 + 1 \\ N_2 \rightarrow N_2 + 1 \end{array}} \eta_{R} = \frac{Z_1 - Z_2}{Z_1 + Z_2} & D_{N,Z}^{(1)} = \lambda_{N,Z} (A \rightarrow A + 1) - \lambda_{N,Z} (A \rightarrow A - 1) \\ D_{N,Z}^{(2)} = \frac{1}{2} [\lambda_{N,Z} (A \rightarrow A + 1) + \lambda_{N,Z} (A \rightarrow A - 1)] \\ \lambda_{N,Z}^{(\pm)} = \lambda_{N,Z}^{0} \sqrt{\frac{\rho(A \pm 1)}{\rho(A)}} P_{tr}(R; A \rightarrow A \pm 1) \end{array}$$

 $\frac{dR}{dR} = \frac{p_R}{p_R}$ Variables: {R, θ , φ_1 , φ_2 , β_1 , β_2 , η_7 , η_N } $\frac{\frac{d\theta}{d\theta}}{\frac{d\theta}{dt}} = \frac{\frac{\mu_R}{\mu_R}}{\frac{\ell}{\mu_R R^2}}$ Most uncertain parameters: μ_0, γ_0 - nuclear viscosity and friction, λ_Z^0 , λ_N^0 - nucleon transfer rate $\frac{d\varphi_1}{dt} = \frac{L_1}{\mathfrak{I}_1}, \ \frac{d\varphi_2}{dt} = \frac{L_2}{\mathfrak{I}_2}$ $\eta = \frac{A_{1} - A_{2}}{A_{1} + A_{2}}$ $\eta_{Z} = \frac{Z_{1} - Z_{2}}{Z_{1} + Z_{2}}$ φ1 $\frac{d\beta_1}{dt} = \frac{p_{\beta 1}}{\mu_{\beta 1}}$ R A₁ μ_{B1} b θ. $\frac{d\beta_2}{dt} = \frac{p_{\beta 2}}{\mu_{\beta 2}}$ $\eta_{N} = \frac{N_{1} - N_{2}}{N_{1} + N_{2}}$ $\langle \varphi_2 \rangle$ Α2 $\frac{d\eta_{z}}{dt} = \frac{2}{Z_{\rm CN}} D_{\rm Z}^{(1)} + \frac{2}{Z_{\rm CN}} \sqrt{D_{\rm Z}^{(2)}} \Gamma_{\rm Z} (t)$ $\lambda_{\mathbf{Z}}^{\mathbf{0}} = \lambda_{\mathbf{N}}^{\mathbf{0}} = \frac{\lambda_{\mathbf{Q}}^{\mathbf{0}}}{2}$ $\frac{d\eta_{\rm N}}{dt} = \frac{2}{N_{\rm CN}} D_{\rm N}^{(1)} + \frac{2}{N_{\rm CN}} \sqrt{D_{\rm N}^{(2)}} \Gamma_{\rm N} (t)$ $\frac{dp_R}{dt} = -\frac{\partial V}{\partial R} + \frac{\ell^2}{\mu_R R^3} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2}\right)\frac{\partial \mu_R}{\partial R} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2}\frac{\partial \mu_{\beta_1}}{\partial R} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2}\frac{\partial \mu_{\beta_2}}{\partial R} - \gamma_R \frac{p_R}{\mu_R} + \sqrt{\gamma_R T}\Gamma_R(t)$ $\frac{d\ell}{dt} = -\frac{\partial V}{\partial \vartheta} - \gamma_{\text{tang}} \left(\frac{\ell}{\mu_{R}R} - \frac{L_{1}}{\Im_{1}}a_{1} - \frac{L_{2}}{\Im_{2}}a_{2} \right) R + \sqrt{\gamma_{\text{tang}}T}\Gamma_{\text{tang}}(t)$ $\frac{dL_1}{dt} = -\frac{\partial V}{\partial \varphi_1} + \gamma_{\text{tang}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\Im_1} a_1 - \frac{L_2}{\Im_2} a_2 \right) a_1 - \frac{a_1}{R} \sqrt{\gamma_{\text{tang}} T} \Gamma_{\text{tang}}(t)$ $\frac{dL_2}{dt} = -\frac{\partial V}{\partial \varphi_2} + \gamma_{\text{tan}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\Im_1} a_1 - \frac{L_2}{\Im_2} a_2 \right) a_2 - \frac{a_2}{R} \sqrt{\gamma_{\text{tang}} T} \Gamma_{\text{tang}}(t)$ $\frac{dp_{\beta_1}}{dt} = -\frac{\partial V}{\partial \beta_1} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_1} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_1} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2}\right) \frac{\partial \mu_R}{\partial \beta_1} - \gamma_\beta \frac{p_{\beta_1}}{\mu_{\beta_1}} + \sqrt{\gamma_{\beta_1} T} \Gamma_{\beta_1}(t)$ $\frac{dp_{\beta_2}}{dt} = -\frac{\partial V}{\partial \beta_2} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_2} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_2} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2}\right) \frac{\partial \mu_R}{\partial \beta_2} - \gamma_\beta \frac{p_{\beta_2}}{\mu_{\beta_2}} + \sqrt{\gamma_{\beta_2} T} \Gamma_{\beta_2}(t)$

Good agreement with experiment: e.g. ⁴⁸Ca + ²⁴⁸Cm



Transfer reactions in damped collision of very heavy nuclei ?



Most probable way of evolution of the giant nuclear system



238U + 248Cm. Primary fragments



238U + 248Cm. Excitation energies and survival probability



Isotopic yield of SHE in collisions of heavy actinide nuclei



How much is a role of the shell effects in damped collisions ? ¹⁶⁰Gd + ¹⁸⁶W

(proposal for a new experiment)



Non-accelerative production of superheavy nuclei

Nucleogenesis under the influence of neutron flux



time of neutron capture

$$\tau_n^{cap} = \frac{1}{n_0^x \sigma(n,\gamma)} \qquad \begin{array}{l} n_0 \text{ is the neutron flux } (\frac{1}{cm^2 \cdot sec}), \\ \sigma(n,\gamma) \text{ is the n-capture cross section } (\sim 1 \text{ barn} = 10^{-24} \text{ cm}^2, \text{ E}_n = 0.5 \text{ MeV}) \end{array}$$

the shift to the right stops

when
$$\mathbf{T}_{1/2}(\mathbf{Z}, \mathbf{A}) < \tau_n^{cap}$$

 $n_0 (reactor) < 10^{19} \frac{1}{cm^2 \cdot sec}$, $\tau_n^{cap} > 10^5 \sec (1 \text{ day})$
 $n_0 (explosion) \sim 10^{30} \frac{1}{cm^2 \cdot sec}$, $\tau_n^{cap} \sim 1 \,\mu s$

Nucleogenesis in reactors and in nuclear explosion



Rapid neutron capture in nuclear explosion



Multiple nuclear explosions



Pulsed reactors: Bypassing the gaps of instability



New generation of Pulsed Reactors ?



Summary

- Element 120 is reachable in the Ti and/or Cr fusion reactions at the level of 0.02 pb
- Multi-nucleon transfer reactions are to be used for synthesis of neutron rich long living SH nuclei
- A macroscopic amount of the long-living SH nuclei located at the island of stability may be really produced in the multiple (rather "soft") nuclear explosions
- This goal could be also reached by using the pulsed nuclear reactors of next generation

Problems

- Where is the island of stability ? What is the most stable SH element to find it in Nature ?
- How much is the shell-effect enhancement in transfer reactions?
- How extensive is the Fm gap ? What is the nearest "blue" (beta-decayed) Fm isotope ? How deep (short-living) is the gap in the region of Z~108, A~270 ?
- Is it possible to construct desired pulsed reactor to overcome the both gaps ? Or "soft" multiple explosions are still cheaper?