

Production of heavy neutron-rich nuclei by multi-nucleon transfer reactions

- **State of the art: neutron rich heavy nuclei were not synthesized yet**
- **Outline of the model (3 slides only)**
- **Low-energy multi-nucleon transfer reactions:**
 - **Production of trans-target nuclei**
 - **Shell effects in damped collisions of heavy ions ?**
 - **Production of neutron rich nuclei located along the neutron closed shell $N=126$**
 - **Synthesis of neutron rich transfermium nuclei**
- **New setup for selective laser ionization of heavy neutron rich nuclei**
- **Summary**

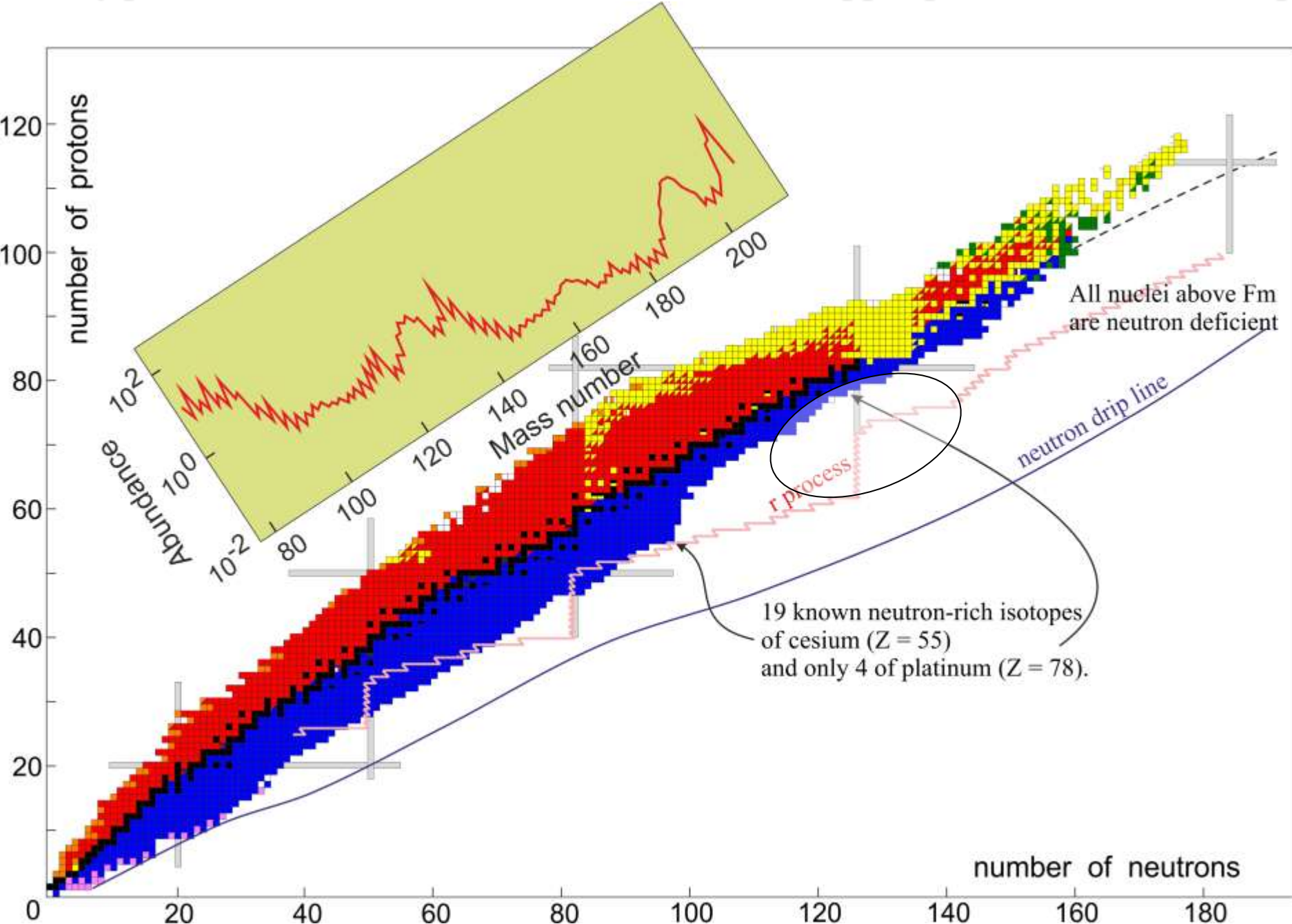
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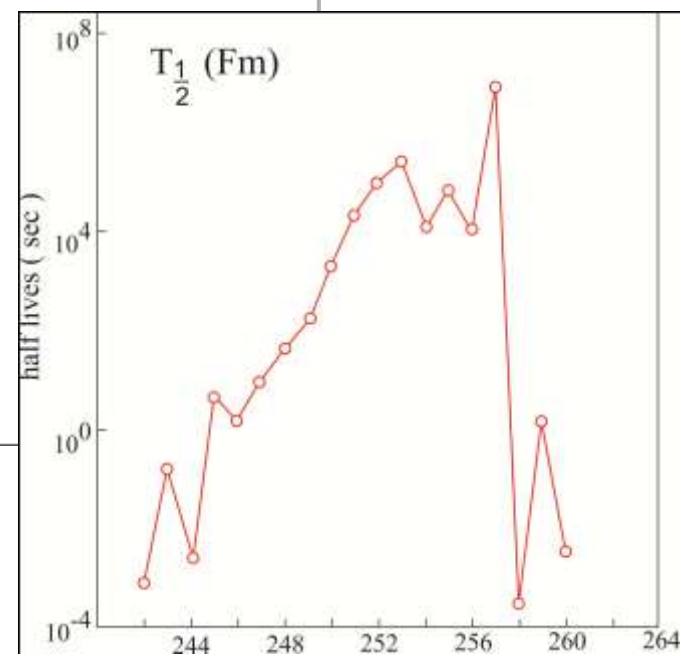
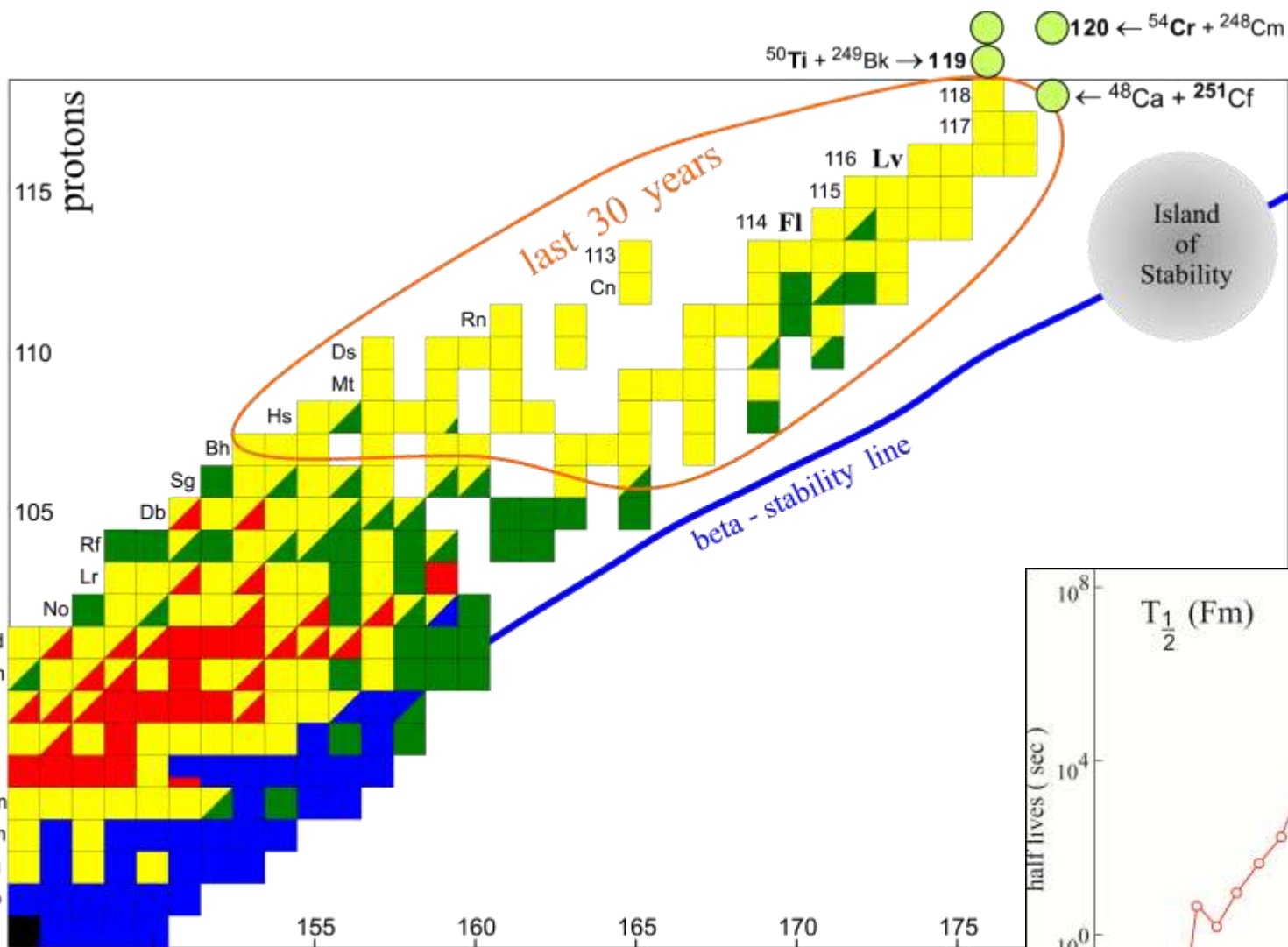
for “Colloque GANIL 2013”, September 25, 2013, Port-en-Bessin, France



Mostly proton-rich nuclei were studied so far in the upper part of the nuclear map

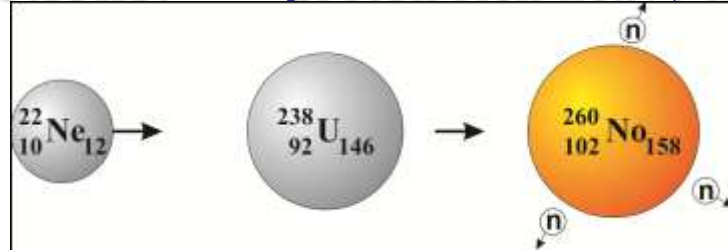


SHE: we are far from the stability line and from the Island of Stability

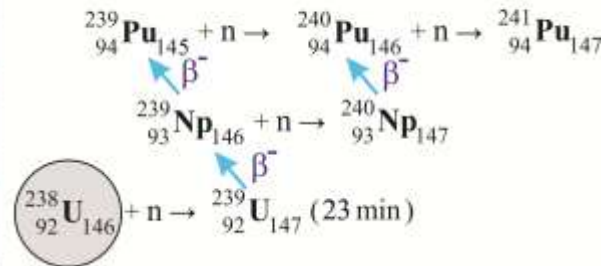


There are only 3 methods for synthesis of heavy nuclei

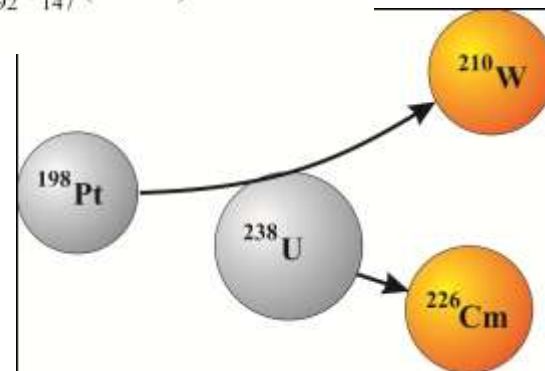
1. Fusion reactions: → proton rich heavy nuclei



2. Sequence of neutron capture and beta(-) decay processes: neutron fluxes in reactors are too low, nuclear explosions are forbidden



3. Multi-nucleon transfer reactions



Theoretical models of transfer reactions

Multi-nucleon transfers in damped collisions

Master equation

L.G. Moretto and J.S. Sventek, Phys. Lett. B **58**, 26 (1975)

Fokker-Plank equation

W. Norenberg, Phys. Lett. B **52**, 289 (1974)

Langevin equations

P. Frobrich and S.Y. Xu, Nucl. Phys. **A477**, 143 (1988)

Semi-classical approaches

E. Vigezzi and A. Winther, Ann. Phys. (N.Y.) **192**, 432 (1989).

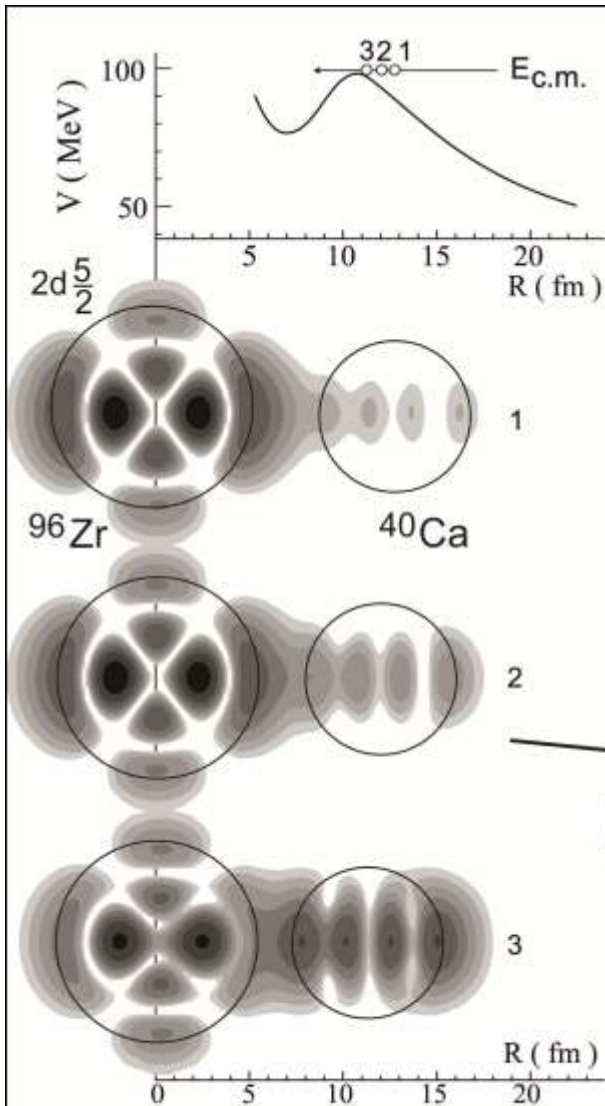
V.I. Zagrebaev, Ann. Phys. (N.Y.) **197**, 33 (1990).

Few-nucleon transfers (GRAZING)

A. Winther, Nucl. Phys. **A594**, 203 (1995)

<http://personalpages.to.infn.it/nanni/grazing>

Adiabatic dynamics of low-energy heavy ion collisions and nucleon transfers



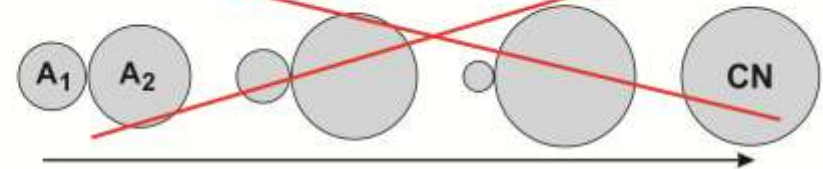
Appropriate description:

time-dependent Schrödinger equation for single particle wave functions
 (VZ, Samarin, Greiner, 2007);

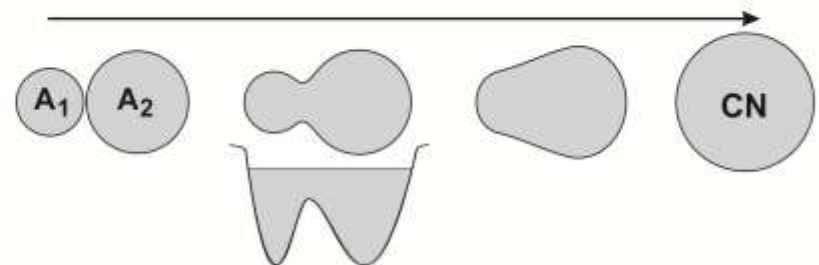
~~Overlapping of single particle wave functions:~~

$$\langle \psi_{A_1}^i(\mathbf{R}+\mathbf{r}_1) | \Delta V | \psi_{A_2}^k(\mathbf{R}+\mathbf{r}_k) \rangle$$

~~DNS. frozen nuclei, isolated mean fields~~



- overlapped mean fields
- two-center shell model
- rearrangement of nucleons



System of coupled Langevin type Equations of Motion

$$\frac{dR}{dt} = \frac{p_R}{\mu_R} \quad \text{Variables: } \{R, \theta, \varphi_1, \varphi_2, \beta_1, \beta_2, \eta_Z, \eta_N\}$$

$$\frac{d\vartheta}{dt} = \frac{\ell}{\mu_R R^2}$$

$$\frac{d\varphi_1}{dt} = \frac{L_1}{\mathfrak{I}_1}, \quad \frac{d\varphi_2}{dt} = \frac{L_2}{\mathfrak{I}_2}$$

$$\frac{d\beta_1}{dt} = \frac{p_{\beta_1}}{\mu_{\beta_1}}$$

$$\frac{d\beta_2}{dt} = \frac{p_{\beta_2}}{\mu_{\beta_2}}$$

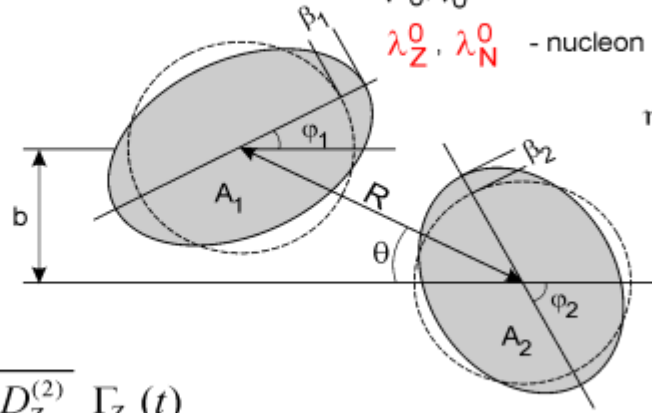
$$\frac{d\eta_Z}{dt} = \frac{2}{Z_{CN}} D_Z^{(1)} + \frac{2}{Z_{CN}} \sqrt{D_Z^{(2)}} \Gamma_Z(t)$$

$$\frac{d\eta_N}{dt} = \frac{2}{N_{CN}} D_N^{(1)} + \frac{2}{N_{CN}} \sqrt{D_N^{(2)}} \Gamma_N(t)$$

Most uncertain parameters:

μ_0, γ_0 - nuclear viscosity and friction,

λ_Z^0, λ_N^0 - nucleon transfer rate



$$\eta = \frac{A_1 - A_2}{A_1 + A_2}$$

$$\eta_Z = \frac{Z_1 - Z_2}{Z_1 + Z_2}$$

$$\eta_N = \frac{N_1 - N_2}{N_1 + N_2}$$

$$\lambda_Z^0 = \lambda_N^0 = \frac{\lambda^0}{2}$$

$$\frac{dp_R}{dt} = -\frac{\partial V}{\partial R} + \frac{\ell^2}{\mu_R R^3} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial R} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial R} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial R} - \gamma_R \frac{p_R}{\mu_R} + \sqrt{\gamma_R} T \Gamma_R(t)$$

$$\frac{d\ell}{dt} = -\frac{\partial V}{\partial \vartheta} - \gamma_{\text{tang}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) R + \sqrt{\gamma_{\text{tang}}} T \Gamma_{\text{tang}}(t)$$

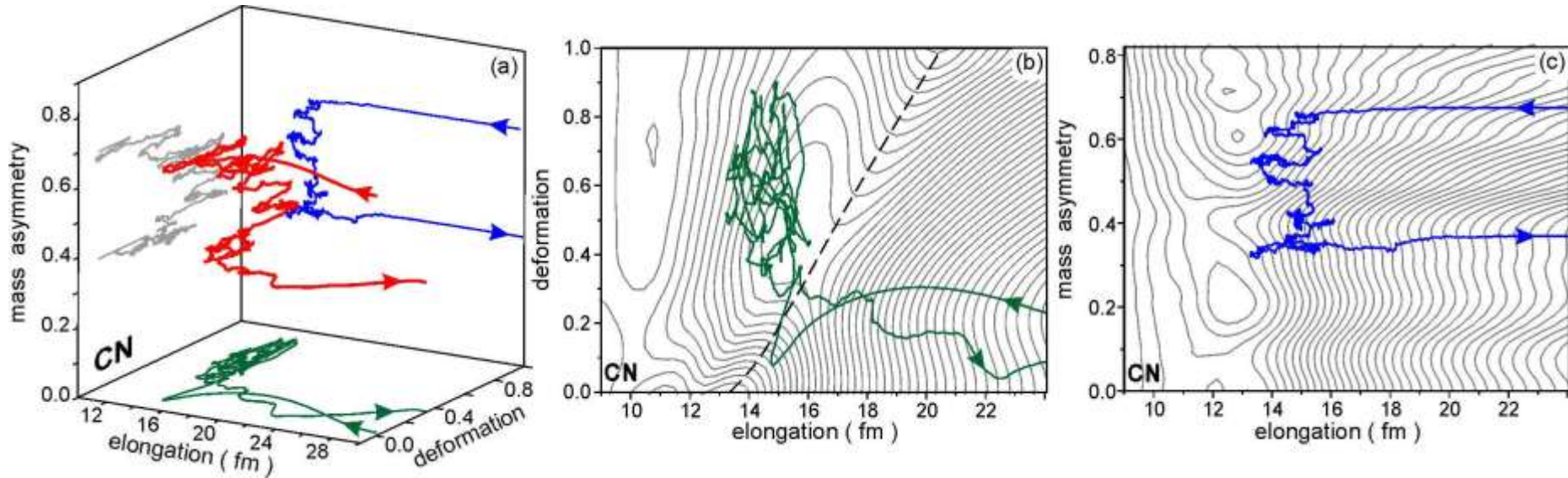
$$\frac{dL_1}{dt} = -\frac{\partial V}{\partial \varphi_1} + \gamma_{\text{tang}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) a_1 - \frac{a_1}{R} \sqrt{\gamma_{\text{tang}}} T \Gamma_{\text{tang}}(t)$$

$$\frac{dL_2}{dt} = -\frac{\partial V}{\partial \varphi_2} + \gamma_{\text{tan}} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{\mathfrak{I}_1} a_1 - \frac{L_2}{\mathfrak{I}_2} a_2 \right) a_2 - \frac{a_2}{R} \sqrt{\gamma_{\text{tang}}} T \Gamma_{\text{tang}}(t)$$

$$\frac{dp_{\beta_1}}{dt} = -\frac{\partial V}{\partial \beta_1} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_1} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_1} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial \beta_1} - \gamma_{\beta} \frac{p_{\beta_1}}{\mu_{\beta_1}} + \sqrt{\gamma_{\beta_1}} T \Gamma_{\beta_1}(t)$$

$$\frac{dp_{\beta_2}}{dt} = -\frac{\partial V}{\partial \beta_2} + \frac{p_{\beta_1}^2}{2\mu_{\beta_1}^2} \frac{\partial \mu_{\beta_1}}{\partial \beta_2} + \frac{p_{\beta_2}^2}{2\mu_{\beta_2}^2} \frac{\partial \mu_{\beta_2}}{\partial \beta_2} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{p_R^2}{2\mu_R^2} \right) \frac{\partial \mu_R}{\partial \beta_2} - \gamma_{\beta} \frac{p_{\beta_2}}{\mu_{\beta_2}} + \sqrt{\gamma_{\beta_2}} T \Gamma_{\beta_2}(t)$$

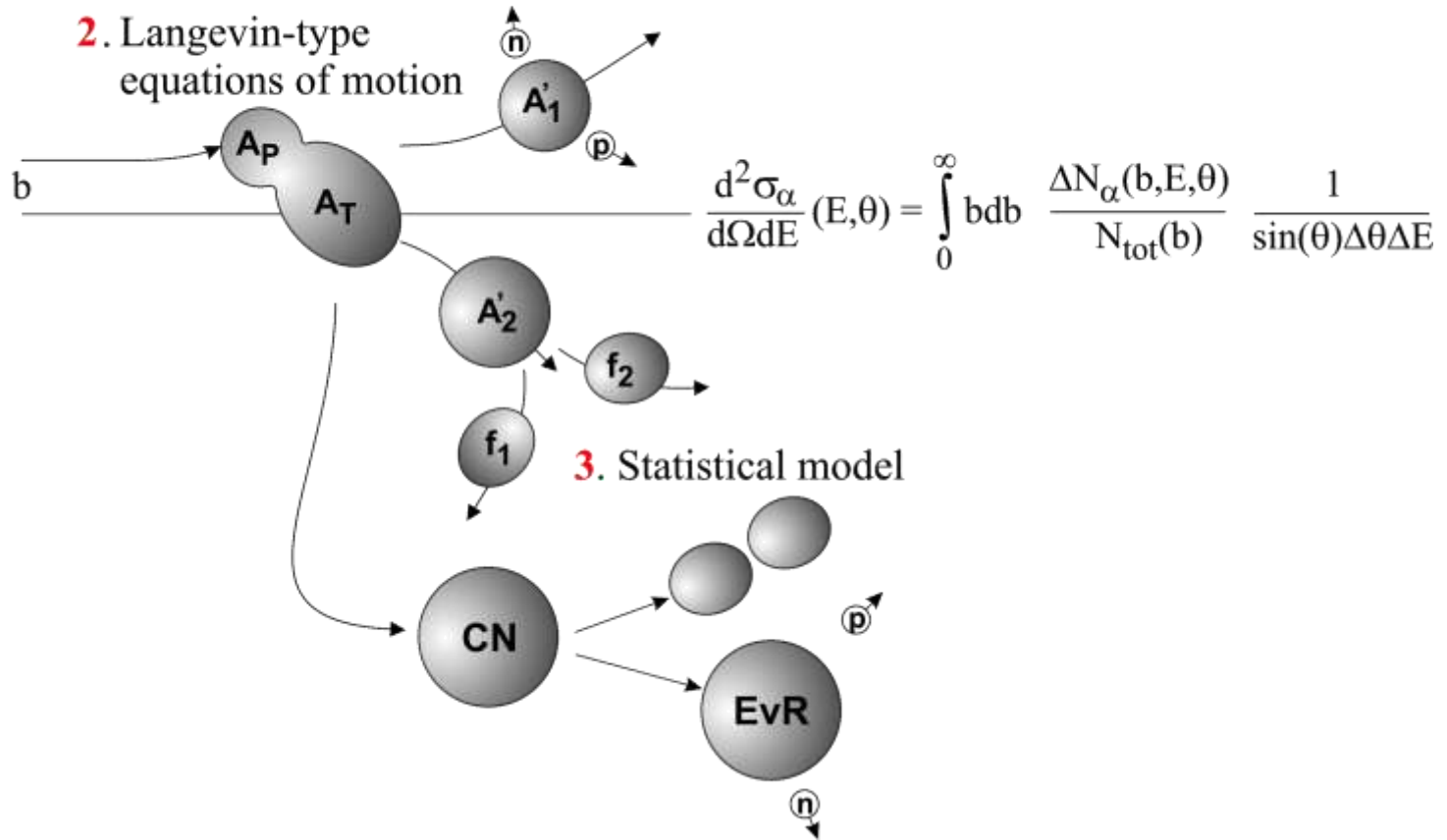
Typical trajectory in the “distance-deformation-mass asymmetry” space ($48\text{Ca} + 248\text{Cm}$, $E=210$ MeV)



Simulation of experiment and cross sections

1. Time-dependent driving potential $V(r, \xi; t)$:
Folding \rightarrow Adiabatic Two-Center Shell Model

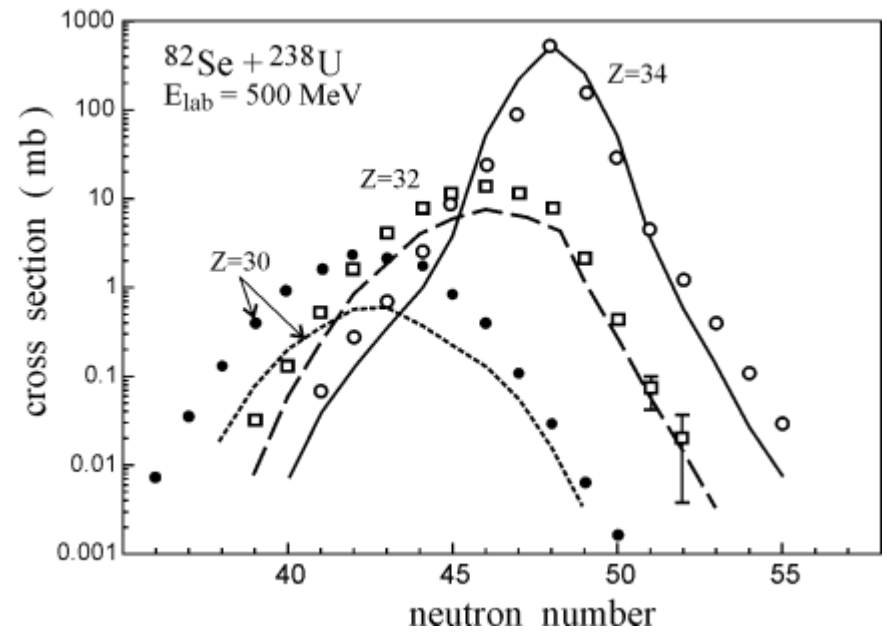
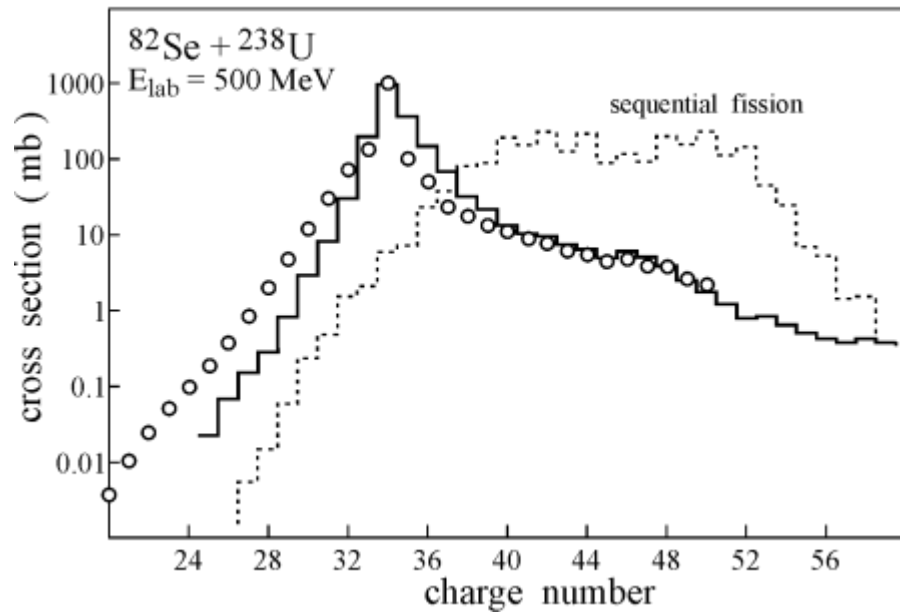
2. Langevin-type equations of motion



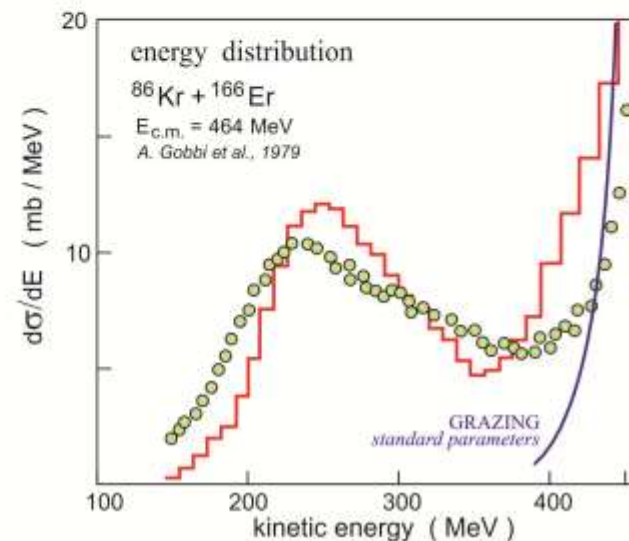
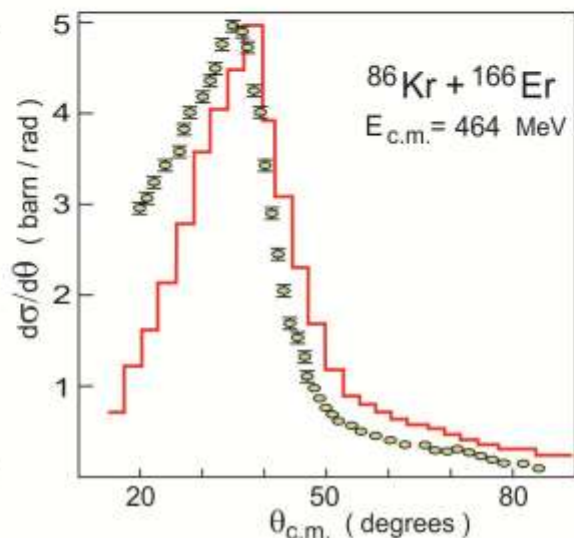
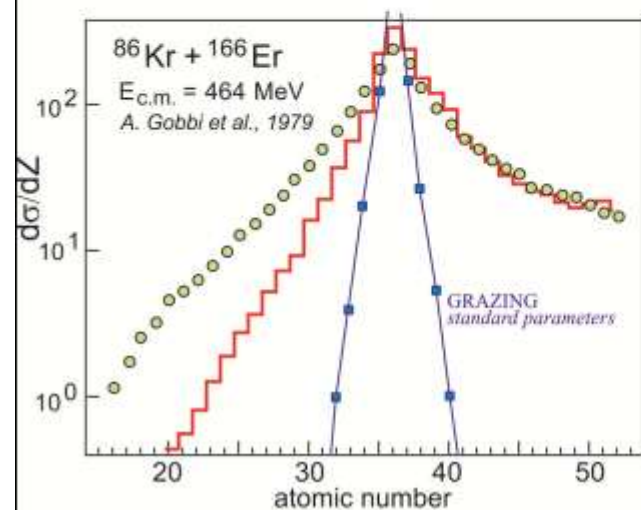
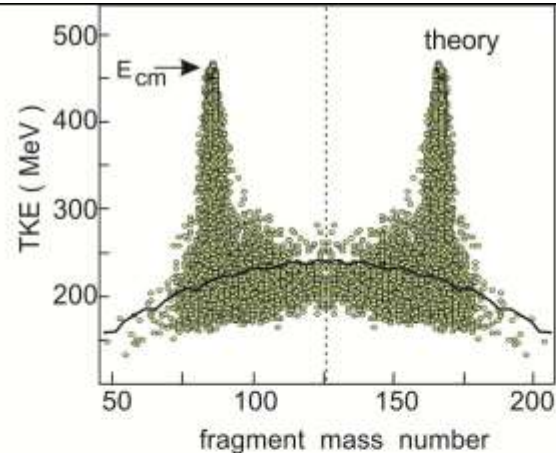
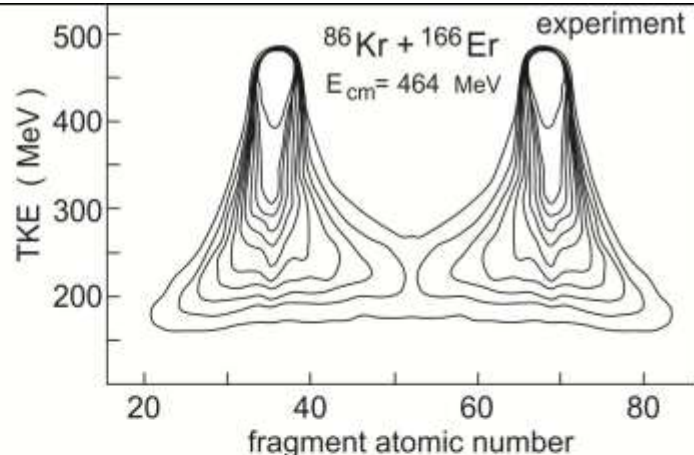
Dynamics: 10^6 tested events (trajectories),
 Statistical model: 10^{-6} ($3n$), 10^{-7} ($4n$) survival probability
 cross sections up to **0.1 pb** can be calculated

Comparison with experiments on multi-nucleon transfer

Experiment: L. Corradi et al., 2006

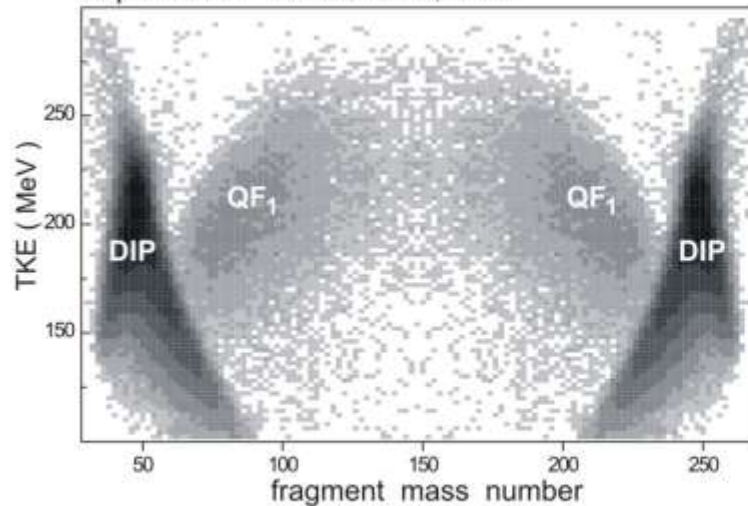


For the first time quantitative description of all the features of Deep Inelastic scattering was obtained

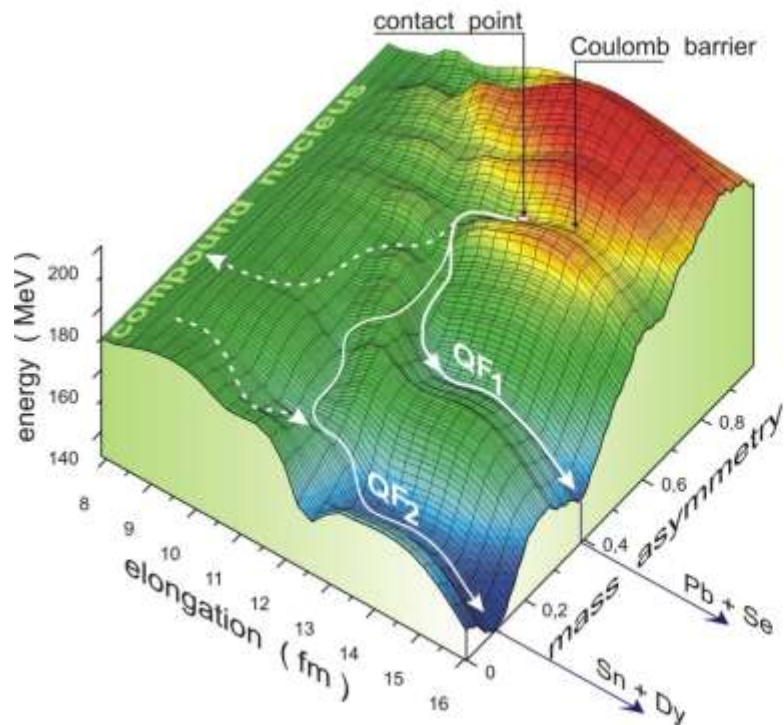
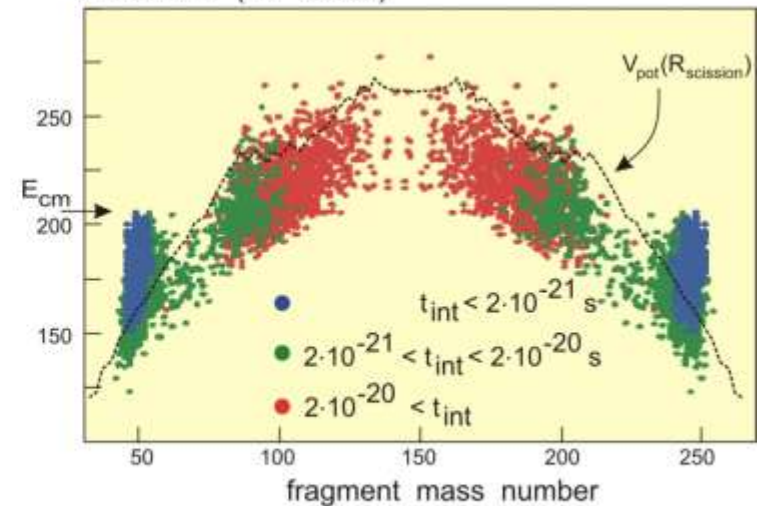


Quasi-Fission process: $^{48}\text{Ca} + ^{248}\text{Cm}$

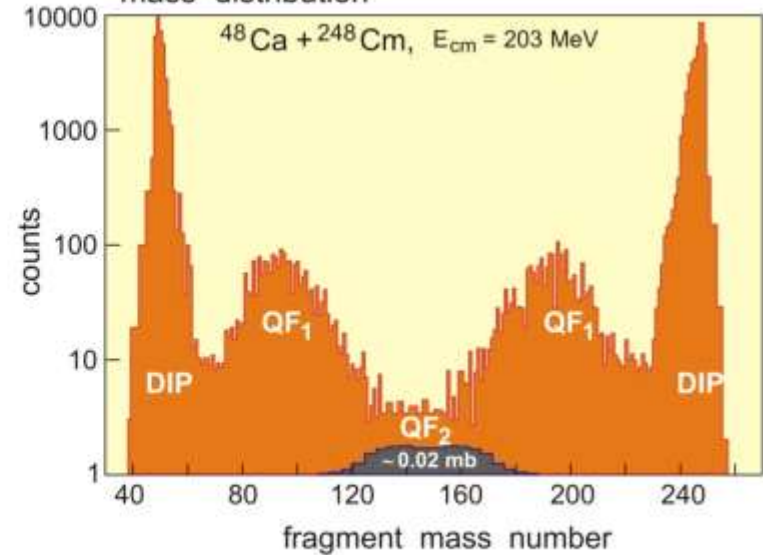
experiment: M. Itkis et al., 2000



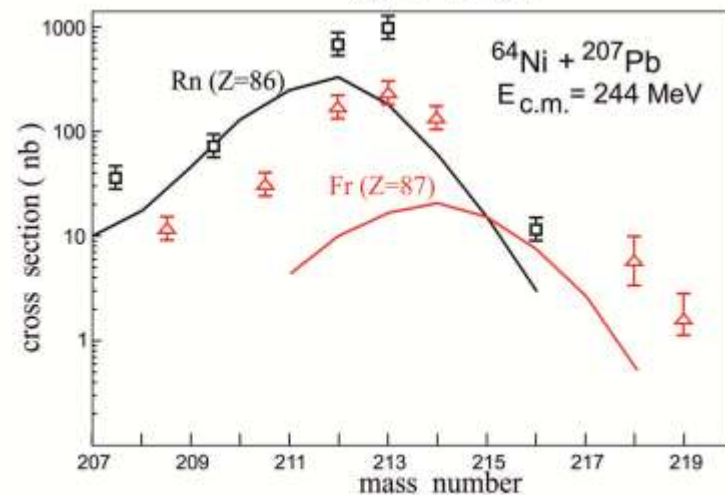
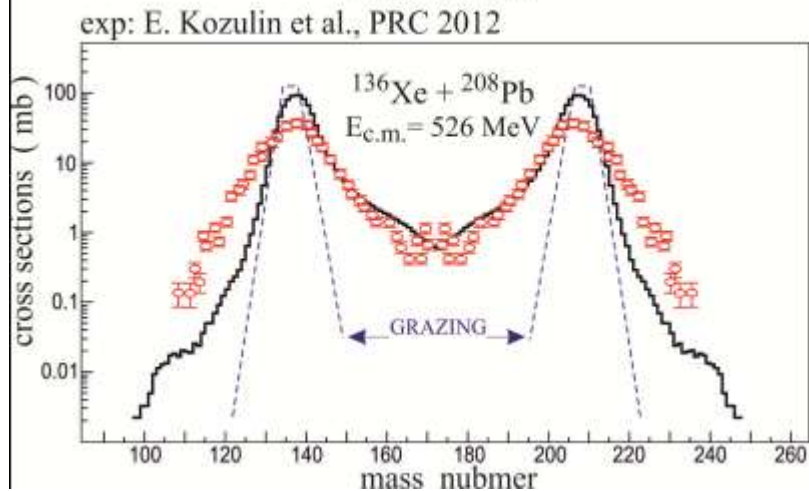
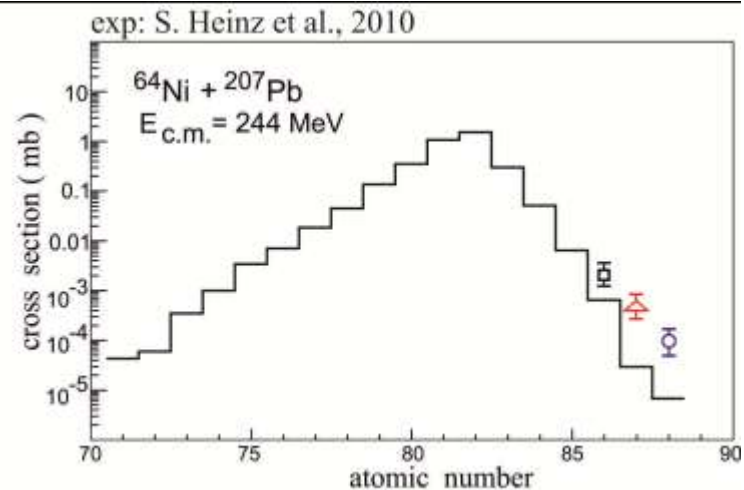
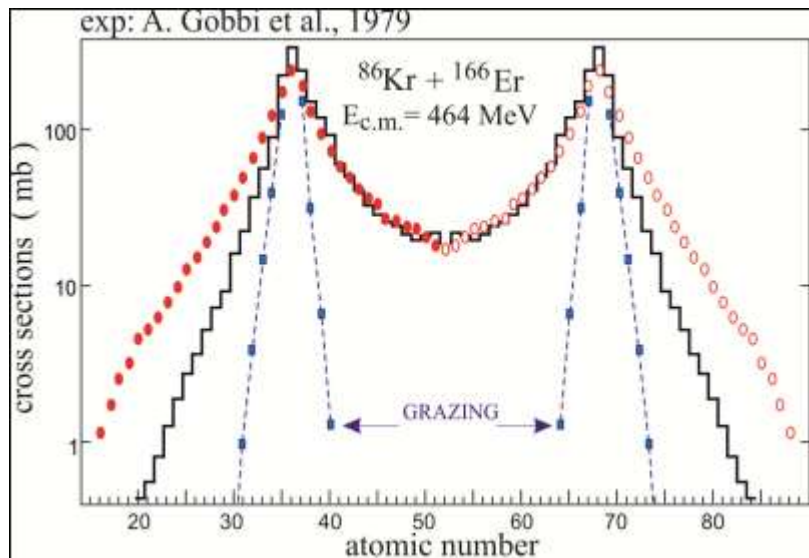
calculation (10^5 events)



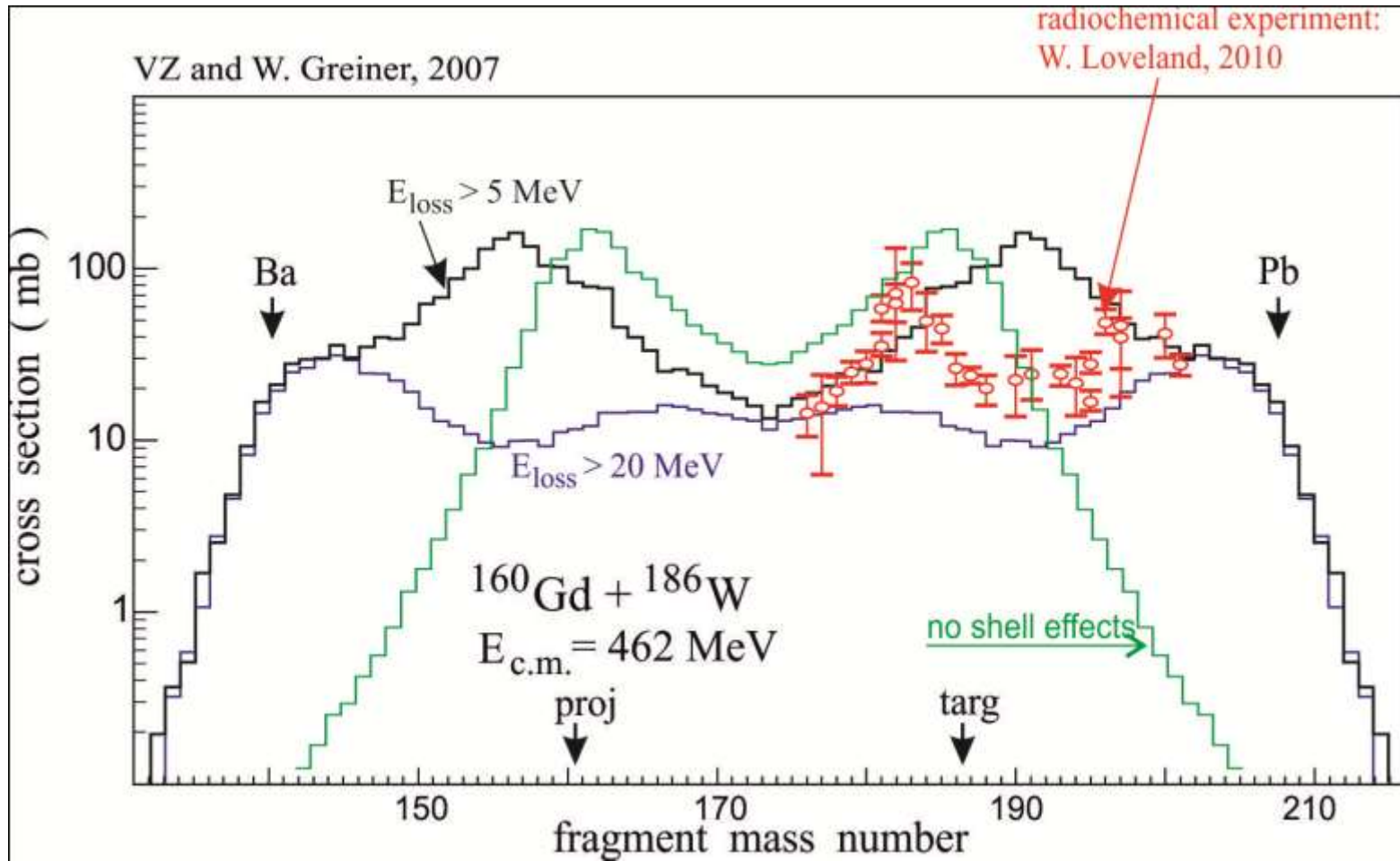
mass distribution



Underestimation of the yield of trans-target nuclei ?



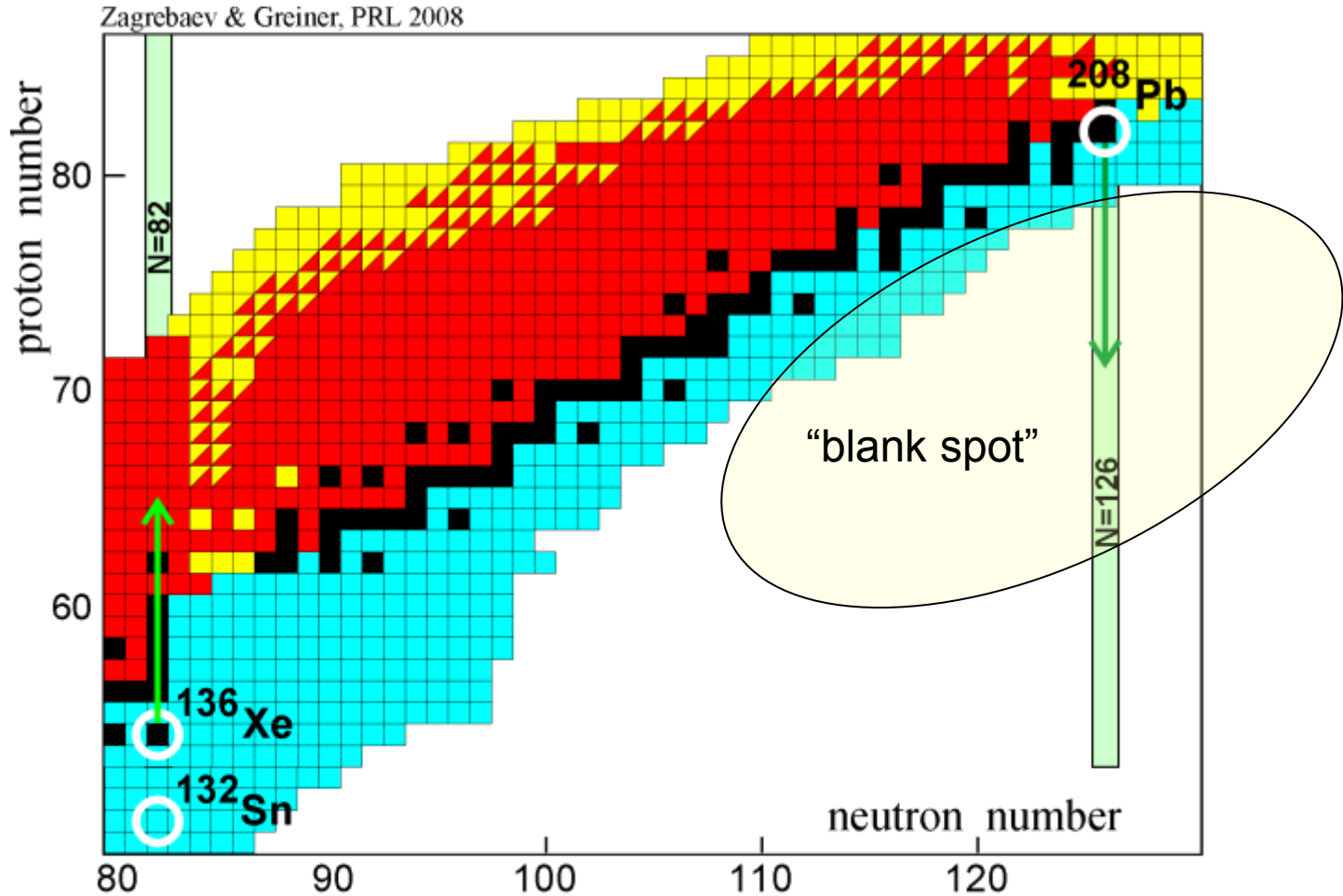
Shell effects in low-energy multi-nucleon transfer reactions ?



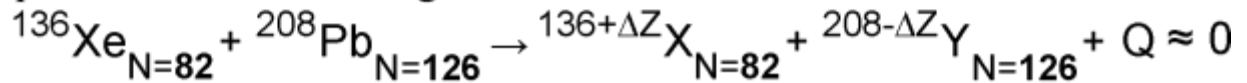
on-line Gd+W experiment started in Dubna on September 20

If the shell effects in transfer reactions will be confirmed
it will open a door for the production of neutron rich Superheavy Elements
in U+Cm collisions (inverse quasi-fission)

Production of new heavy nuclei in the region of N=126



proton transfer along the neutron closed shells:

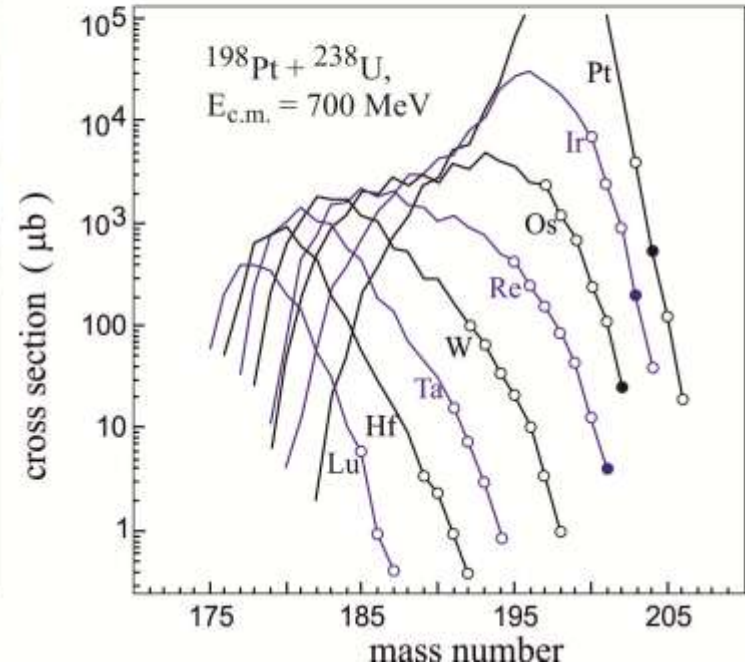
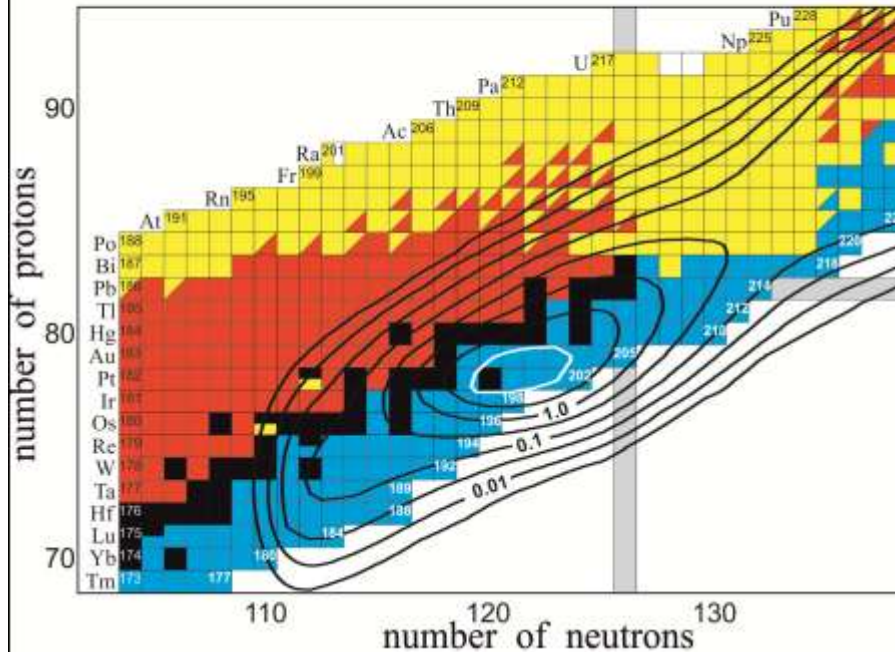
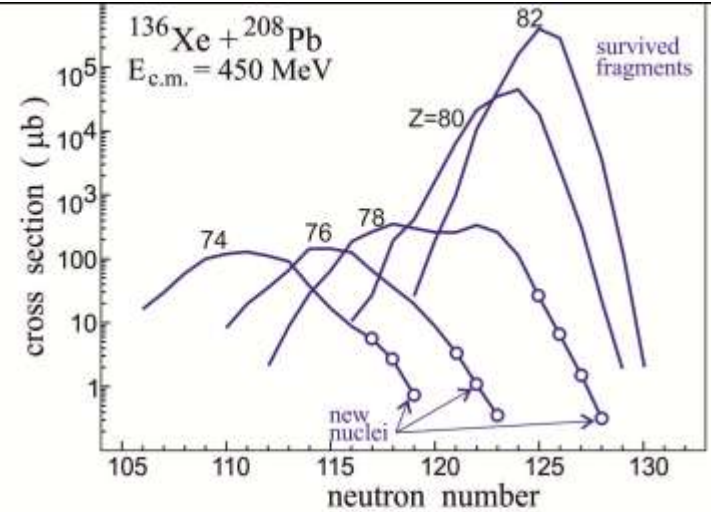
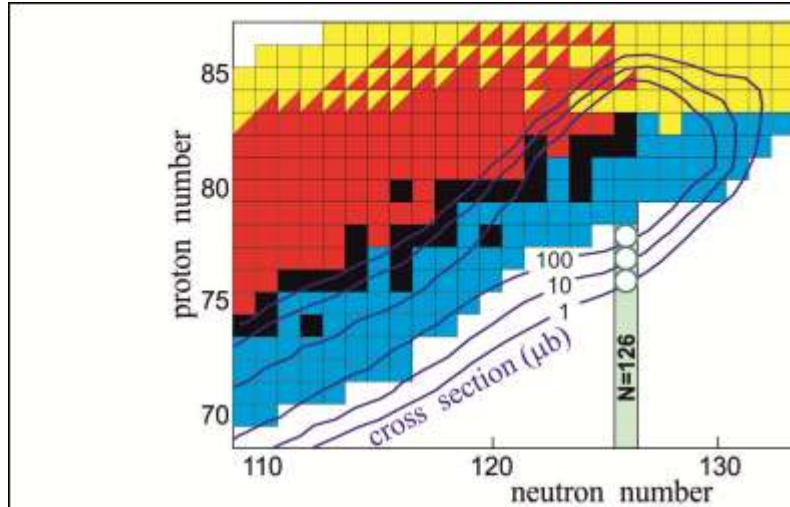


Reactions with $Q \approx 0$ are very favorable for proton transfer

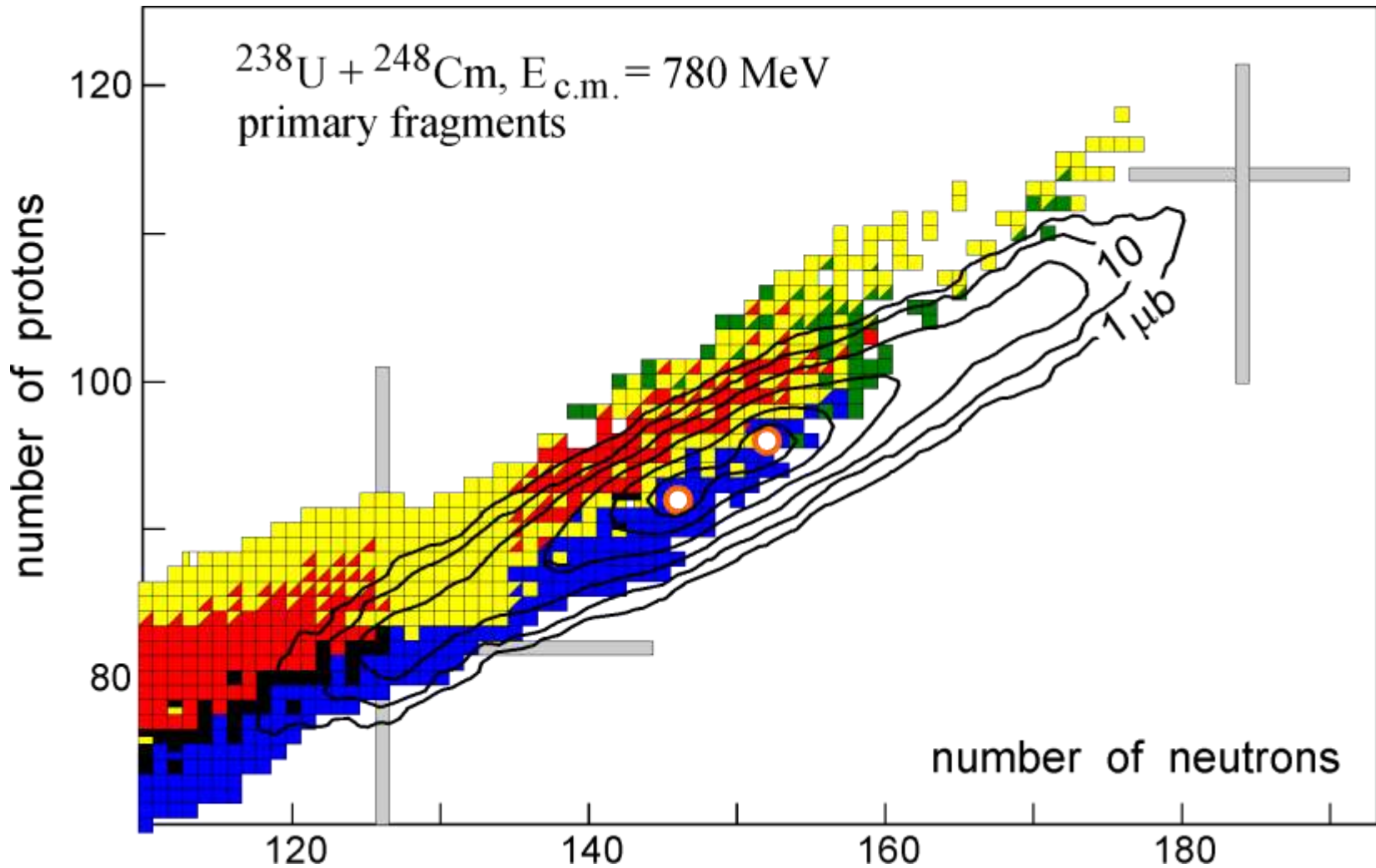
The use of ${}^{132}\text{Sn}$ is even better !

*Isotope production
with radioactive beams,
Dasso, Pollarolo, Winther,
PRL 1994*

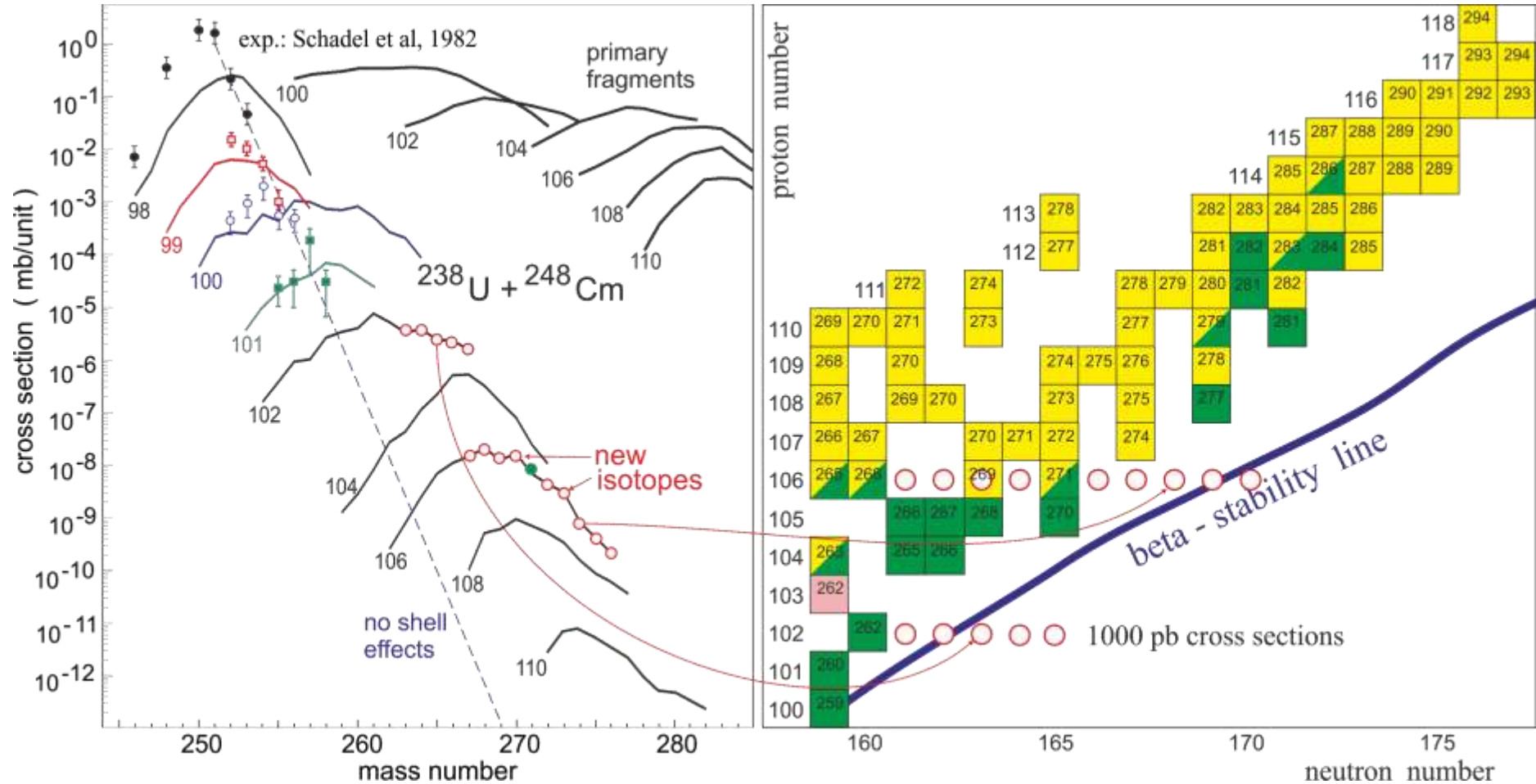
Production of neutron rich heavy nuclei located along the last “waiting point” of astrophysical nucleosynthesis



$^{238}\text{U} + ^{248}\text{Cm}$. Primary fragments



Production of transfermium nuclei along the line of stability looks quite possible



Rather wide angular distribution of reaction fragments:
separators of a new kind are needed

Selective laser ionization of Au & Hg atoms

Ionization Schemes

Au I

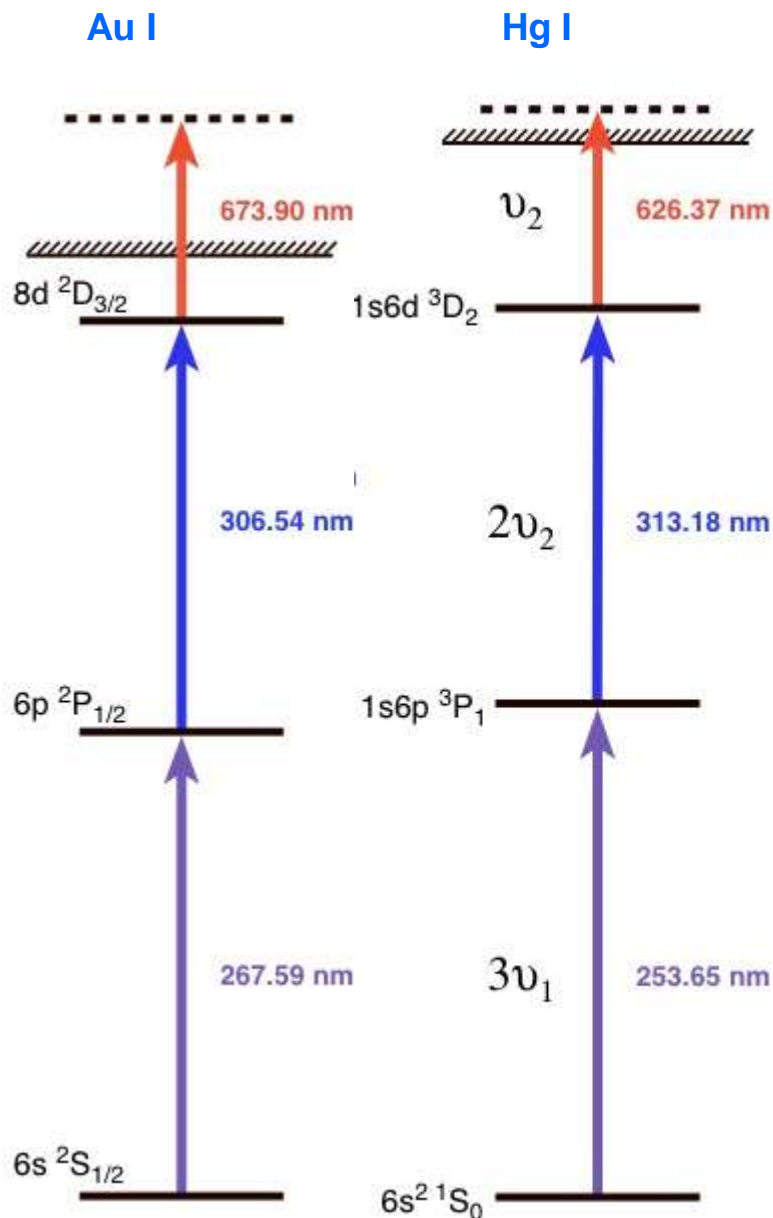
Level Energy, cm ⁻¹	Configuration	Wavelength, nm
E_0 0	$6s^2 2S_{1/2}$	λ_1 267.6
E_1 37358.99	$6p^2 P_{1/2}$	λ_2 306.5
E_2 69971.42	$8d^2 D_{3/2}$	λ_3 673.9

Chemical series: Transition metals
 Group, Period, Block: 11, 6, *d*
 Atomic mass: 196.96655(2) g/mol
 Electron configuration: [Xe] 4f¹⁴ 5d¹⁰ 6s
[Ionization potential](#): 74408.88 cm⁻¹ (9.22553 eV)

Hg I

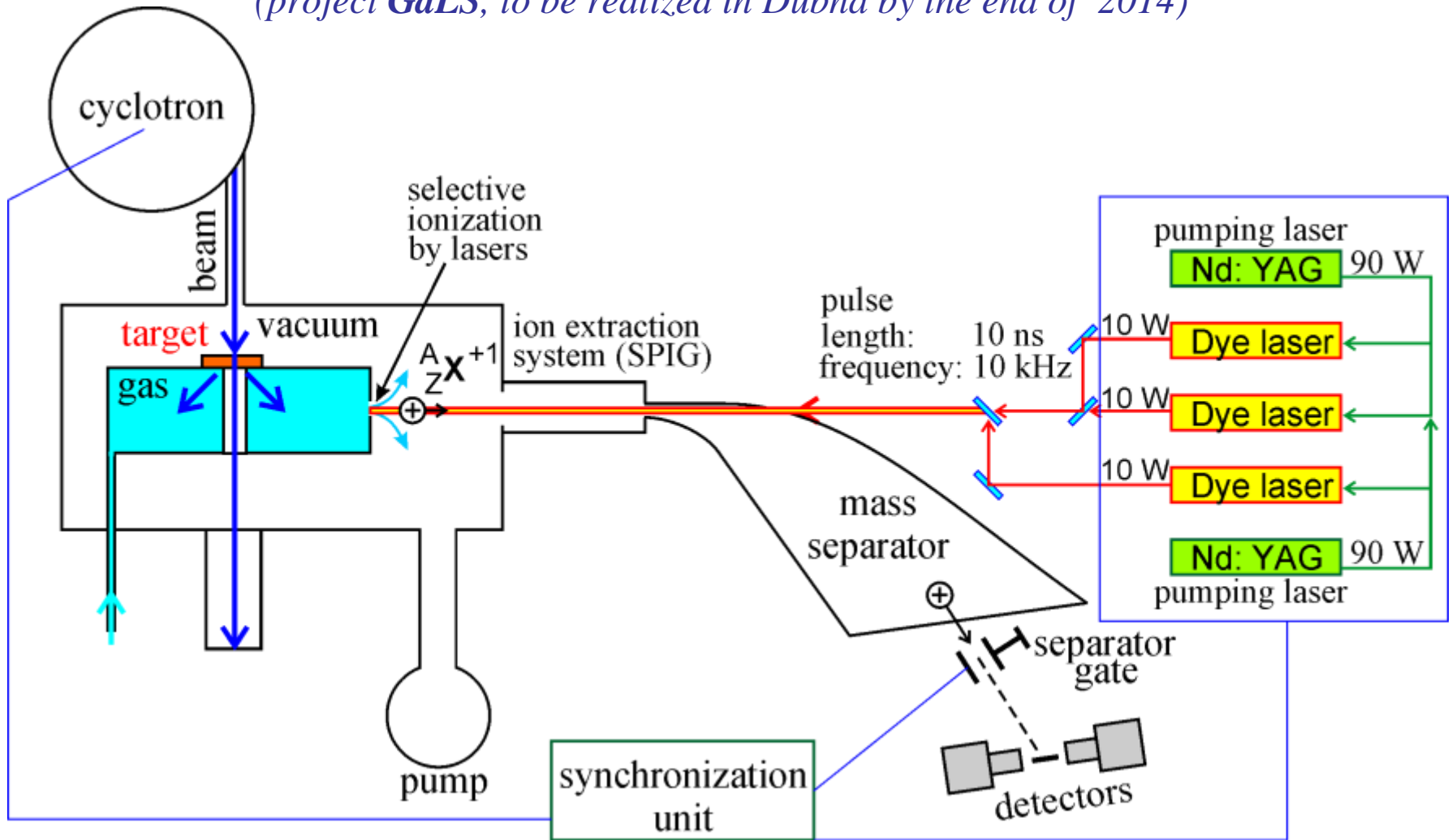
Level Energy, cm ⁻¹	Configuration	Wavelength, nm
E_0 0	$6s^2 1S_0$	λ_1 253.65
E_1 39412.30	$6s 6p^3 P_1$	λ_2 313.18
E_2 71396.22	$6s 6d^3 D_2$	λ_3 626.37

Chemical series: *d*
 Group, Period, Block: 12, 6, *d*
 Atomic mass: 200.59(2) g/mol
 Electron configuration: [Xe] 4f¹⁴ 5d¹⁰ 6s²
[Ionization potential](#): 84184.1 cm⁻¹ (10.4375 eV)

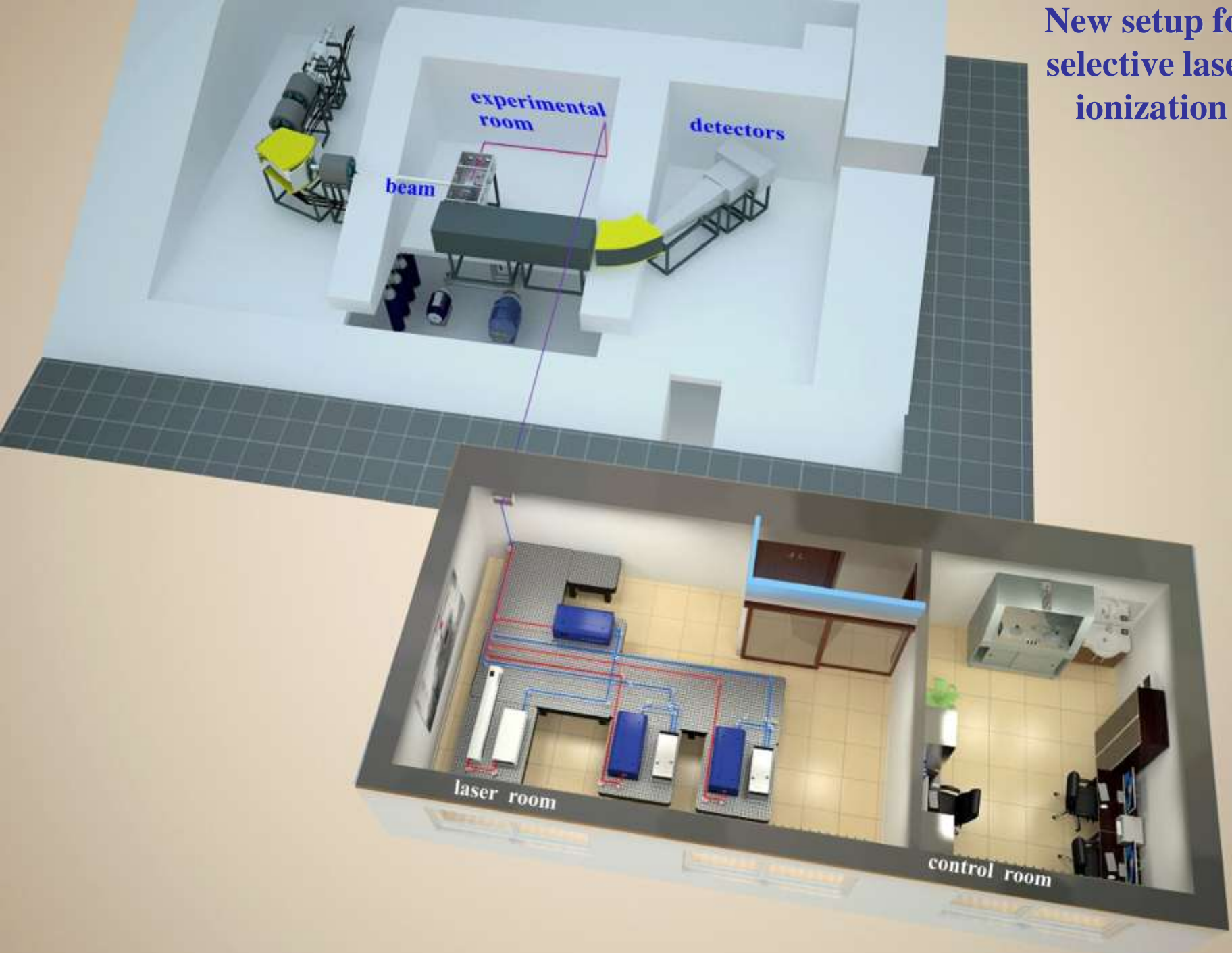


New setup for selective laser ionization of multi-nucleon transfer reaction products stopped in gas

(project *GaLS*, to be realized in Dubna by the end of 2014)



New setup for selective laser ionization



Summary

- North-east part of the nuclear map is still “**terra incognita**”. Heavy neutron rich nuclei are not synthesized and studied yet.
- Multi-nucleon transfer reactions can be used for synthesis of new **neutron rich nuclei** located along the closed neutron shell **N=126** having the largest impact on the astrophysical r-process. Cross sections are higher than 1 μb .
- Multi-nucleon transfer reactions can be used also for synthesis of new **neutron enriched transfermium nuclei** located along the beta-stability line. U-like beams are needed as well as a new kind of separators!
- Shell effects and dynamics of the “**inverse quasi-fission processes**” in HI damped collisions should be studied much better. There are several quite promising experiments to be perform just now.



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