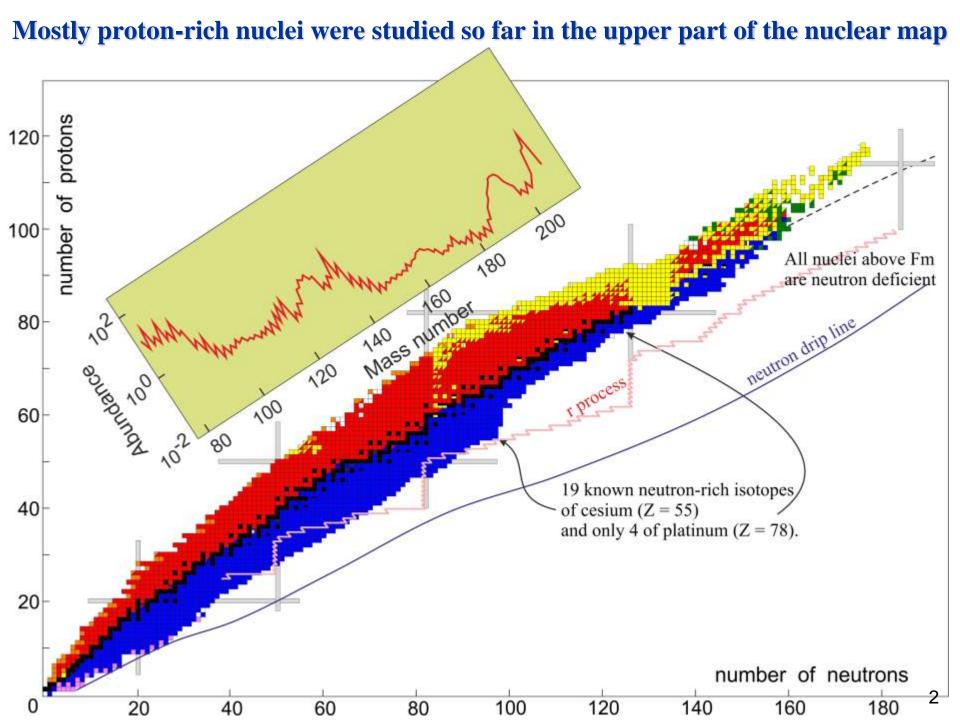
### Production of heavy and superheavy neutron rich nuclei

- Unexplored north-east area of the Nuclear Map
- Our predictions and proposals:
  - Elements 119 and 120 are on the way. What's the next?
  - Filling the gap of not-yet-synthesized isotopes of SH elements (Z=106 116)
  - Narrow (hypothetical) pathway to the Island of Stability
  - Production of new neutron enriched SH nuclei in transfer reactions
  - Shell effects in damped collisions of heavy ions?
  - Production of trans-target nuclei (inverse quasi-fission process)
  - Production of neutron rich nuclei located along the neutron closed shell N=126
- New facilities are needed. They are coming



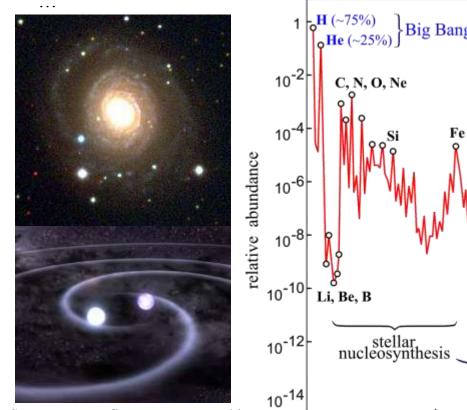
Valeriy Zagrebaev



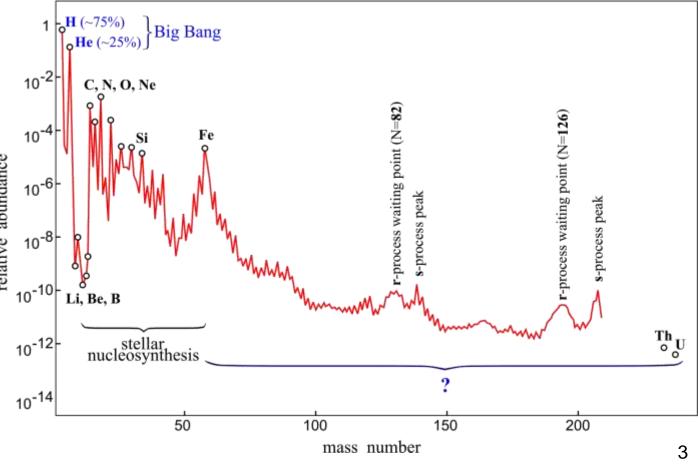
#### **Abundance of the elements in the Universe**

## The 11 Greatest Unanswered Questions of Physics (National Research Council, NAS, USA, 2002):

- 1. What is dark matter?
- 2. What is dark energy?
- 3. How were the heavy elements from iron to uranium made?
- 4. Do neutrinos have mass?



Strong neutron fluxes are expected in **core-collapse supernova explosions** or in the **mergers of neutron stars.** 

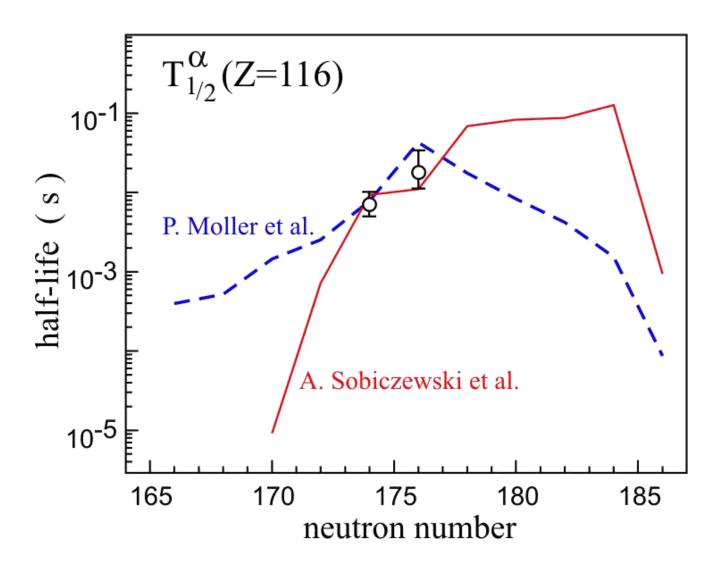


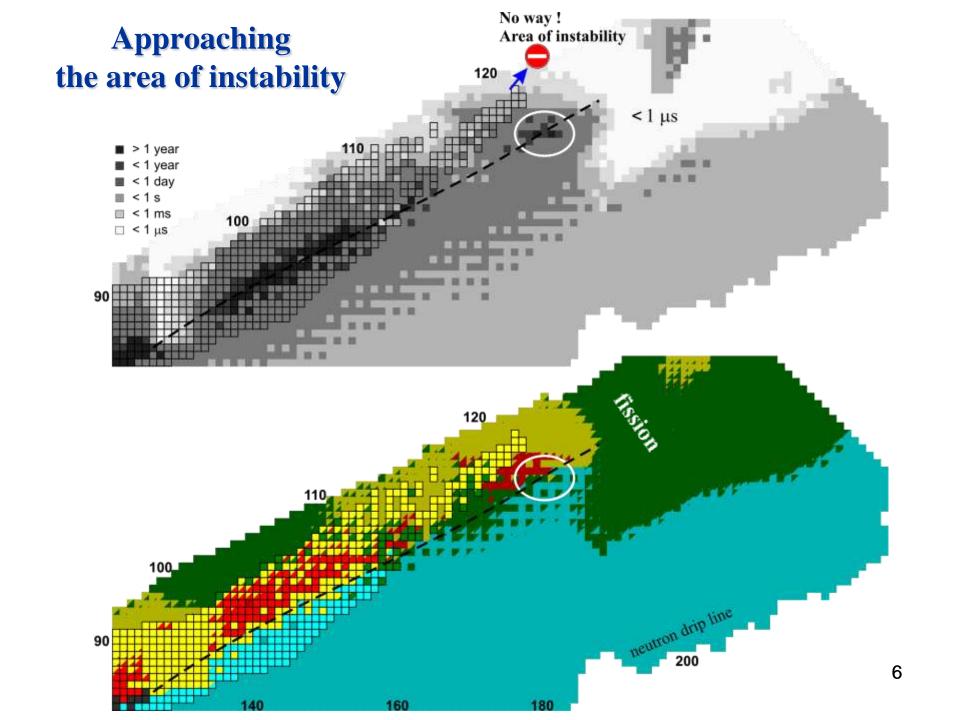
SHE: we are far from the stability line and from the Island of Stability <sup>50</sup>Ti + <sup>249</sup>Cf → **120** 120 ← <sup>54</sup>Cr + <sup>248</sup>Cm 50Ti + 249Bk → 119 protons 118 48Ca + 239Pu <sup>48</sup>Ca + <sup>243</sup>Cm 48Ca + 241Am 117 116 Lv 115 115 Island 35 40 45 50 excitation energy ( MeV ) 35 40 45 50 excitation energy (MeV) 114 FI of Stability Rn 110 Ds Mt Hs beta-stability line Bh Sg 105 Db Rf No Md Cr beam: Ti beam: Fm SHIP (May, 2011) TASCA (August, 2011) TASCA (October, 2012) Es 50Ti + 249Bk → 299119 50Ti + 249Cf → 299120 54Cr + 248Cm → 302120 ( gd o ~ 50 fb o ~ 40 fb o ~ 25 fb Cf Bk cross section Cm Am Pu E\* (MeV) 50 Np E\* (MeV) 50 40 E\* (MeV) 50 155 160 165 170 180 175

U

neutrons

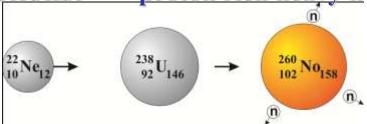
#### Our ability of predictions in superheavy mass area



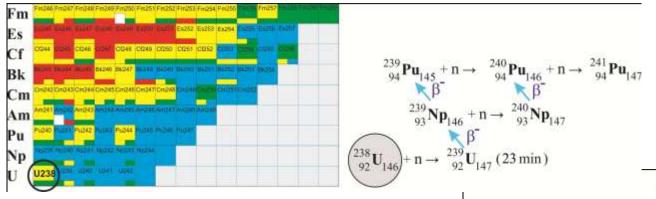


### There are only 3 methods for synthesis of heavy nuclei

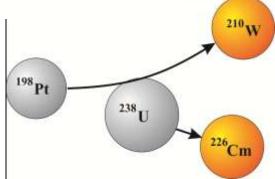
1. Fusion reactions:  $\rightarrow$  proton rich heavy nuclei



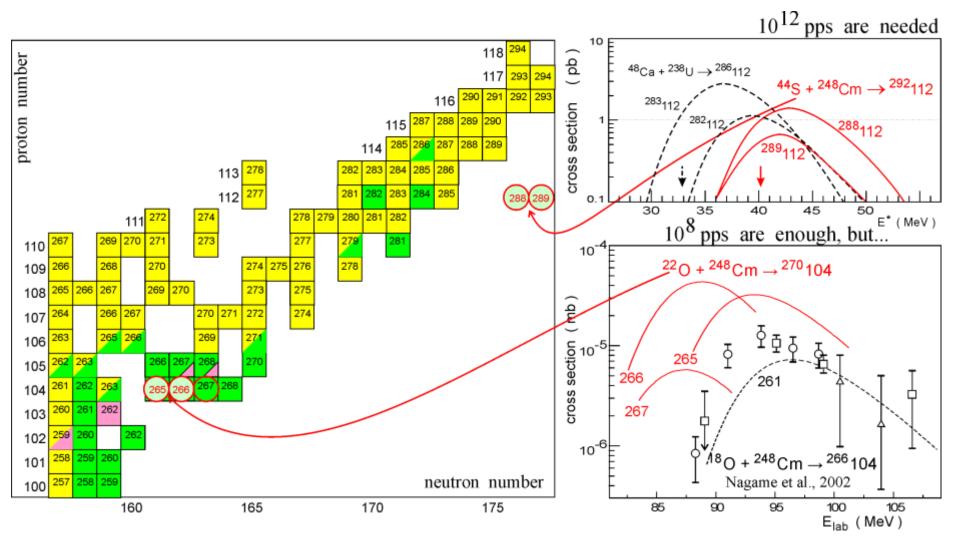
2. Sequence of neutron capture and beta(-) decay processes: neutron fluxes in reactors are too low, nuclear explosions are forbidden



3. Multi-nucleon transfer reactions



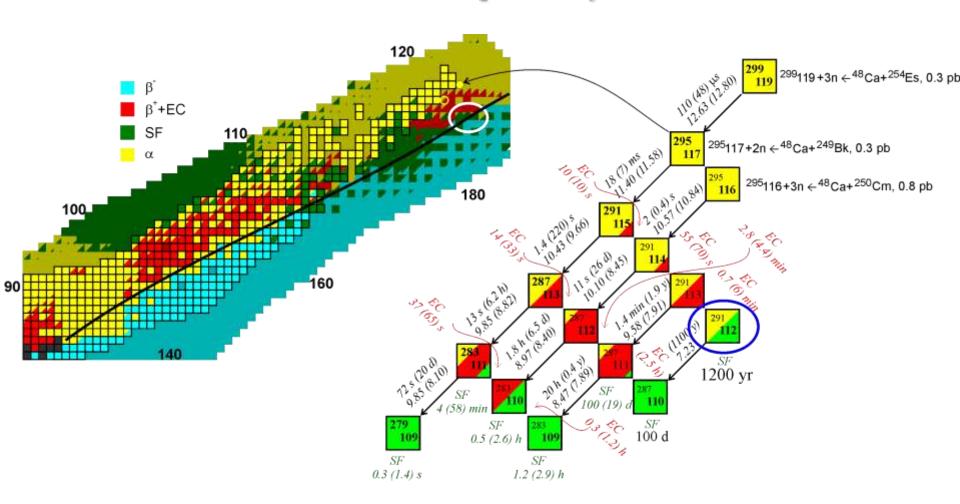
## The use of low-energy Radioactive Ion Beams for the production of neutron rich superheavy nuclei?





No chances today. But in the nearest future ...

## Narrow (hypothetical) pathway to the Island of Stability still exists probably!



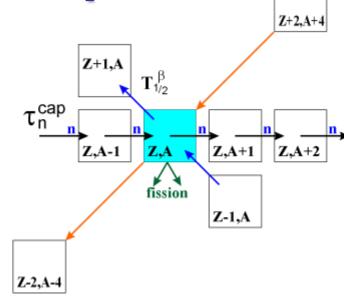
### Nucleosynthesis by neutron capture

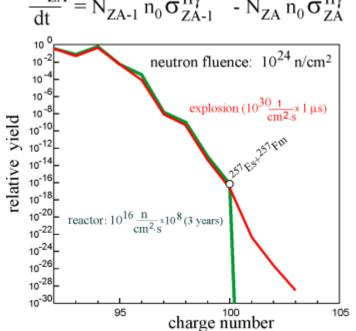
 $n_0$  is the neutron flux time of neutron capture  $\tau_n^{\text{cap}} = \frac{1}{n_0^x \sigma(n,\gamma)}$ 

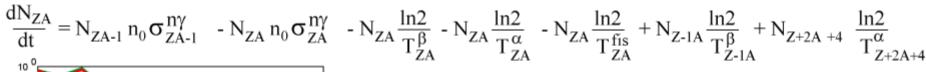
$$(Z,A) \rightarrow (Z,A+1)$$
 if  $T_{1/2} > \tau_n^{cap}$ 

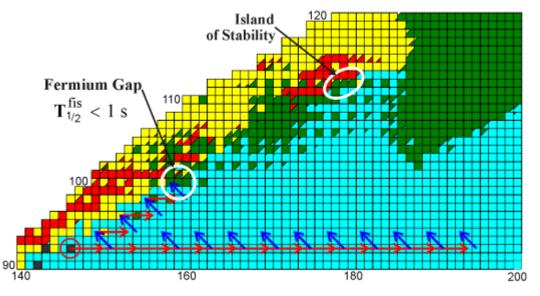
nuclear reactor:  $\tau_n^{cap} \sim 1 \text{ year}$ 

nuclear explosion:  $\tau_n^{cap} \sim 1 \ \mu s$ 









#### Theoretical models of transfer reactions

#### Multi-nucleon transfers in damped collisions

#### **Master equation**

L.G. Moretto and J.S. Sventek, Phys. Lett. B **58**, 26 (1975)

#### **Fokker-Plank equation**

W. Norenberg, Phys. Lett. B **52**, 289 (1974)

#### **Langevin equations**

P. Frobrich and S.Y. Xu, Nucl. Phys. **A477**, 143 (1988)

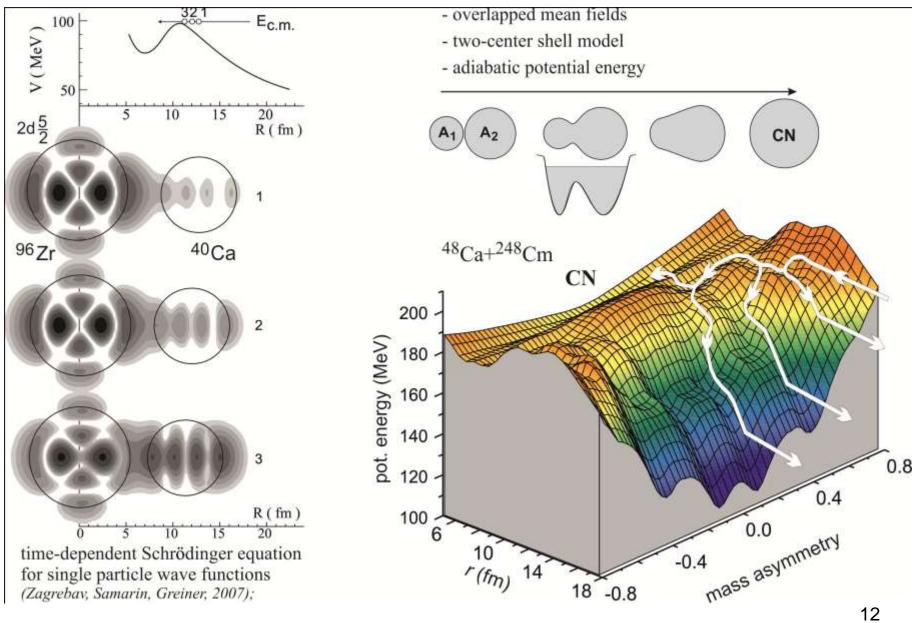
#### **Semi-classical approaches**

E. Vigezzi and A. Winther, Ann. Phys. (N.Y.) 192, 432 (1989).V.I. Zagrebaev, Ann. Phys. (N.Y.) 197, 33 (1990).

#### Few-nucleon transfers (GRAZING)

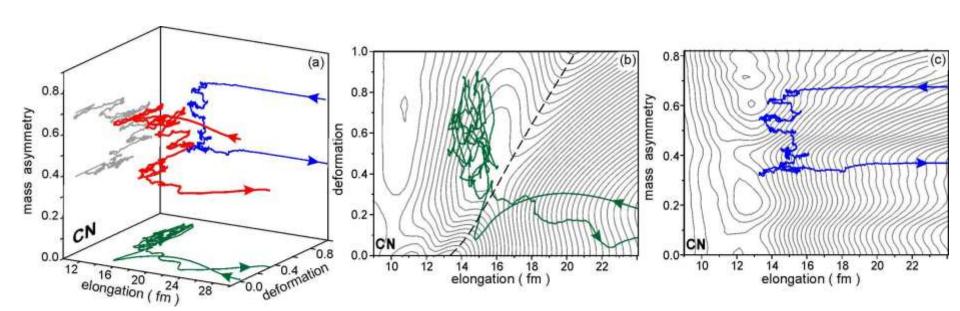
A. Winther, Nucl. Phys. **A594**, 203 (1995) http://personalpages.to.infn.it/nanni/grazing

#### Adiabatic dynamics of low-energy heavy ion collisions and nucleon transfers



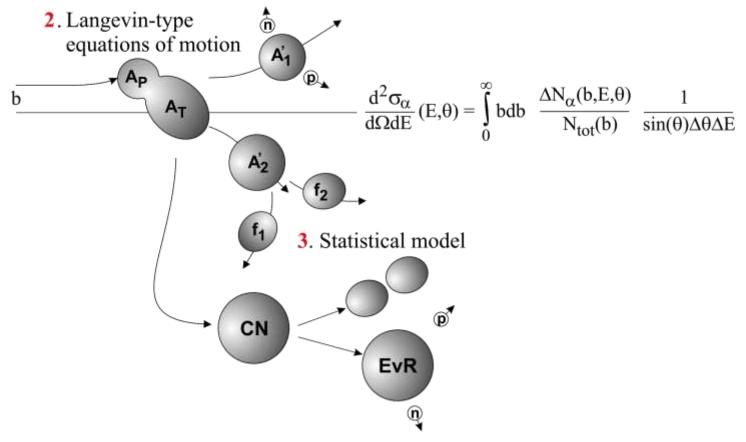
$$\frac{dR}{dt} = \frac{P_R}{\mu_R}$$
 Variables:  $\{R, \theta, \phi_1, \phi_2, \beta_1, \beta_2, \eta_Z, \eta_N\}$  Most uncertain parameters:  $\mu_0, \gamma_0$  - nuclear viscosity and friction,  $\frac{d\phi_1}{dt} = \frac{L_1}{3_1}, \frac{d\phi_2}{dt} = \frac{L_2}{3_2}$   $\frac{d\beta_1}{dt} = \frac{P_{\beta 1}}{\mu_{\beta 1}}$   $\frac{d\theta_2}{dt} = \frac{P_{\beta 2}}{\mu_{\beta 2}}$   $\frac{d\theta_1}{dt} = \frac{P_{\beta 2}}{\mu_{\beta 2}}$   $\frac{d\eta_N}{dt} = \frac{2}{N_{CN}} D_N^{(1)} + \frac{2}{N_{CN}} \sqrt{D_N^{(2)}} \Gamma_N(t)$  
$$\frac{d\rho_N}{dt} = \frac{\partial V}{\partial R} + \frac{\ell^2}{\mu_R R^3} + \left(\frac{\ell^2}{2\mu_R^2 R^2} + \frac{P_R^2}{2\mu_R^2}\right) \frac{\partial \mu_R}{\partial R} + \frac{P_{\beta 1}^2}{2\mu_{\beta 1}^2} \frac{\partial \mu_{\beta 1}}{\partial R} + \frac{P_{\beta 2}^2}{2\mu_{\beta 2}^2} \frac{\partial \mu_{\beta 2}}{\partial R} - \gamma_R \frac{P_R}{\mu_R} + \sqrt{\gamma_R} T \Gamma_R(t)$$
 
$$\frac{d\ell}{dt} = -\frac{\partial V}{\partial \Phi} + \gamma_{tang} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{3_1} a_1 - \frac{L_2}{3_2} a_2\right) R + \sqrt{\gamma_{tang}} T \Gamma_{tang}(t)$$
 
$$\frac{dL_1}{dt} = -\frac{\partial V}{\partial \phi_1} + \gamma_{tang} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{3_1} a_1 - \frac{L_2}{3_2} a_2\right) a_2 - \frac{a_2}{R} \sqrt{\gamma_{tang}} T \Gamma_{tang}(t)$$
 
$$\frac{dL_2}{dt} = -\frac{\partial V}{\partial \phi_2} + \gamma_{tan} \left(\frac{\ell}{\mu_R R} - \frac{L_1}{3_1} a_1 - \frac{L_2}{3_2} a_2\right) a_2 - \frac{a_2}{R} \sqrt{\gamma_{tang}} T \Gamma_{tang}(t)$$
 
$$\frac{d\rho_{\beta 1}}{dt} = -\frac{\partial V}{\partial \phi_1} + \frac{P_{\beta 1}^2}{2\mu_{\beta 1}^2} \frac{\partial \mu_{\beta 1}}{\partial \beta_1} + \frac{P_{\beta 2}^2}{2\mu_{\beta 2}^2} \frac{\partial \mu_{\beta 2}}{\partial \beta_1} + \frac{\ell^2}{2\mu_R^2 R^2} + \frac{P_R^2}{2\mu_R^2} \frac{\partial \mu_R}{\partial \beta_1} - \gamma_\beta \frac{P_{\beta 1}}{\mu_{\beta 1}} + \sqrt{\gamma_{\beta 1}} T \Gamma_{\beta 1}(t)$$
 
$$\frac{d\rho_{\beta 1}}{dt} = -\frac{\partial V}{\partial \phi_2} + \frac{P_{\beta 1}^2}{2\mu_{\beta 1}^2} \frac{\partial \mu_{\beta 1}}{\partial \beta_1} + \frac{P_{\beta 2}^2}{2\mu_{\beta 2}^2} \frac{\partial \mu_{\beta 2}}{\partial \beta_1} + \frac{\ell^2}{2\mu_R^2 R^2} + \frac{P_R^2}{2\mu_R^2} \frac{\partial \mu_R}{\partial \beta_2} - \gamma_\beta \frac{P_{\beta 1}}{\mu_{\beta 1}} + \sqrt{\gamma_{\beta 1}} T \Gamma_{\beta 1}(t)$$
 
$$\frac{d\rho_{\beta 1}}{dt} = -\frac{\partial V}{\partial \phi_2} + \frac{P_{\beta 1}^2}{2\mu_{\beta 1}^2} \frac{\partial \mu_{\beta 1}}{\partial \beta_1} + \frac{P_{\beta 2}^2}{2\mu_{\beta 2}^2} \frac{\partial \mu_{\beta 2}}{\partial \beta_1} + \frac{\ell^2}{2\mu_R^2 R^2} + \frac{P_R^2}{2\mu_R^2} \frac{\partial \mu_R}{\partial \beta_2} - \gamma_\beta \frac{P_{\beta 1}}{\mu_{\beta 1}} + \sqrt{\gamma_{\beta 1}} T \Gamma_{\beta 2}(t)$$
 
$$\frac{d\rho_{\beta 1}}{dt} = -\frac{\partial V}{\partial \phi_2} + \frac{P_{\beta 1}^2}{2\mu_{\beta 1}^2} \frac{\partial \mu_{\beta 1}}{\partial \beta_2} + \frac{P_{\beta 2}^2}{2\mu_{\beta 2}^2} \frac{\partial \mu_{\beta 2}}{\partial \beta_2} + \frac{\ell^2}{2\mu_R^2 R^2} \frac{\partial \mu_R}{\partial \beta_2} - \gamma_\beta \frac{P_{\beta 1}}{2\mu_R^2} + \sqrt{\gamma_{\beta 2}} T \Gamma_{\beta 2}(t)$$

## Typical trajectory in the "distance-deformation-mass asymmetry" space (48Ca + 248Cm, E=210 MeV)



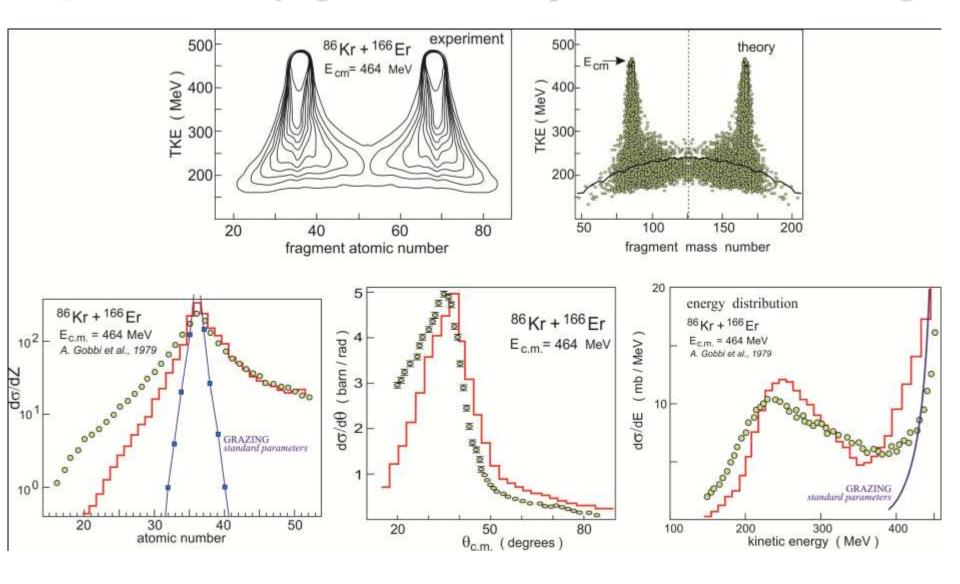
### Simulation of experiment. Cross sections

Time-dependent driving potential V(r,ξ;t):
Folding → Adiabatic Two-Center Shell Model



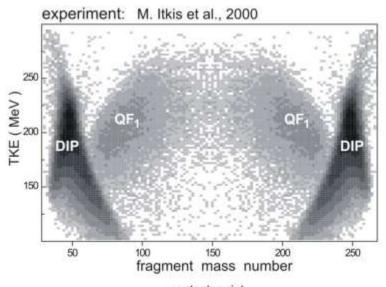
Dynamics: **10<sup>6</sup>** tested events (trajectories), Statistical model: **10<sup>-6</sup>**(3n), **10<sup>-7</sup>**(4n) survival probability cross sections up to **0.1 pb** can be calculated

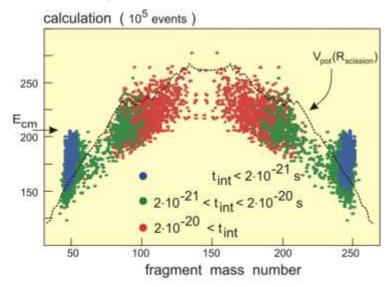
### Quite satisfactory agreement with experiments on DI scattering

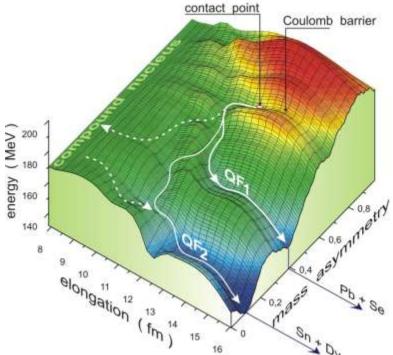


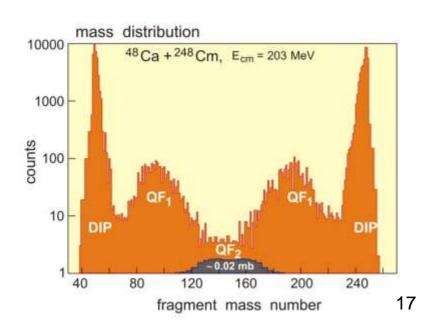
### Quasi-Fission process is understood quite well

(example:  ${}^{48}Ca + {}^{248}Cm$ )

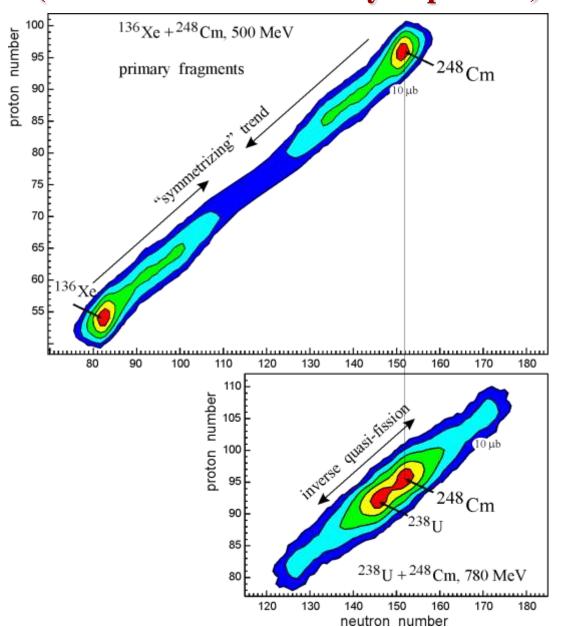




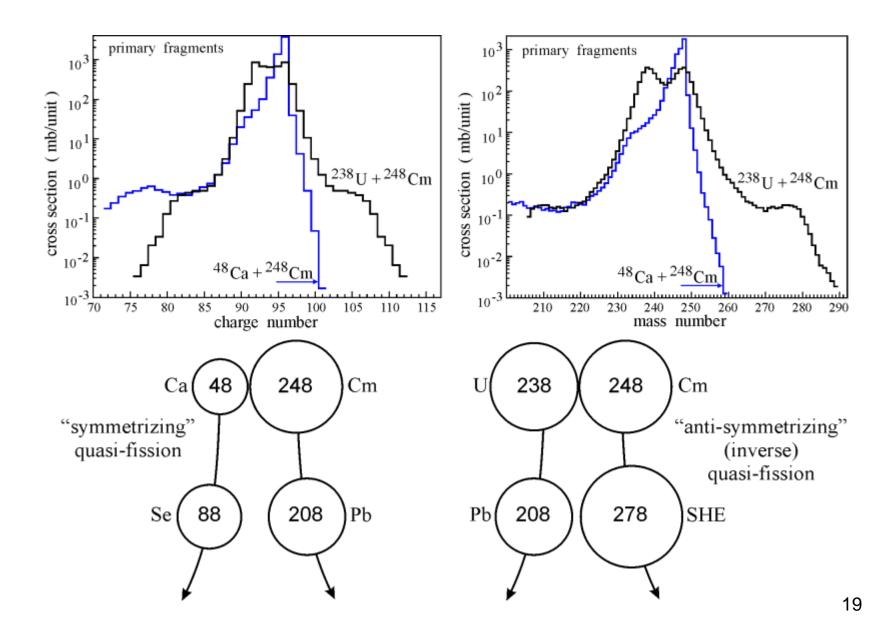




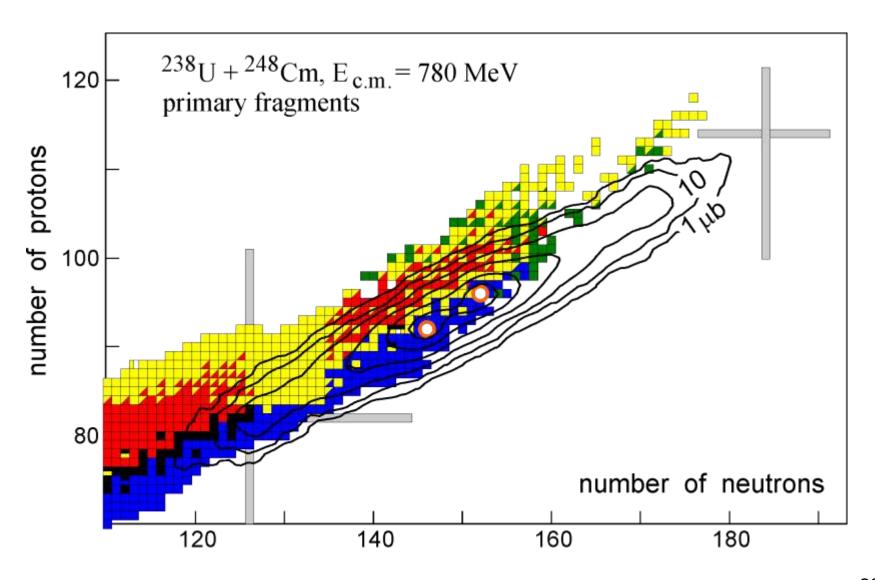
## Multi-nucleon transfers for production of Super-Heavy Elements (choice of reaction is very important)



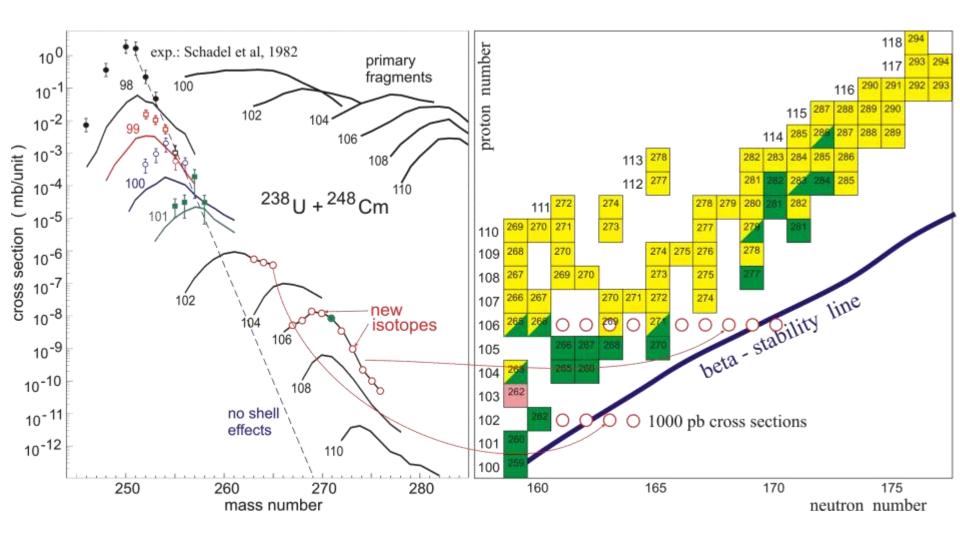
### "Inverse quasi-fission" reactions



### 238U + 248Cm. Primary fragments

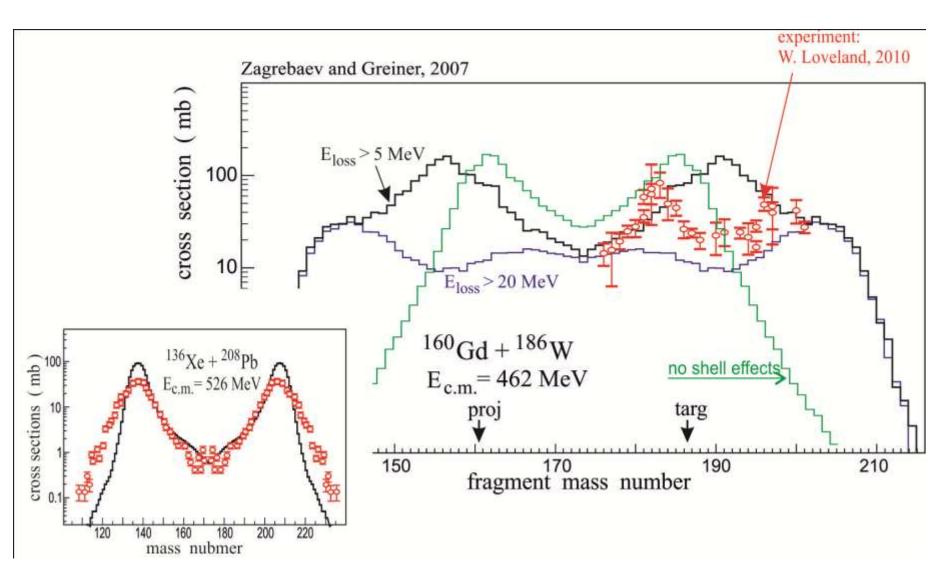


## Production of transfermium nuclei along the line of stability looks quite possible owing to shell effects

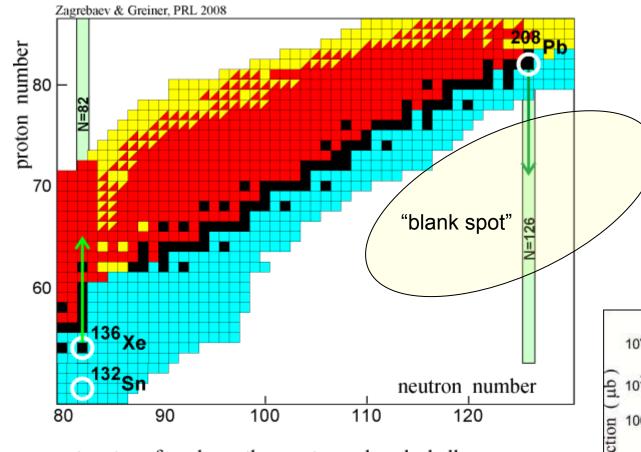


Rather wide angular distribution of reaction fragments: a new kind of separator is needed

## Shell effects in low-energy multi-nucleon transfer reactions is not studied yet!



### Production of new heavy nuclei in the region of N=126

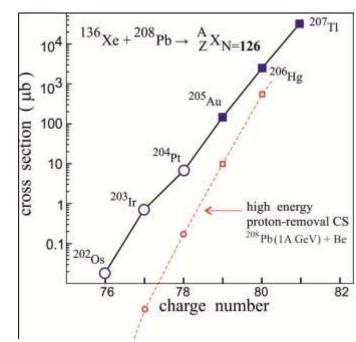


proton transfer along the neutron closed shells:

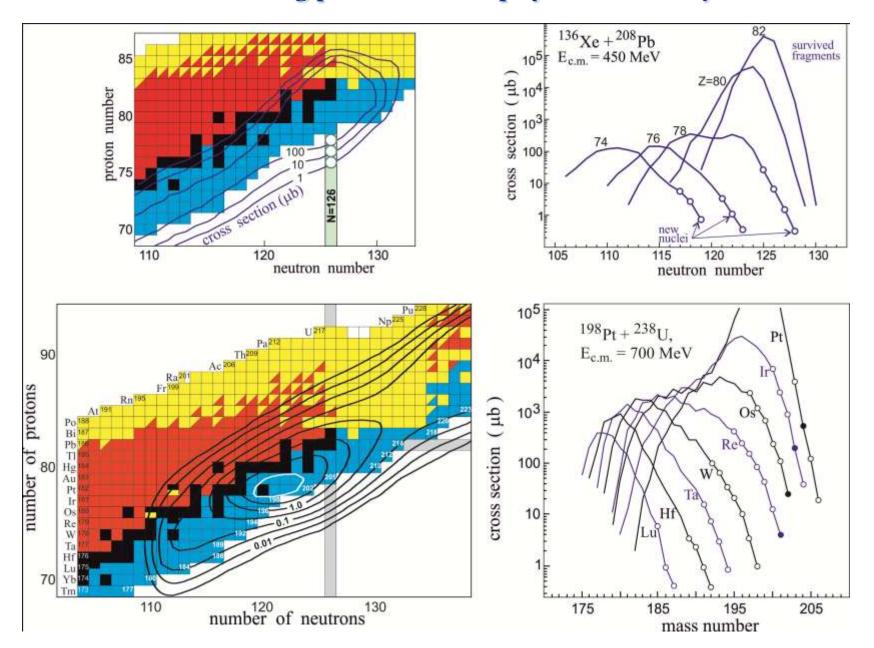
$$^{136}$$
Xe<sub>N=82</sub> +  $^{208}$ Pb<sub>N=126</sub>  $\rightarrow ^{136+\Delta Z}$ X<sub>N=82</sub> +  $^{208-\Delta Z}$ Y<sub>N=126</sub> + Q ≈ 0

Reactions with  $Q \approx 0$  are very favorable for proton transfer The use of  $^{132}$ Sn is even better!

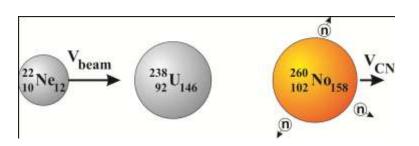
Isotope production with radioactive beams(54Ca, 154Xe), Dasso, Pollarolo, Winther, PRL 1994



## Production of neutron rich heavy nuclei located in the region of the last "waiting point" of astrophysical nucleosynthesis

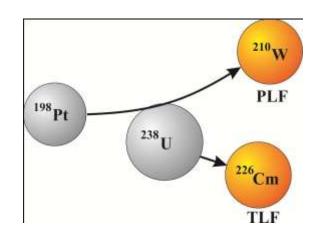


**Fusion reactions** 

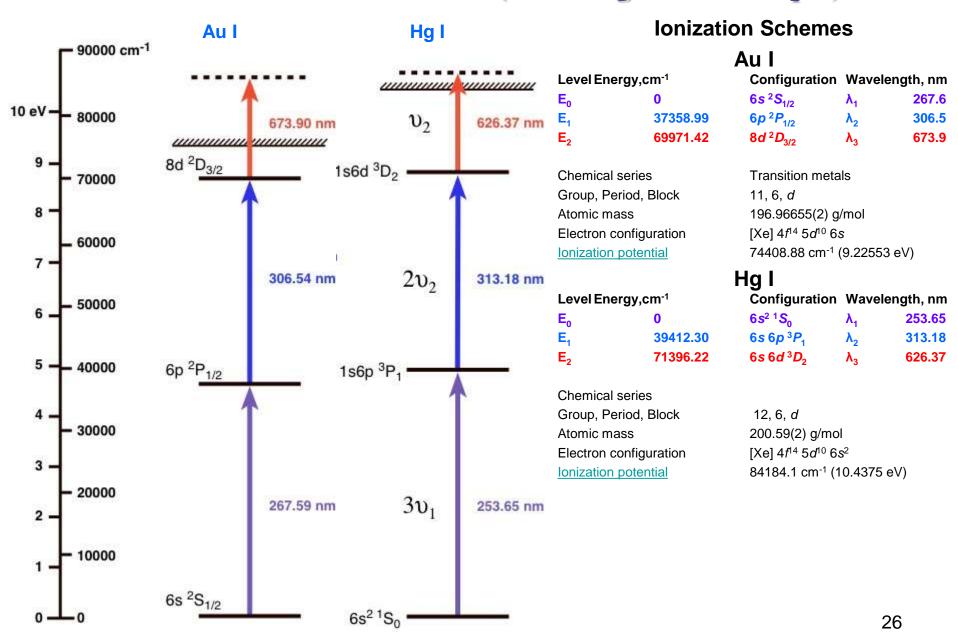


# How to separate a given nucleus from all the other transfer reaction products?

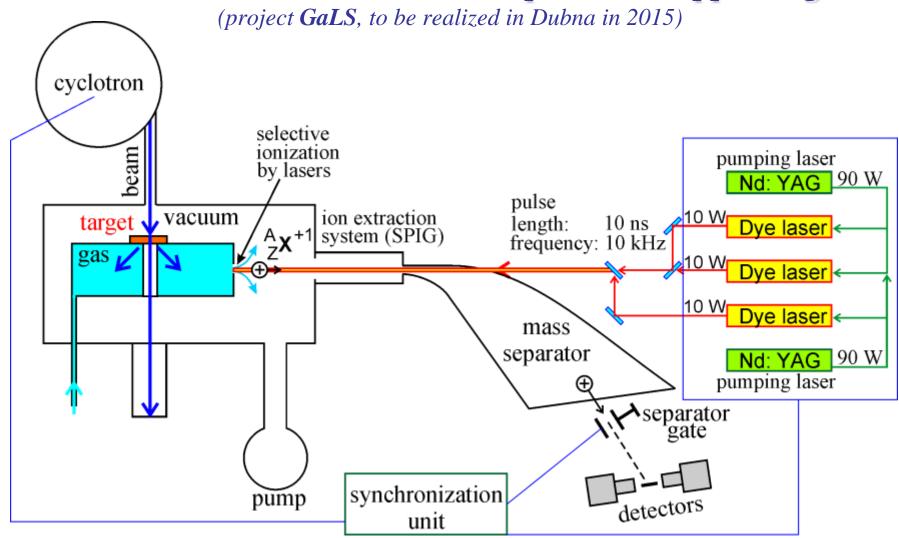
Transfer reactions

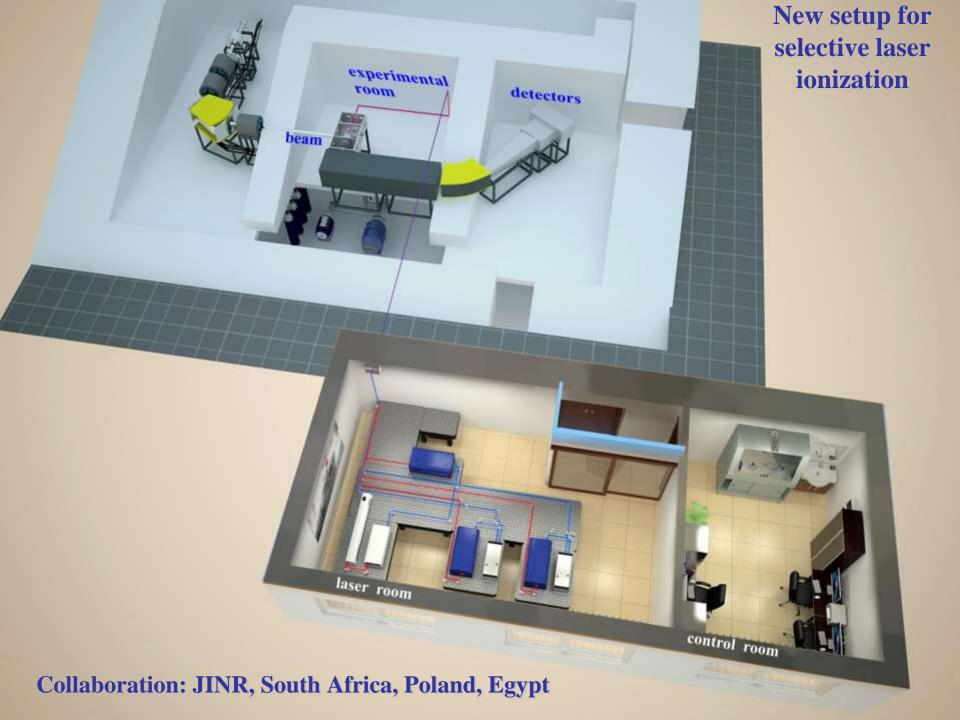


### Selective laser ionization! (Au & Hg as an example)



## New setup for selective laser ionization of multi-nucleon transfer reaction products stopped in gas





### Summary

- North-east part of the nuclear map is still "terra incognita". Heavy neutron rich nuclei are not synthesized and studied yet. It's time to develop these new lands.
- Multi-nucleon transfer reactions can be used for synthesis of new neutron enriched transfermium nuclei located along the beta-stability line. U-like beams are needed as well as a new kind of separators!
- Multi-nucleon transfer reactions can be also used for synthesis of new neutron rich nuclei located along the closed neutron shell N=126 having the largest impact on the astrophysical r-process. Cross sections are higher than 1  $\mu$ b.
- Shell effects and dynamics of the "inverse quasi-fission processes" in heavy ion damped collisions should be studied much better.
- Selective laser separation of reaction products is a very promising tool.



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